

A MATHEMATICAL MODEL FOR THE STARTING PROCESS OF A TRANSONIC LUDWIG TUBE WIND TUNNEL

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) A simplified mathematical model is presented for the unsteady flow process of starting a transonic Ludwig tube wind tunnel. The hardware modeled consists of a porous-walled test section surrounded by a plenum chamber with an exhaust system independent of the tunnel's main starting valves, which are located downstream of the diffuser-test section. In the present method, the hardware is modeled as three control volumes: the plenum, the test section,		

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20. ABSTRACT (Continued)

and the diffuser. The plenum is treated with the unsteady integral continuity equation with one-dimensional influx or outflux through the porous wall, through the plenum exhaust system, and through the flaps, which exhaust into the diffuser. The other two control volumes are treated with the steady integral continuity equation and a steady, adiabatic, one-dimensional energy equation whose stagnation conditions vary in time according to the classical solution for an unsteady expansion wave. Numerical solutions are compared with experimental pressure-time histories of a small, transonic, high Reynolds number tunnel referred to as HIRT. Agreement between the model and experiment is good.

PREFACE

The work reported herein was conducted by the Arnold Engineering Development Center (AEDC), Air Force Systems Command (AFSC), under Program Element 65807F. The results were obtained by ARO, Inc. (a subsidiary of Sverdrup & Parcel and Associates, Inc.), contract operator of AEDC, AFSC, Arnold Air Force Station, Tennessee. The research was done under ARO Project No. V37A-32A in support of the High Reynolds Number Wind Tunnel (HIRT) project. The author of this report was Frederick L. Shope, ARO, Inc. The manuscript (ARO Control No. ARO-VKF-TR-75-147) was submitted for publication on September 26, 1975.

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1.0 INTRODUCTION

This report documents an effort to mathematically model the aerodynamics involved in the unsteady process of starting a Ludwig Tube wind tunnel. In essence, the model represents the end product of many people assimilating a large amount of experimental data obtained from a transonic Ludwig tube facility and, thus, depends on several experimentally derived parameters and assumptions. The wind tunnel configuration studied here consists of a very long, circular supply tube which contracts to a rectangular, porous-walled test section. The test section expands through a diffuser into a valve manifold. Surrounding the test section is a plenum chamber with exhaust valves which can be controlled independently of the main valves. In addition, the plenum contains a set of ejector flaps which allow the plenum to exhaust itself into the diffuser.

When one considers that larger scale transonic Ludwig tube facilities would have a price of order \$10,000,000 and would produce a usable run time of only a few seconds per run, it is clear that considerable effort must be concentrated to ensure that the tunnel can be started rapidly under a wide range of operating conditions. A laboratory scale pilot facility (Ref. 1) (known as "Pilot HIRT") at Arnold Engineering Development Center provides an experimental vehicle to measure the effects of many of the important parameters in the tunnel starting process and to provide basic experimental data for verification of math models.

To clarify the need for a mathematical model of starting such a device, a brief explanation of the tunnel operation is required. Prior to a run, the tunnel is pumped to the desired charge pressure and temperature. A tunnel run is initiated by first opening the main valves downstream of the diffuser. This opening process sends unsteady expansion waves up the tunnel to the supply tube. Were it not for the plenum, the flow in the test section would become steady soon after the trailing edge of the unsteady wave from the valve, initiated by the valve area becoming steady, passed the test section into the supply tube. The test section flow cannot become steady until the plenum volume has been exhausted to the point where the summation of mass flow across the porous wall, through the flaps, and out the plenum exhaust (dumped to atmosphere) becomes zero and allows the plenum pressure to become steady. Since current state-of-the-art, fast-opening valves easily reach the required flow area in advance of the plenum becoming steady, the plenum is the primary limitation upon how quickly the tunnel can be started and steady flow established in the test section.

The present model assumes that the unsteady expansion wave emanating from the main valves propagates instantaneously to all parts of the wind tunnel and that property variation within the wave at any location in the diffuser, test section, nozzle, or supply tube is totally controlled by the area-time curve of the main valve. While partially retaining

the effect of the unsteady wave, this assumption allows use of the steady continuity equation in the test section coupled with the well-known exact solution for one-dimensional, variable area, isentropic flow (Ref. 2). Use of these equations at any instant requires a knowledge of stagnation conditions driving the flow, which vary through the nonisentropic expansion wave. Variation of the stagnation properties is computed via the exact solution for a one-dimensional unsteady wave in a variable area duct (Ref. 3). The unsteadiness of the plenum is handled via the unsteady continuity equation by equating the rate of mass accumulation in the plenum to the summation of all the flow rates entering and leaving the plenum. The air in the plenum is assumed to be a calorically perfect gas and its temperature is assumed either isentropic or equal to the stagnation temperature of the flow in the test section (whichever is greater), an experimentally based assumption. The main valves are treated as one-dimensional sonic orifices driven by the stagnation pressure and temperature of the unsteady wave. The plenum exhaust valves are handled similarly by assuming that the flow in the plenum is stagnant. Flow through the ejector flaps and across the porous wall is computed via an adaptation of the work of Ref. 4, which empirically corrected the flow rates with the pressure drops across these devices.

In the discussion which follows, the mathematical model will first be presented, including a more detailed description of the physical situation, the assumptions underlying the model, the mathematical formulation, and the solution procedure. Next, the model will be compared with a sample of experimental data from the Pilot HIRT facility. The appendixes contain some of the mathematical details and a brief user's manual for the computer program.

2.0 THE MATHEMATICAL MODEL

2.1 DESCRIPTION OF THE PHYSICAL SITUATION TO BE MODELED

All of the essential features of the proposed HIRT facility which are to be modeled are given in Fig. 1. The overall length of the facility is 1,880 ft, and the supply tube has an inside diameter of 15 ft. The main valve system consists of a number of fast-acting valves, and the plenum exhaust also requires a multiple valve system. The pilot facility provides a precisely scaled (1/13) flow envelope but has a single sliding sleeve valve in place of the valve manifold of the full-scale tunnel and a single plenum exhaust valve fed by multiple tubes from the plenum.

A tunnel run is initiated by opening the main valves and possibly the plenum valves, not necessarily together or in the same length of time. Both sets of valves send nonisentropic expansion waves throughout the tunnel and primarily up the charge tube. The main valve system produces the steepest (or strongest) wave because it handles a much greater portion

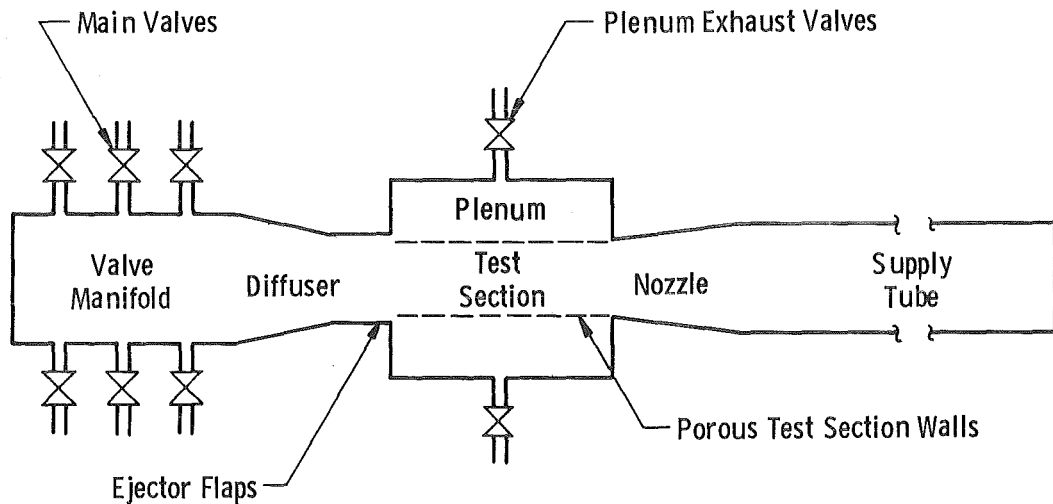


Figure 1. Major components of the High Reynolds Number Wind Tunnel.

of the flow rate than the plenum exhaust. At any point in the supply tube, the gas remains totally stagnant until the first expansion wave reaches that point; and the flow at that point does not become steady until the last expansion wave passes the point. The main valve system sends out its last expansion wave when the flow area becomes constant. The plenum also continues to send out expansion (and sometimes compression) waves until the plenum pressure becomes steady. But the plenum does not become steady until the sum of all the flows into and out of it are zero (Fig. 2), and it invariably controls the start time of the tunnel. Since the main valves are much faster than the plenum response, the pressure in the test section drops rapidly below the plenum pressure, causing mass flow to enter the test section from the plenum. As the plenum gradually catches up to the test section, the wall crossflow (across the porous test section wall) gradually decreases and, in some cases, reverses. This process, coupled with the increasing main valve area, gradually increases the flow rate drawn from the supply tube. However, the flow rate from the tube may continue to increase only until the nozzle exit becomes choked, after which the supply tube flow becomes steady since the choke point will no longer pass additional expansion or compression waves (unless the compression wave is strong enough to unchoke the nozzle). Whether the nozzle eventually chokes and whether the test section eventually steadies out to supersonic or subsonic flow depends on the relative flow areas of the main valves, the plenum exhaust valves, and the test section, the direction of the flap and wall crossflows, and how the various steady conditions are approached in time relative to each other. Subsonic and very slightly supersonic test section Mach numbers can be obtained without steady-state plenum exhaust, though the plenum exhaust may be opened temporarily and then closed in order to reduce the starting time. For subsonic flows, the steady main valve area - in terms of the ideal, one-dimensional area

at the choke point - must be as much less than the nozzle area (where the nozzle meets the entrance to the test section) as is dictated by the steady test section Mach number to be attained (neglecting diffuser losses). A slightly supersonic test section can be obtained with a steady main valve area greater than or equal to the nozzle area if the flaps and porosity are set properly, giving a flow situation as follows: with the nozzle choked and the plenum steadied at a pressure very near the static test section pressure such that the static pressure and dynamic heads of the main flow force a small crossflow into the plenum, the net test section flow decreases from the choked flow rate at the nozzle. The slightly subcritical flow rate leaving the test section thus produces a slightly supersonic condition, resulting in a favorable pressure gradient for the plenum to exhaust its incoming crossflow out the flaps and hence become steady. Normally, however, supersonic conditions (up to Mach 1.3 in the pilot) are obtained by having the plenum exhaust area become steady at a flow area sufficient to pass all of the mass flow rate entering the plenum via wall crossflow and sometimes via reverse flap flow.

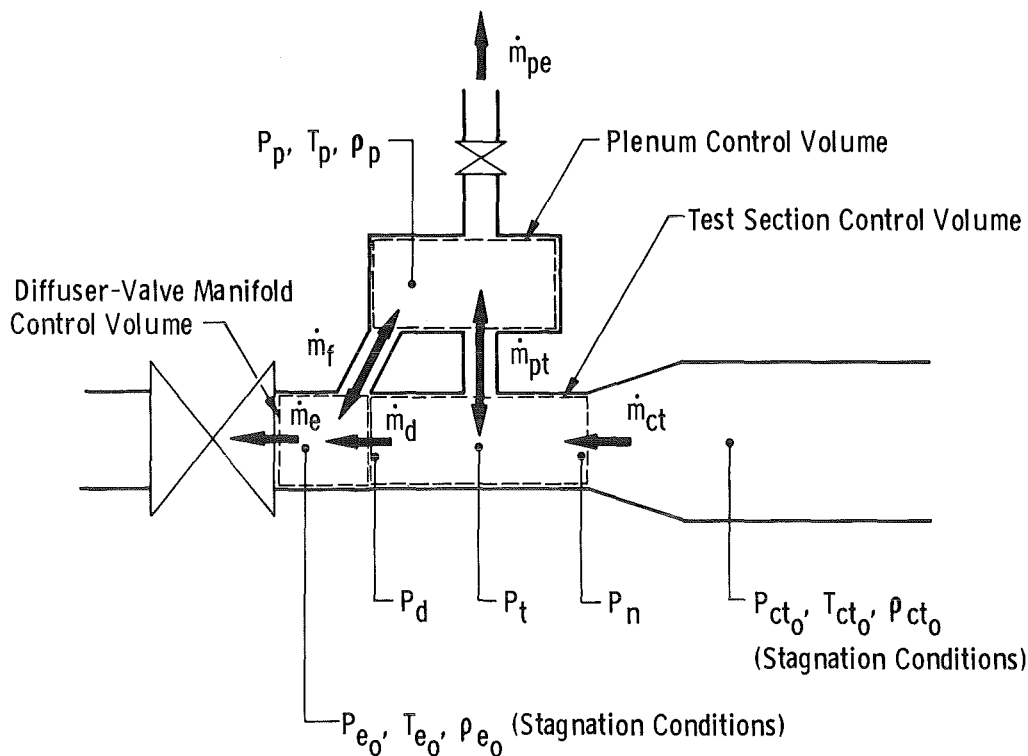


Figure 2. Schematic illustration of the flow process during start.

To understand the flow in terms of the mathematical model, the various flow configurations might be best thought of in terms of the steady energy equation relating the local pressure to the mass flux (Fig. 3). Subsonic flows fall on the branch to the

right of the choke point, supersonic flows to the left. In general, all points in the tunnel are initially at point A, which corresponds to no flow. Higher flow rates with correspondingly lower static pressures are illustrated by movement from point A to B

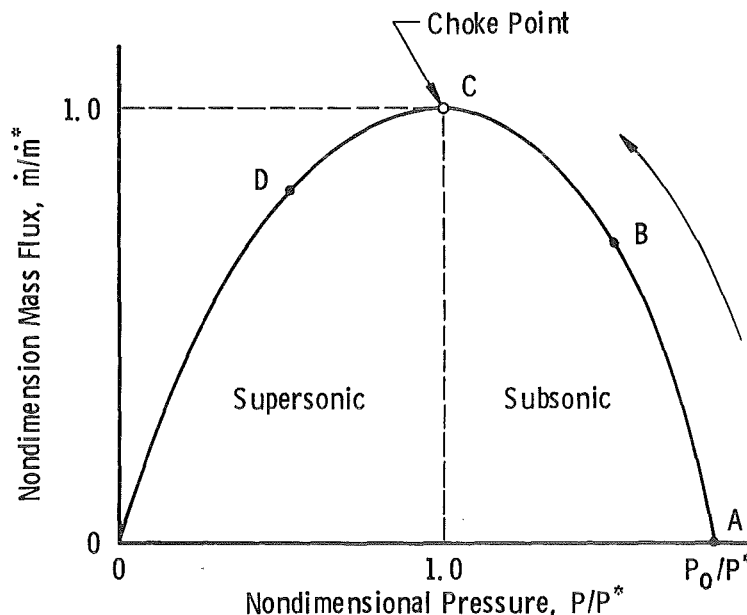


Figure 3. Qualitative plot of the energy equation.

on the energy equation. Flows which become subsonically steady would halt to the right of C; while for supersonic flows, some portions of the tunnel would proceed beyond C to D. If the energy dome is then plotted versus axial position in the test section as shown in Figs. 4, 5, and 6, the importance of the wall crossflow and the relative timewise approach of various components to their steady conditions may be made clearer. For a normal subsonic run, Fig. 4 shows the energy dome at the entrance and exit of the test section. The constant time contours are shown as straight lines for purposes of illustration, though in reality they would have to be nonlinear to some degree in order for all points on the contour to fall on the surface of the dome cylinder and because the wall crossflow does not necessarily vary linearly along the test section. As the flow begins, the constant time contours do not remain parallel because the flow rate leaving the test section will not balance that at the entrance, the difference being the wall crossflow. For the most probable case of the plenum lagging the test section pressure, the crossflow will be into the test section, giving a greater flow rate at the exit than at the entrance. As time proceeds, however, the plenum pressure eventually catches up to the test section so that the contours do become nearly straight and parallel as the crossflow becomes insignificant. This process assumes that the plenum exhaust, if opened, is eventually closed.

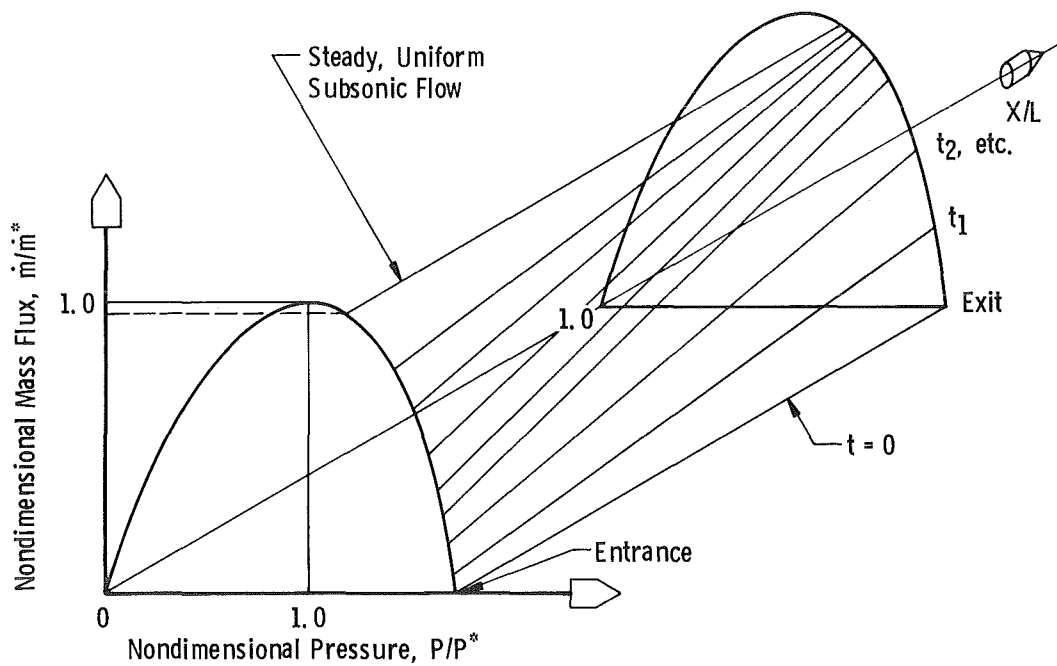


Figure 4. Energy dome versus position in test section for normal subsonic flow.

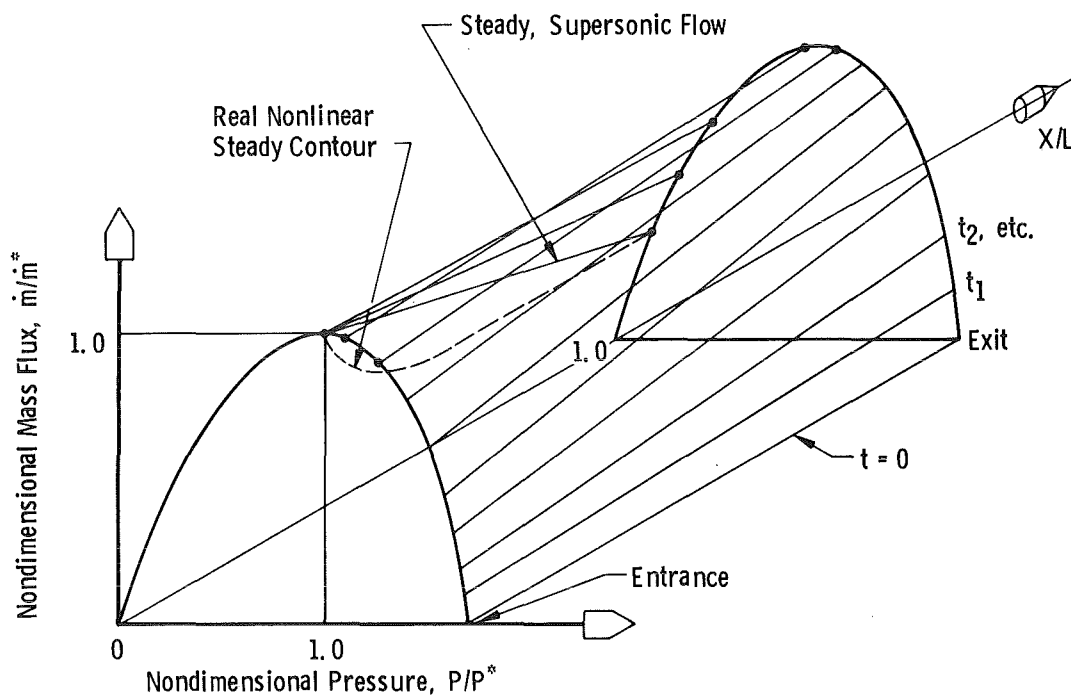


Figure 5. Energy dome versus position in test section for normal supersonic flow.

If the plenum exhaust is not closed and the steady main valve area is sufficiently large, the supersonic case of Fig. 5 may result. The initial constant time contours are similar to the subsonic case. However, the origins of the contours at the entrance eventually stop at the peak of the dome while at the exit they proceed over the choke point downward on the supersonic branch as the crossflow reverses from entering to leaving the test section. The contours, however well approximated by straight lines in the subsonic case, become significantly nonlinear for the higher supersonic Mach numbers, as illustrated by the dotted "real nonlinear steady contour" in Fig. 5. This results from a combination of the nonlinear variation of the wall crossflow and boundary layer growth. These nonlinear effects, though certainly present in the subsonic case, are more pronounced in the supersonic case because the pressure at the nozzle must remain unchanged at the choke value while the pressure at the exit varies significantly with the exit Mach number.

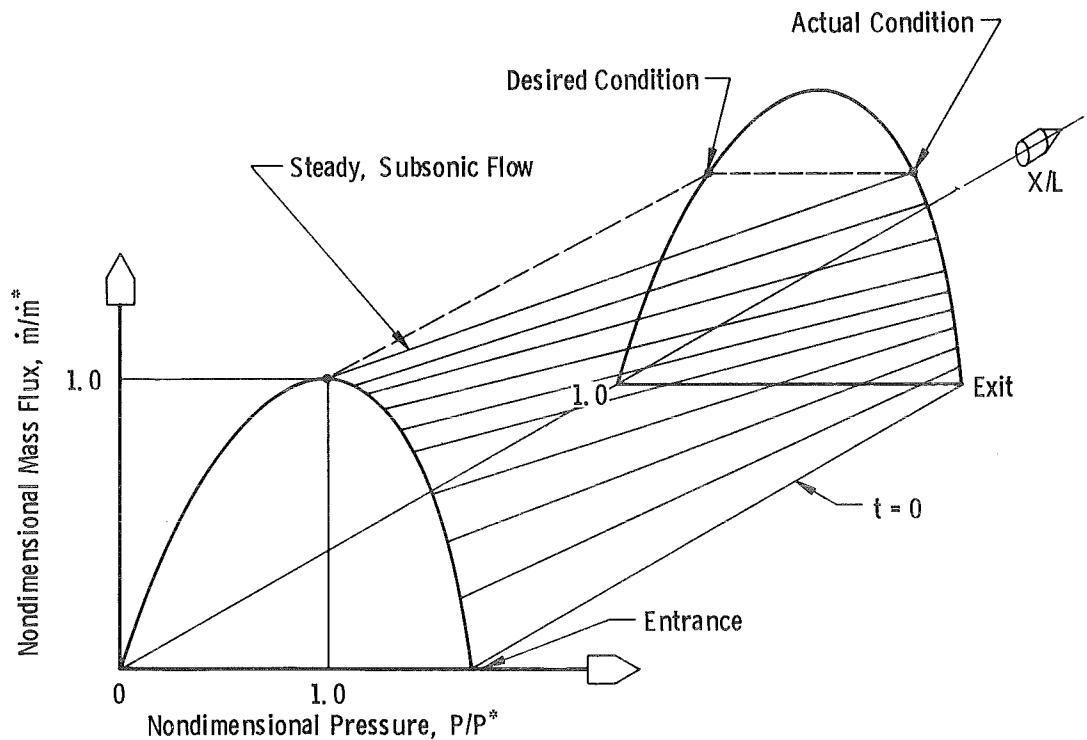


Figure 6. Energy dome versus position in test section for subsonic flow with choked nozzle.

The slopes of the constant time contours in Figs. 4 and 5 depend on the magnitude of the wall crossflow, which in turn depends partially on the pressure difference between the plenum and the test section. Since the timewise variation of the plenum pressure can be controlled by controlling the flow area-time curve of the plenum exhaust valves,

it appears that the shortest starting time for the tunnel would be obtained by controlling the plenum pressure to precisely follow the test section pressure so that the plenum would reach its steady conditions simultaneously with the main valve system. This would result in the constant time contours remaining parallel right up to their final position, or up to the choke point for a supersonic run. In Fig. 6, the plenum is exhausted fast enough so that the wall crossflow is always out of the test section, resulting in less flow rate leaving the exit of the test section than entering. Thus, the constant time contour at the entrance dome reaches its peak while the point on the exit dome is forced by the plenum exhaust to become steady before reaching the peak though the desired steady condition lies on the other side of the dome and cannot be reached. Hence, it appears that the manner in which the various portions of the tunnel approach their steady conditions in time relative to each other can affect the final outcome of a run.

The foregoing discussion of the test section flow in terms of the energy domes serves as an introduction to one of the key elements of the mathematical model, namely, the steady energy equation in an unsteady environment. The domes also provide graphic visualization for the flow process.

2.2 GOAL OF THE MODELING

The purpose of this mathematical model is to study the starting process, controlled by the plenum, in order to size the plenum exhaust system; determine the effect upon start time of the interaction of the area-time curves of the main valves, flaps, and plenum exhaust; and, in general, to provide the essential information necessary for trading off facility cost and start time. To provide this information, the model must accept the following input data. The gross level mass flow rate depends upon the cross-sectional area of the supply tube and nozzle exit. The geometric factor, on which the wall crossflow primarily depends, is the porosity, the fraction of the total surface area of the test section walls drilled out to allow flow between the test section and plenum. Thus, the dimensions of the test sections and porosity must be provided along with the experimentally derived coefficients for the flow model. A key design parameter having first-order impact on the start time is the plenum volume ratioed to the test section volume. The area-time curves of the main valves, plenum exhaust valves, and the flaps are required along with the experimental coefficients for the flap flow model. Finally, the characteristics of the gas must be provided in terms of the ideal gas constant and the specific heat ratio.

This input to the model is then used to compute the following data concerning the flow. As functions of time the static and stagnation properties - pressure, density, and temperature - along with mass flow rate and Mach number are computed for three stations along the tunnel circuit: the supply tube at the nozzle entrance, the test section entrance,

and the test section exit. The plenum properties along with the mass flux through the porous wall, flaps, plenum exhaust, and main valves are computed as functions of time.

There are many other considerations, neglected herein, which might be of interest for other applications. One of the most important is the boundary layer, whose growth on the walls of the supply tube and test section varies with time. This unsteadiness occurs because at any given station along the tunnel, the particles of air passing that station at succeeding times into the run have travelled over successively longer lengths of tube from their starting points. If the effect of the boundary-layer growth on the local mass flow rate is thought of in terms changing the effective flow area, one might suspect that the test section would never become steady. In reality, however, the boundary-layer growth, sufficiently late in the starting process, varies with approximately the same proportion in the nozzle and test section so that, though the effective flow areas may be varying, the area ratios (A/A^*) are not. As experimentally documented in Refs. 1 and 5, this results in essentially constant Mach number once the plenum has become steady, thus justifying the neglect of the boundary layer herein.

Neglect of the boundary layer means that no prediction is made of property variation over the cross section of the flow area. Similarly, detailed variation of properties along the length of the test section is not predicted. Such information would be useful for studying wall loading or flow uniformity but is of secondary importance for present purposes. Very severe nonuniformity occurs in the diffuser section (connecting the test section and main valve manifold), which has been subjected to a detailed experimental study in Ref. 4. The complexity of the diffuser flow results from a combination of effects: shock waves, flow separation, flap exhaust, and the presence of the model or probe support sector. The performance of the diffuser is important because of its effect on the noise environment in subsonic flow in the test section and because its stagnation losses significantly impact the sizing of the real flow area of the main valve system. However, for purposes of the starting model, diffuser losses may be neglected if the main valve area is assumed to be the ideal, one-dimensional flow area needed to pass a given mass flow rate for a given set of driving stagnation conditions as determined from wave mechanics.

Three additional effects neglected herein deserve mention. First, wave spreading is neglected. This phenomenon is due to the difference in propagation speed between the leading and trailing edges of the unsteady wave. Since the wave propagation speed (equal to the local speed of sound minus the local velocity) is less for the trailing edge than the leading edge, the time delay between a change in main valve area and the sensing of this change in the supply tube is greater for the last area change than the first. In fact, this delay is different for each position along the tunnel. However, over the greatest

distance of importance in the pilot facility, this difference in delay is less than 0.5 msec and is neglected in the model. Besides wave spreading, the model also neglects the finite time required for a disturbance to travel from one point to another. Such a consideration is important for determining the relative times for first motion of main valves and plenum exhaust valves; but for purposes of the starting model, the tunnel components determining start time - plenum, test section, and supply tube exit - are sufficiently close together that the propagation times (on the order of one millisecond in the pilot) are small compared with the starting time under study. However, neglect of the propagation time and wave spreading should not be construed to mean that the finite wave width is neglected. This width, or time difference between passage of a given point of the leading and trailing edges of the wave, depends primarily on the opening time of the valve but is also increased by the nonideal flow processes in the diffuser. Such effects are accounted for herein by correction of the area-time curve of the main valve. A final additional effect, accounted for empirically but not modeled in detail, is the nonisentropicity of the thermodynamics of the plenum. It has been experimentally observed that the temperature in the plenum approximates an isentropic process only during the initial portion of the starting process, but over the entire start time for the tunnel, the asymptotic plenum temperature is much closer to the stagnation temperature in the test section than that for a completely isentropic expansion. A good model of this process would have to include the mixing of the virgin plenum air with that entering from the test section as well as account for the heat transfer from the walls of the plenum. This possible refinement to the present model is not yet included.

2.3 FORMAL ASSUMPTIONS

Before proceeding to the equations comprising the mathematical model, the following list of assumptions should be reviewed:

- a. Flow across all control volume surfaces is one dimensional.
- b. The fluid is assumed to be a calorically perfect gas (constant specific heats).
- c. Flow within the envelope comprised of the supply tube, test section, and main valves is inviscid, adiabatic, and irrotational except as accounted for by the unsteady wave equations.
- d. Within this envelope and at a constant time, property variation from point to point is isentropic. Entropy variation with time is governed by the wave equations. Thus, at any given instant, the one-dimensional, variable area, isentropic equations of gas dynamics (Ref. 2) are applicable.

- e. Wave propagation time and wave spreading are zero. This justifies the steady assumption needed to invoke the equations of Ref. 2.

2.4 MATHEMATICAL FORMULATION

The set of equations comprising the model naturally divides into two groups, one for subsonic flow and one for supersonic flow. Since the set of equations for supersonic flow is nearly an exact subset of the subsonic case, the latter will be presented first, followed by a discussion of the changes needed for supersonic flow. The subsonic model is in the form of 19 algebraic equations, not necessarily linear, involving 19 unknowns. This system of equations must be solved numerically at successive points in time until all properties have approached their asymptotic values. The solution at any time t depends entirely on the property values obtained for the solution at $t - \Delta t$, a short time earlier, as well as the given valve area-time curves, which may be thought of as forcing functions. Quantities which vary between $t - \Delta t$ and t are usually evaluated at an intermediate time t^* such that $(t - \Delta t) < t^* < t$. The time t^* is usually taken as the midpoint of the time interval.

The model is based on mass conservation for three control volumes as illustrated in Fig. 2. Conservation of mass for the plenum is derived from the unsteady integral continuity equation for a control volume to give

$$\rho_p(t) = \rho_p(t - \Delta t) + [\dot{m}_{pt}(t^*) + \dot{m}_{pe}(t^*) + \dot{m}_f(t^*)] \frac{\Delta t}{V_p} \quad (1)$$

Here ρ_p is the mass density in the plenum, assumed uniform throughout, and V_p is the volume of the plenum. The quantities $\dot{m}_{pt}(t^*)$, $\dot{m}_{pe}(t^*)$, and $\dot{m}_f(t^*)$ represent, respectively, the mass flow rates between the plenum and test section (pt), out the plenum exhaust (pe), and through the flaps (f). The formal continuity equation can not be precisely integrated because the dependence of the mass flow rates on t or ρ_p can not be written in simple closed form. However, the law of the mean provides that $\rho_p(t)$ may still be precisely computed if the flow rates are treated as constant but evaluated at a suitable intermediate point t^* . If Δt is now chosen sufficiently small so that the flow rates may be suitably approximated by linear functions of time, t^* can obviously be chosen as $t - 1/2\Delta t$, the midpoint. For the other two control volumes in Fig. 2, the steady continuity equation is used, having been justified by assumption (e) of the last section. By noting that Eq. (1) assumes the flap and wall crossflows are positive when flow is into the plenum, continuity for the test section becomes

$$\dot{m}_{ct}(t^*) = \dot{m}_{pt}(t^*) + \dot{m}_d(t^*) \quad (2)$$

and for the diffuser-valve manifold control volume

$$\dot{m}_d(t^*) = \dot{m}_e(t^*) + \dot{m}_f(t^*) \quad (3)$$

The three new mass flow rates introduced here are, in terms of the subscripts, that leaving the supply tube (ct, for charge tube, as it is often called), the primary tunnel exit (e) provided by the main valves, and the diffuser-end (d) of the test section. It should be noted from Fig. 2 that \dot{m}_d corresponds to a point upstream of where the flap flow enters the main stream.

Proceeding next to model each of these six mass flow rates, consider first the flow through the plenum exhaust and the main valves, which are both treated as single one-dimensional sonic orifices driven by the stagnation conditions.

$$\dot{m}_e(t^*) = a \frac{P_{e_o}(t^*) A_{e_o}(t^*)}{\sqrt{T_{e_o}(t^*)}} \quad (4)$$

$$\dot{m}_{pe}(t^*) = a \frac{P_p(t^*) A_{pe}(t^*)}{\sqrt{T_p(t^*)}} \quad (5)$$

In Eq. (4), P_{e_o} and T_{e_o} are the stagnation pressure and temperature in the valve manifold; and in Eq. (5), P_p and T_p are the pressure and temperature in the plenum, approximated as stagnant. The quantities A_e and A_{pe} are the total flow areas of the main valves and plenum exhaust valves. These areas are assumed to be the ideal, one-dimensional flow areas of a sonic orifice. If the real valve areas are used, discharge coefficients must be included in Eqs. (4) and (5). The constant a is given by

$$a = \sqrt{\left(\frac{2}{\gamma+1}\right)^{\frac{\gamma+1}{\gamma-1}} \frac{\gamma}{R}} \quad (6)$$

where R is the ideal gas constant and γ is the ratio of the specific heats.

Consider next the flap and wall crossflows, which have been neatly modeled by Varner (Ref. 2) as simply proportional to the pressure drop across the devices. With a second order adaptation added here, Varner's model takes the following form

$$\dot{m}_{pt}(t^*) = - \frac{A_w}{k_w} [P_p(t^*) - A_{15} P_t(t^*)] \quad (7)$$

$$\dot{m}_f(t^*) = - \frac{A_f(t^*)}{k_f} [P_p(t^*) - A_{16} P_d(t^*)] \quad (8)$$

Here A_w and A_f are the effective flow areas through the porous wall and through the flaps. While A_f is the actual geometric area, A_w depends on the total surface area of the test section walls (A_{tsw}), the porosity (τ), and a flow coefficient. Varner gives this relationship as

$$A_w = 0.17 \tau A_{tsw} \quad (9)$$

The flow coefficients k_w and k_f were determined by Varner from experimental data from Pilot HIRT and are given in Fig. 7. The values of k_w in Fig. 7 are for the porosity shown in Fig. 8. The coefficients¹ A_{15} and A_{16} multiplying, respectively, the mean test section pressure P_t and the diffuser end test section pressure P_d were added in an effort to improve the accuracy of the asymptotic values of the numerical solution. The rationale for each of these constants is different. Rigorous modeling of the crossflow must include not only the effect of pressure forces but also the momentum of the fluid as it moves along the test section wall. The coefficient A_{15} thus represents an attempt to include momentum effects as a small correction to the existing crossflow model. Experimental evidence from the pilot facility indicates that this small momentum effect can make the difference between choking and not choking when the desired steady conditions are very near sonic flow. In particular, it has been observed that during supersonic flow, where the net crossflow must be from the test section to the plenum, the test section pressure is actually slightly less than the plenum pressure.

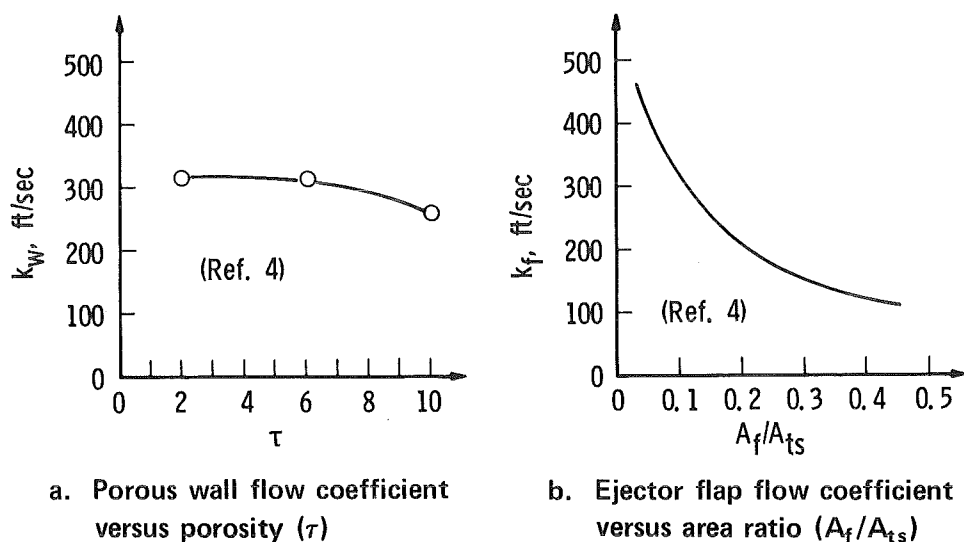


Figure 7. Porous wall and flap flow coefficients.

¹The subscripts 15, 16, and 17 have no significance beyond consistency with variable names in the computer program.

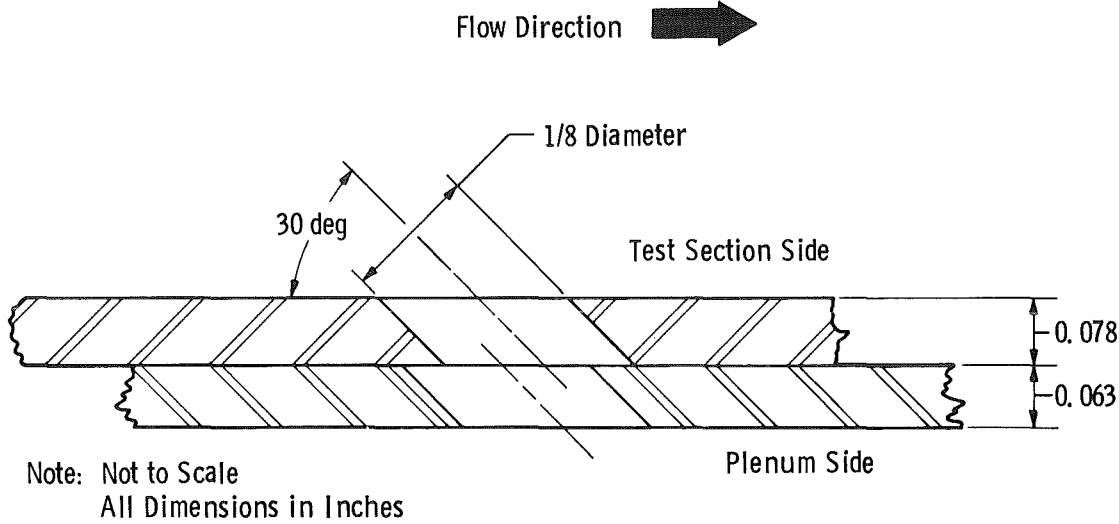


Figure 8. Wall porosity in Pilot HIRT.

This has been attributed to the fluid momentum in the test section overcoming the slightly adverse pressure gradient. The other constant, A_{16} in the flap model, was added to account for some of the losses in the upstream portion of the diffuser. Unfortunately, both of these constants were found to be functions of the test section Mach number, thus indicating the need for more accurate modeling.

The mean test section pressure P_t in Eq. (7) is computed from a weighted average of the pressure at the nozzle-end of the test section P_n and at the diffuser end P_d . That is,

$$P_t(t^*) = (1 - A_{17})P_n(t^*) + A_{17}P_d(t^*) \quad (10)$$

where $0 \leq A_{17} \leq 1$. Since a detailed model of axial property variation in the test section has not yet been included in the start model, properties are computed only at the nozzle and diffuser ends of the test section. For subsonic flows, the value of A_{17} was not found critical to the accuracy of the solution and was thus taken as 0.5, assuming a linear variation. For supersonic flow, a value of 0.9 was used to account for the more pronounced axial gradients.

The remaining two mass flow rates (\dot{m}_{ct} and \dot{m}_d) may be related to pressures already introduced above using the steady energy equation discussed earlier and shown in Fig. 3. At the diffuser end of the test section, the energy equation is

$$\left[\frac{\dot{m}_d(t^*)}{\dot{m}_o(t^*)} \right]^2 = \frac{2}{\gamma - 1} \left\{ \left[\frac{P_d(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{2}{\gamma}} - \left[\frac{P_d(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{\gamma+1}{\gamma}} \right\} \quad (11)$$

where as before the subscript "ct" refers to the charge tube and the subscript "o" indicates stagnation properties. The quantity \dot{m}_o is defined as

$$\dot{m}_o(t^*) \equiv \sqrt{\frac{\gamma}{R}} \frac{P_{ct_o}(t^*)}{\sqrt{T_{ct_o}(t^*)}} A_{ts} \quad (12)$$

where A_{ts} is the cross-sectional area of the test section. The stagnation properties (P_{ct_o} and T_{ct_o}) are thought of as originating from the unsteady wave when it reaches the charge tube and are assumed the same, for any given time, throughout all of the flow envelope except the plenum. At the nozzle end of the test section, the flow rate is equal to that in the charge tube, since its value has not yet been modified by any wall crossflow. At this station, the energy equation is, therefore,

$$\left[\frac{\dot{m}_{ct}(t^*)}{\dot{m}_o(t^*)} \right]^2 = \frac{2}{\gamma - 1} \left\{ \left[\frac{P_n(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{2}{\gamma}} - \left[\frac{P_n(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{\gamma+1}{\gamma}} \right\} \quad (13)$$

To complete the portion of the model not arising from the unsteady wave, the thermodynamic equations of state for the plenum are needed. To compute the properties at t^* for use in Eqs. (5), (7), and (8) while Eq. (1) gives the density at t , the density at t^* is computed from

$$\rho_p(t^*) = 1/2[\rho_p(t) + \rho_p(t - \Delta t)] \quad (14)$$

The plenum temperature is assumed equal to the greater of the isentropic temperature and the stagnation temperature in the test section.

That is,

$$T_p(t^*) = \max \left\{ T_p(t^* - \Delta t) \left[\frac{\rho_p(t^*)}{\rho_p(t^* - \Delta t)} \right]^{\gamma-1}, T_{ct_o}(t^*) \right\} \quad (15)$$

In either event, the pressure may then be obtained from the perfect gas law:

$$P_p(t^*) = \rho_p(t^*) R T_p(t^*) \quad (16)$$

Closing the system of equations presented so far requires relationships for how the stagnation properties vary in time. A careful accounting of the number of equations and the number of unknowns to this point would reveal that, given values of P_{ct_o} and T_{ct_o} and assuming $P_{e_o} = P_{ct_o}$ and $T_{e_o} = T_{ct_o}$ (which is what is done for the subsonic case),

it is possible to compute the value of \dot{m}_{ct} . This value of the flow rate from the charge tube represents that required by the sum total of all the expansion waves which at a given time have reached the charge tube from all parts of the tunnel. That is, \dot{m}_{ct} identifies an intermediate point within the entire unsteady wave, which begins with the first motion of a valve somewhere in the tunnel and ends when the plenum reaches its asymptotic pressure. Thus, \dot{m}_{ct} may be used to compute all other stagnation properties for that point in the unsteady wave. By using the equations of Ref. 3 and after some algebra, the charge tube Mach number at the desired point in the wave may be related to \dot{m}_{ct} by the equation:

$$\dot{m}_{ct}(t^*) = M_{ct}(t^*) \left[1 + \frac{\gamma-1}{2} M_{ct}^2(t^*) \right]^{-\frac{\gamma+1}{\gamma-1}} \dot{m}_c \quad (17)$$

where \dot{m}_c is defined from

$$\dot{m}_c = \sqrt{\frac{\gamma}{R}} \frac{P_c}{\sqrt{T_c}} A_{ct} \quad (18)$$

Here A_{ct} is the cross-sectional area of the charge tube, and P_c and T_c are the charge conditions, that is, the air pressure and temperature after the tunnel has been pumped up but before any valves are opened. These charge conditions are assumed to apply uniformly throughout the envelope, including the plenum. After obtaining the charge tube Mach number, the stagnation pressure and temperature are readily computed from the following equations from Ref. 3:

$$P_{ct_o}(t^*) = \left\{ \frac{1 + \frac{\gamma-1}{2} M_{ct}^2(t^*)}{\left[1 + \frac{\gamma-1}{2} M_{ct}^2(t^*) \right]^2} \right\}^{\frac{\gamma}{\gamma-1}} P_c \quad (19)$$

$$T_{ct_o}(t^*) = \left\{ \frac{1 + \frac{\gamma-1}{2} M_{ct}^2(t^*)}{\left[1 + \frac{\gamma-1}{2} M_{ct}^2(t^*) \right]^2} \right\} T_c \quad (20)$$

Equations (1) through (20) thus comprise the subsonic portion of the starting model and are summarized in Table 1. The supersonic case is physically different from the subsonic case and requires solution of a different set of equations as noted in Table 1. The distinguishing factor of the supersonic case is that the nozzle exit is choked, making the flow rate and stagnation conditions steady there. Once the nozzle chokes, the charge tube Mach number is a constant depending only on the area ratio between the charge tube and nozzle exit. From Ref. 2, the steady Mach number can be obtained by reverting the equation:

Table 1. List of Exact Simultaneous Equations

Equation	Independent Variable to be Computed	Included in Supersonic Case?	Text Equation Number	Program Equation Number
$\rho_p(t) = \rho_p(t - \Delta t) + [\dot{m}_{pt}(t^*) + \dot{m}_{pe}(t^*) + \dot{m}_f(t^*)] \frac{\Delta t}{V_p}$	$\rho_p(t)$	Yes	1	5
$\dot{m}_{ct}(t^*) = \dot{m}_{pt}(t^*) + \dot{m}_d(t^*)$	$\dot{m}_{ct}, M < 1$ $\dot{m}_d, M > 1$	Yes	2	6
$\dot{m}_d(t^*) = \dot{m}_e(t^*) + \dot{m}_f(t^*)$	\dot{m}_d	No	3	7
$\dot{m}_e(t^*) = a \frac{P_{e_o}(t^*) A_{e_o}(t^*)}{\sqrt{T_{e_o}(t^*)}}$	\dot{m}_e	No	4	1
$\dot{m}_{pe}(t^*) = a \frac{P_p(t^*) A_{pe}(t^*)}{\sqrt{T_p(t^*)}}$	\dot{m}_{pe}	Yes	5	2
$\dot{m}_{pt}(t^*) = - \frac{A_w}{k_w} [\dot{P}_p(t^*) - A_{15} \dot{P}_t(t^*)]$	\dot{m}_{pt}	Yes	7	4
$\dot{m}_f(t^*) = - \frac{A_f(t^*)}{k_f} [\dot{P}_p(t^*) - A_{16} \dot{P}_d(t^*)]$	\dot{m}_f	Yes	8	3
$\dot{P}_t(t^*) = (1 - A_{17}) \dot{P}_n(t^*) + A_{17} \dot{P}_d(t^*)$	\dot{P}_t	Yes	10	11
$\left[\frac{\dot{m}_d(t^*)}{\dot{m}_{ct}(t^*)} \right]^2 = \frac{2}{\gamma - 1} \left\{ \left[\frac{P_d(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{2}{\gamma}} - \left[\frac{P_d(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{\gamma+1}{\gamma}} \right\}$	P_d^a	Yes	11	12
$\left[\frac{\dot{m}_{ct}(t^*)}{\dot{m}_{ct_o}(t^*)} \right]^2 = \frac{2}{\gamma - 1} \left\{ \left[\frac{P_n(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{2}{\gamma}} - \left[\frac{P_n(t^*)}{P_{ct_o}(t^*)} \right]^{\frac{\gamma+1}{\gamma}} \right\}$	P_n^a	No	13	13
$\dot{m}_o(t^*) = \sqrt{\frac{\gamma}{R}} \frac{P_{ct_o}(t^*)}{\sqrt{T_{ct_o}(t^*)}} A_{ts}$	\dot{m}_o	No	12	14
$\rho_p(t^*) = 1/2[\rho_p(t) + \rho_p(t - \Delta t)]$	$\rho_p(t^*)$	Yes	14	18
$T_p(t^*) = \max \left\{ T_p(t^* - \Delta t) \left[\frac{\rho_p(t^*)}{\rho_p(t^* - \Delta t)} \right]^{\gamma-1}, T_{ct_o}(t^*) \right\}$	T_p	Yes	15	17
$P_p(t^*) = \rho_p(t^*) R T_p(t^*)$	P_p	Yes	16	19
$\dot{m}_{ct}(t^*) = M_{ct}(t^*) \left[1 + \frac{\gamma-1}{2} M_{ct}^2(t^*) \right]^{-\frac{\gamma+1}{\gamma-1}} \dot{m}_c$	M_{ct}^a	No	17	8
$P_{ct_o}(t^*) = \left\{ \frac{1 + \frac{\gamma-1}{2} M_{ct}^2(t^*)}{\left[1 + \frac{\gamma-1}{2} M_{ct}(t^*)^2 \right]^2} \right\}^{\frac{\gamma}{\gamma-1}} P_c$	P_{ct_o}	No	19	9
$T_{ct_o}(t^*) = \left\{ \frac{1 + \frac{\gamma-1}{2} M_{ct}^2(t^*)}{\left[1 + \frac{\gamma-1}{2} M_{ct}(t^*)^2 \right]^2} \right\} T_c$	T_{ct_o}	No	20	10
$P_{e_o}(t^*) = P_{ct_o}(t^*), T_{e_o}(t^*) = T_{ct_o}(t^*)$		No	-	-

^aRequire Numerical Reversion

$$\frac{A_{ct}}{A_{ts}} = \frac{1}{M_{ct}(t^*)} \left\{ \frac{2}{\gamma + 1} \left[1 + \frac{\gamma - 1}{2} M_{ct}^2(t^*) \right] \right\}^{\frac{\gamma + 1}{2(\gamma - 1)}} \quad (21)$$

With this final Mach number, the steady stagnation conditions (P_{ct_0} , T_{ct_0} , and \dot{m}_0) along with the steady charge tube flow rate (\dot{m}_{ct}) can be computed one final time from Eqs. (19), (20), (13), and (17), after which these equations and variables may be dropped from the simultaneous solution. Since \dot{m}_{ct} is now constant, the flow rate leaving the test section (\dot{m}_d) is solely dependent on the wall crossflow (\dot{m}_{pt}) according to Eq. (2) and is independent of the flow rate out the main valves (\dot{m}_e), assuming the valve area A_e is sufficient to pass all the charge tube flow not removed by the plenum exhaust. Thus, Eqs. (3) and (4) may also be dropped from the system of equations. This is fortunate since it is no longer true that $P_{e_0} = P_{ct_0}$, which results from the nonisentropic recompression of the supersonic flow entering the diffuser. Thus, the original system of 19 equations and 19 unknowns reduces to 10 equations and 10 unknowns for the supersonic case.

These two sets of equations were solved using an iterative technique which unfortunately failed to converge in the vicinity of the choke point in time. To provide an alternate solution procedure when the iterative technique failed to converge, a small perturbation solution was developed for the original exact equations. The small perturbation solution was then used as an initial guess for the iterative procedure when it converged and as the complete solution when it did not. The results of this lengthy derivation are recorded in Appendix A, but the essential ideas are discussed below.

The exact solution already assumes that Δt is a small quantity. For the small perturbation solution, therefore, any of the 19 variables at time t^* may be assumed to be related to their values at $t^* - \Delta t$ by the general form

$$v_i(t^*) = v_i(t^* - \Delta t) + \epsilon_i(t^*) \quad (22)$$

where ϵ_i is the small increment in the variable and $i = 1, 2, \dots, 19$. If these small perturbation equations are used to expand the original exact equations, a new system of equations involving the increments rather than the variables themselves is obtained. For all exact equations, except the energy equations relating the pressure and mass flux at the entrance and exit of the test section (Eqs. (11) and (13)), only terms of order ϵ_i need be retained in the small perturbation equations. Such is not the case for the energy equations because in the region of the peak (or choke point) in Fig. 3, there is no linear approximation to the function. In the expanded equation, the coefficient of ϵ_i approaches zero as the Mach number approaches one. Thus, the term of order ϵ_i^2 , whose coefficient is nonzero at Mach number one, governs the form of the expansion. The resulting subsonic system

of equations is thus comprised of 17 linear equations and 2 second-degree equations, which can be solved analytically. The supersonic case is composed of 9 linear equations and 1 of second degree.

2.5 SOLUTION PROCEDURE

The procedure used to solve these two systems of equations is discussed in the following section. Included is a discussion of the overall logical procedure, the order in which the equations of the exact solutions are used, convergence considerations, and a general description of the computer program used to accomplish the calculation. The general solution procedure is illustrated by the flow chart in Fig. 9. The decision whether to use the supersonic or subsonic branch is decided by whether $P_d(t^* - \Delta t) < P^*$ or $P_d(t^* - \Delta t) > P^*$, that is whether the diffuser end of the test section was supersonic or subsonic at the midpoint of the previous time interval. If the previous interval was supersonic, the current one is also assumed to be supersonic. If the previous interval was subsonic but $1 - M(t^* - \Delta t) \leq M(t^* - \Delta t) - M(t^* - 2\Delta t)$, then the supersonic branch is used for the current time interval; otherwise the solution is assumed to remain subsonic. This criterion is checked for both ends of the test section, and the switch to the supersonic branch is contingent upon either or both positions satisfying the inequality. In either event, the small perturbation solution is computed to provide a good starting point for the exact iterative procedure. If convergence does not occur before a given number of iterations, the small perturbation solution is used as the final solution, and the next time interval is begun.

The "exact iterative procedure" referred to above is accomplished by taking an initial guess for one of the 19 variables and then proceeding from equation to equation, determining new values for each of the 19 variables until a complete circuit is made and a second value of the variable initially guessed at is obtained. This process is repeated until the difference between two successive values of certain of the variables is within a preset limit. For the subsonic case, the equation order is as follows:

4, 3, 11, 10, 7, 2, 17, 19, 20, 12, 13, 10, 7, 8, 1, 14, 15, 16,
5, ...

The supersonic equation order is

10, 7, 2, 11, 10, 7, 8, 1, 14, 15, 16, 5, ...

Some of these equations (Eqs. (11), (13), and (17)) require reversion from the form given but cannot be reverted analytically in closed form and must be solved numerically. The variable to be solved for in each equation is indicated in Table 1, and the three requiring numerical reversion are marked with an asterisk:

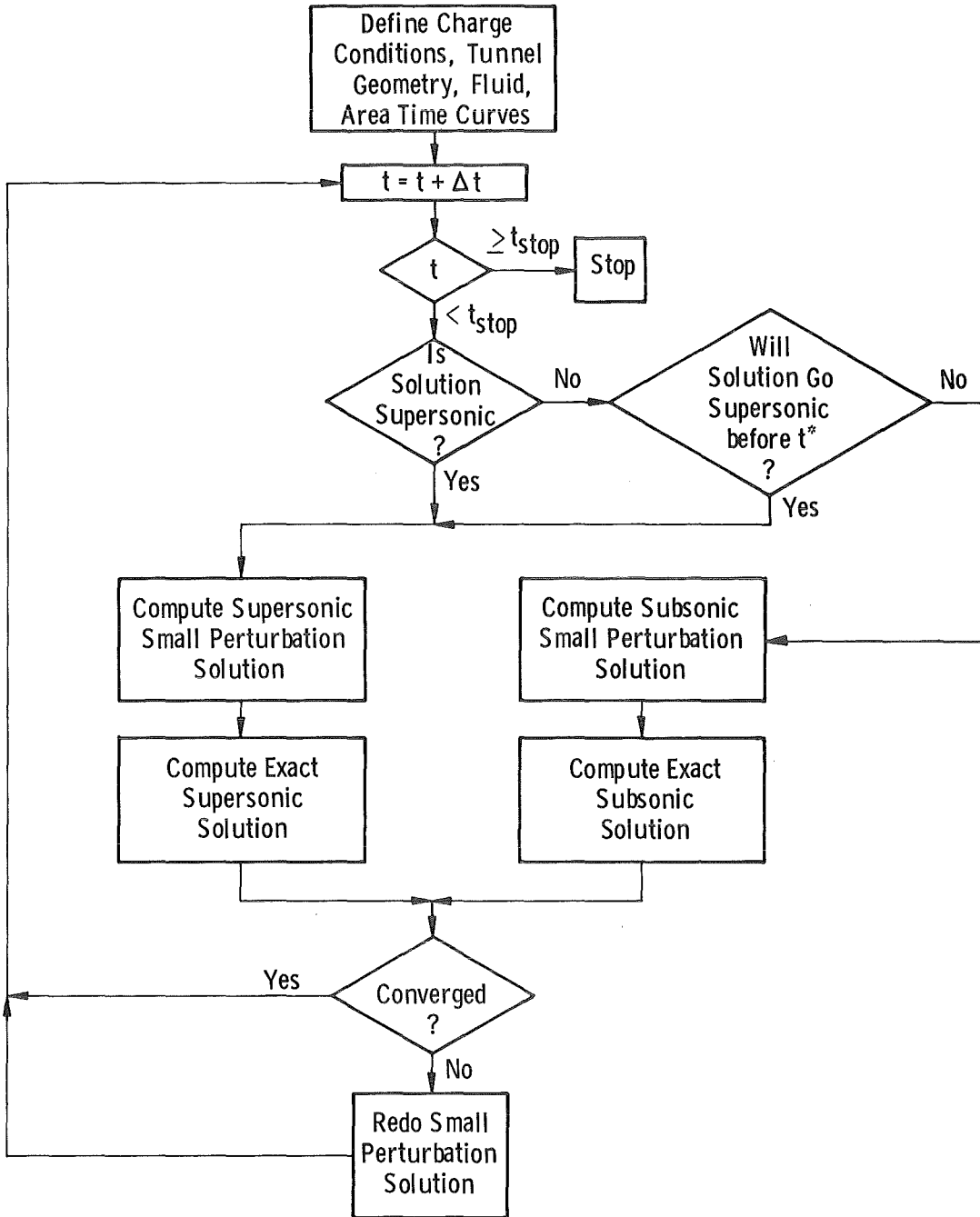


Figure 9. Flow chart of solution procedure.

The complete solution thus requires numerical iteration at three distinct levels, which necessitates careful consideration of convergence criteria as well as what to do when the criteria can not be met because of stability problems. The most basic level of numerical iteration involves reversion of the two energy equations and the mass flux - Mach number

wave equation. Considering the general case where the function $Y = F(X)$ must be solved for X given a value of Y and a guess X_1 , the procedure is simply to adjust X_1 in the direction which reduces the error criterion

$$E_1 = \frac{Y - F(X)}{Y} \quad (23)$$

until $|E_1| \leq |E_{\max}|$, E_{\max} being the present, maximum allowable error. The precise logic of the procedure is illustrated in the flow chart in Fig. 10. Since this procedure must

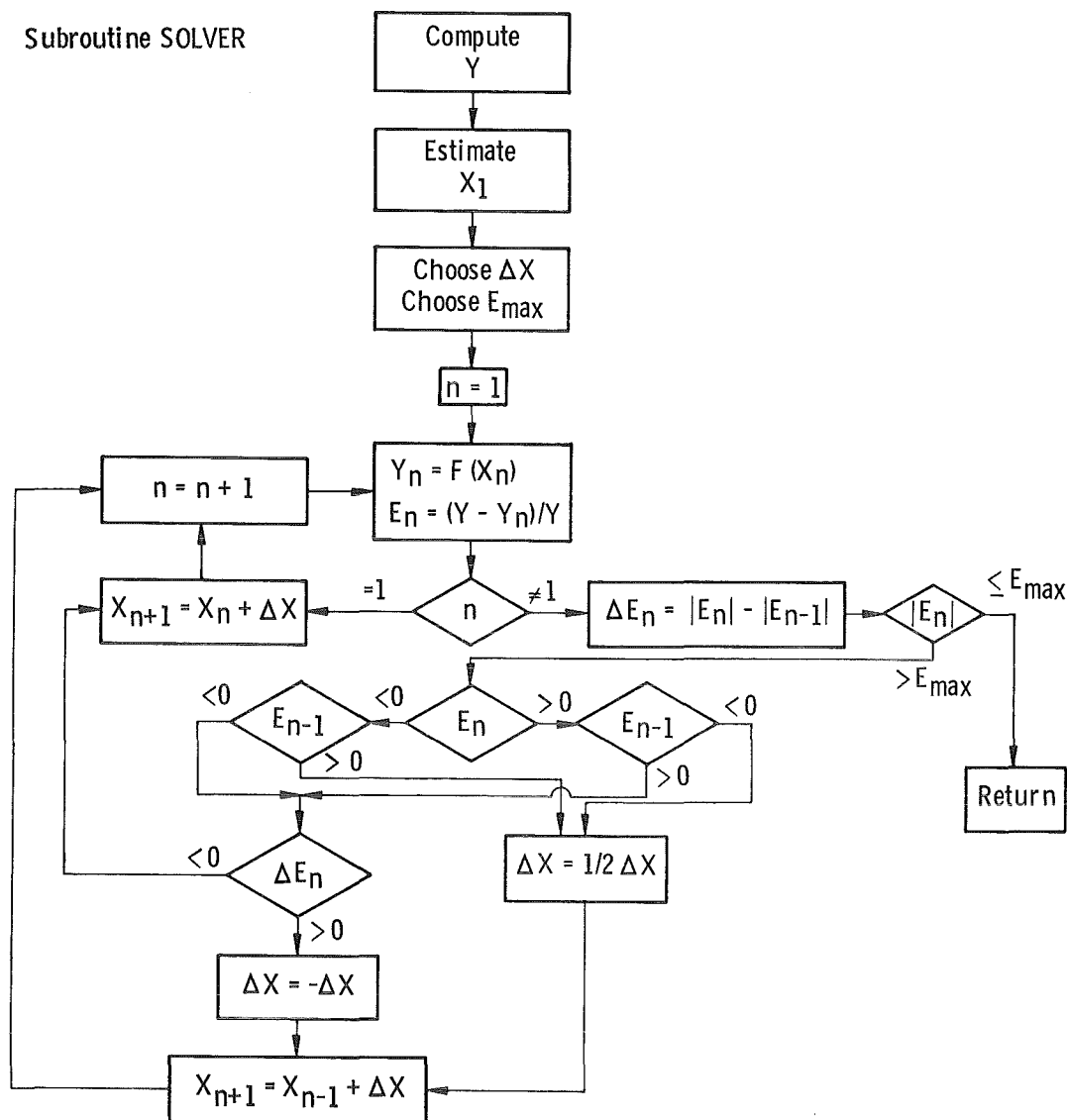


Figure 10. Logic for numerical solution of an algebraic equation $Y = F(X)$ for X , given Y when $X = F^{-1}(Y)$ is not a closed form function.

be repeated many times at each time interval, it is of considerable importance (because of impact on computer time) to achieve a solution with as few iterations as possible. Since the number of iterations depends to a large extent on the accuracy of the guess X_1 , considerable effort was expended in obtaining approximate reversions of the three equations. It was inadvertently discovered that the energy equation may be approximated with surprising accuracy over the entire range of present interest with a single ellipse, the reversion of which is trivial. The wave equation presented more of a problem. Since an easily revertable second-degree expansion around $M_{ct} = 0$ failed to match the accuracy of the elliptic energy equation, the expansion was carried to the seventh degree and then formally reverted according to the procedure of Ref. 6. These expansions are summarized in Appendix B.

The next higher level of iteration is, of course, the simultaneous solution of the exact model equations, during which stability problems were encountered in the vicinity of the choke point. The error criterion for halting the iteration may be generally expressed as

$$\left| \frac{v_i^{(n)} - v_i^{(n+1)}}{\frac{1}{2} [v_i^{(n)} + v_i^{(n+1)}]} \right| \leq P_{err} \quad (24)$$

where test variables (v_i) are the pressures P_n , $P_p(t)$, $P_p(t^*)$, P_d , P_t , and P_{ct0} ; P_{err} is the maximum allowable error; and n is the iteration number. Figure 11 illustrates the stability problem encountered in striving to meet this error limit. Shown is how the plenum pressure $P_p(t^*)$ varied with iteration number at two succeeding time points, one converging and one not. Such stability problems are known to occur in applying the iterational technique to locating the intersection of two curves on a plane when the curves have the same slope (same or opposite sign) at the point of intersection. Whether this simple explanation in 2-space is applicable to 19-space where no two of the 19 functions lie in the same plane is unclear. In any event, improvement in convergence rate was sought via the following procedures, most of which improved the situation:

- a. Relative Errors. It was found that if E_{max} was much greater than $1/10 P_{err}$, the numerical reversions could oscillate enough themselves from one iteration to the next to slow convergence.
- b. Computational Precision. Single precision arithmetic (~ 8 digits on an IBM 370) was found inadequate to achieve errors of $E_{max} = 10^{-5}$ ($P_{err} = 10^{-4}$), and double precision (~ 16 digits) was, therefore, adopted.

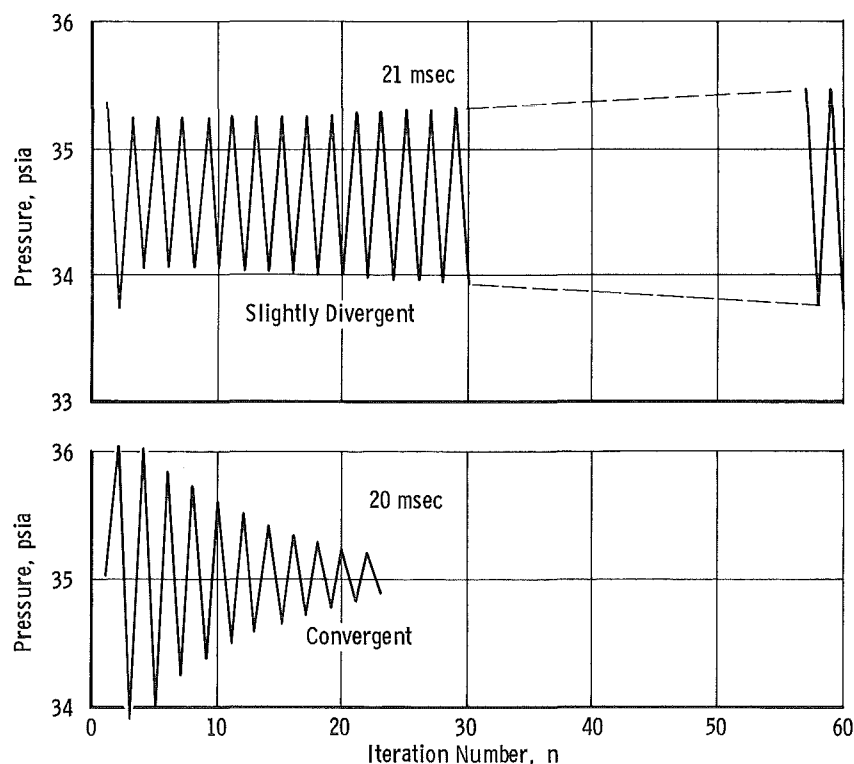


Figure 11. Plenum pressure versus iteration number for convergent and divergent cases.

- c. **Solution Weighting.** The clearly periodic oscillation of Fig. 11 suggests that the average of any two successive values should be closer to the final asymptote than either value. Accordingly, solution weighting,

$$v_i^{(n)} = A_{11}v_i^{(n)} + (1 - A_{11})v_i^{(n+1)} \quad (25)$$

was employed on a regular basis.

- d. **Weight Cutting.** It was further discovered that convergence rate could be greatly improved after the number of iterations reached a certain point if a lesser weight was applied to the current value $v_i^{(n)}$.
- e. **Error Cutting.** It was found that, later in a computation when some of the pressures were very near their asymptotes, the amount of variation from one time point to the next eventually approached the error limit. This in effect allowed these values to vary at random within the error limits and deteriorate the convergence rate. It was thus found prudent to reduce the error limits as necessary so as to maintain

$$P_{err} \leq \left| \frac{v_i(t^*) - v_i(t^* - \Delta t)}{\frac{1}{2}[v_i(t^*) + v_i(t^* - \Delta t)]} \right| \quad (26)$$

and $E_{max} \leq 1/10 P_{err}$.

f. Extrapolation. A second-order extrapolation function

$$v_i(t^*) = 2v_i(t^* - \Delta t) - v_i(t^* - 2\Delta t) \quad (27)$$

was tested in an effort to improve the starting values for iteration through the 19 equations, but this generally produced no improvement in convergence rate. A third-order function

$$v_i(t^*) = 3v_i(t^* - \Delta t) - 3v_i(t^* - 2\Delta t) + v_i(t^* - 3\Delta t) \quad (28)$$

was found not much better. Ultimately, of course, it is illogical to expect any finite order extrapolation scheme to predict the effect of changes in the forcing functions (area-time curves) if those coming changes had not been anticipated by the derivatives of less than that order.

g. Small Perturbation Solution. In place of an extrapolation function, there was used the more logical small perturbation solution. This considerably improved the convergence rate and provided sufficiently accurate results in lieu of the exact solution when it failed to converge in a reasonable length of time.

The complete mathematical model along with the above described convergence enhancement logic have been programmed in Fortran IV for solution on an IBM 370/165. The computer program HIRTSM1 (for HIRT Starting Model) is composed of the normally expected components: the main program (MAIN) containing the exact equations, the convergence control logic, and the overall solution control logic; subroutines to control input (INPUT), output (PRINT and DUMP), and variable definition and initialization (CONST and INIT); and a subroutine which performs the calculation for the analytical solution to the simultaneous small perturbation equations (SMPERT). In addition, the program contains a package of utility subroutines: one routine contains the logic of Fig. 10 to numerically revert any given function (SOLVER); a second expands out the binomial coefficients (BINOM) to give a series which is reverted by a third subroutine (REVERT) to the seventh-degree term; a fourth subroutine (QSIMUL) determines the points of intersection of two conics (the two final energy equations resulting from SMPERT) by converting them to a single fourth-degree polynomial, which has an exact analytical solution for the four roots (QANDC). Use of this program is described in Appendix C.

The program can be run in a partition of 110K bytes and easily completes about 200 time increments in less than a minute of central processor time, though occasionally a run may require up to three minutes. Peripheral storage is not essential, though provisions are made to dump the entire solution on to a direct (random) access data set (such as a disk file) so that the solution may be picked up at any point and continued. The results of calculations with HIRTSM1 are compared in the next section with experimental results from the Pilot HIRT facility.

3.0 RESULTS

Presented below is a comparison between the mathematical model and experimental pressure-time histories from Pilot HIRT. Included is a brief description of those characteristics of the tunnel important to the model. After a comparison of the model and data, some other results of the calculations are shown. The section concludes with a discussion of how the model can be applied in the design of certain portions of the tunnel.

3.1 DESCRIPTION OF PILOT HARDWARE

Figure 12 shows an elevation line drawing of the Pilot HIRT facility, to which the present mathematical model was applied. Figure 13 shows most of the geometric data required by the model and also accurately illustrates the real life hardware, which is simplified in the model. The geometric parameters in the precise form used in the model are summarized in Table 2. The tunnel uses two alternate types of starting devices, the sliding sleeve valve shown in Fig. 13 and, for quicker starts, a Mylar[®] diaphragm and cutter located at the interface of the diffuser and the valve assembly. The plenum exhaust system, shown schematically in Fig. 14, also uses a diaphragm in addition to two valves to control the exhaust flow. The diaphragm initiates the flow, and the ball valve, whose setting cannot be changed during a run, determines the amount of plenum exhaust during the steady portion of the run. The quick-acting valve, however, may be rapidly closed during the run to provide a temporarily elevated plenum exhaust in excess of what the ball valve will pass. The complete system in Fig. 14 is modeled as the area-time curve of a one-dimensional sonic orifice, as is the multiple port system on the main valve.

The portion of the tunnel shown in Fig. 13 was heavily instrumented with pressure taps to measure pressure-time histories at various locations in the nozzle, test section, diffuser, and plenum. Output from the pressure transducers was sampled every 2 msec by a data acquisition system based on a PDP 11/10 digital computer with certain of the signals also displayed on a recording oscillograph. Of primary interest here are the plenum pressure-time histories, which comprise the primary basis for comparison of the theory and experiment.

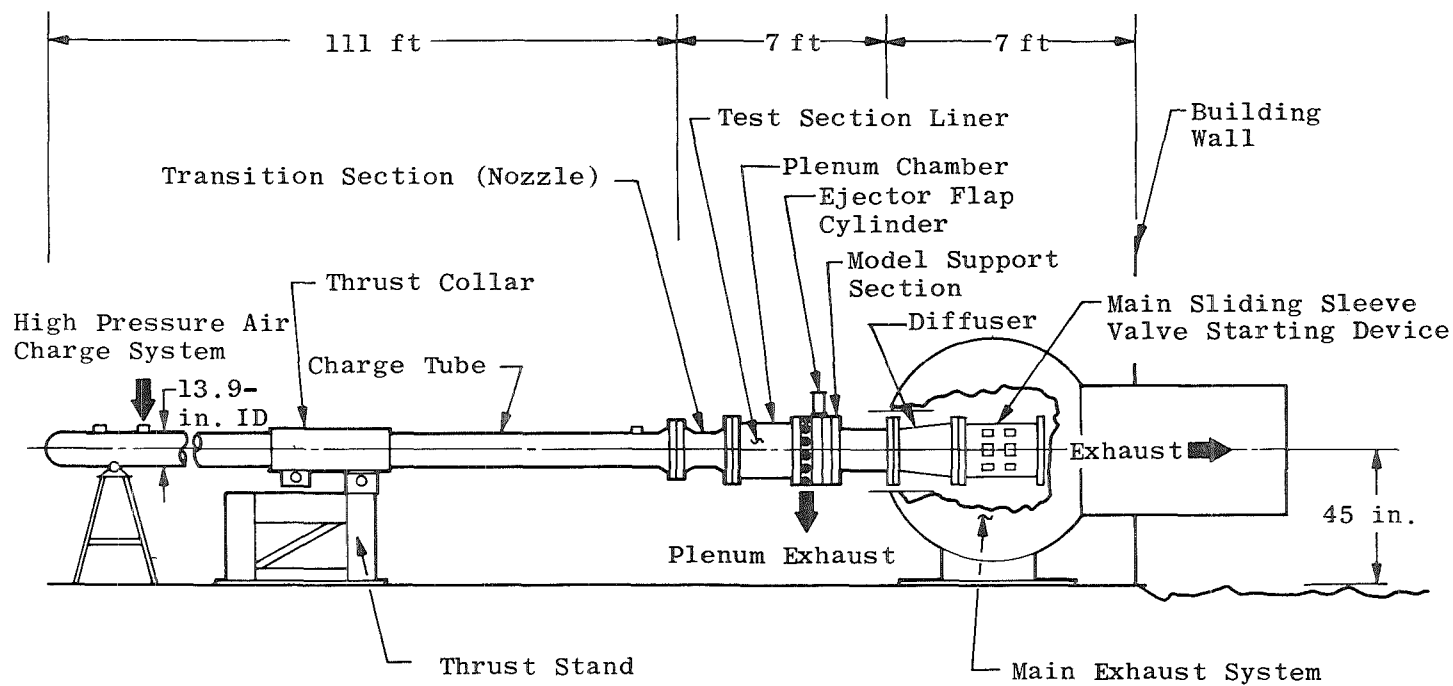


Figure 12. Pilot HIRT elevation line drawing.

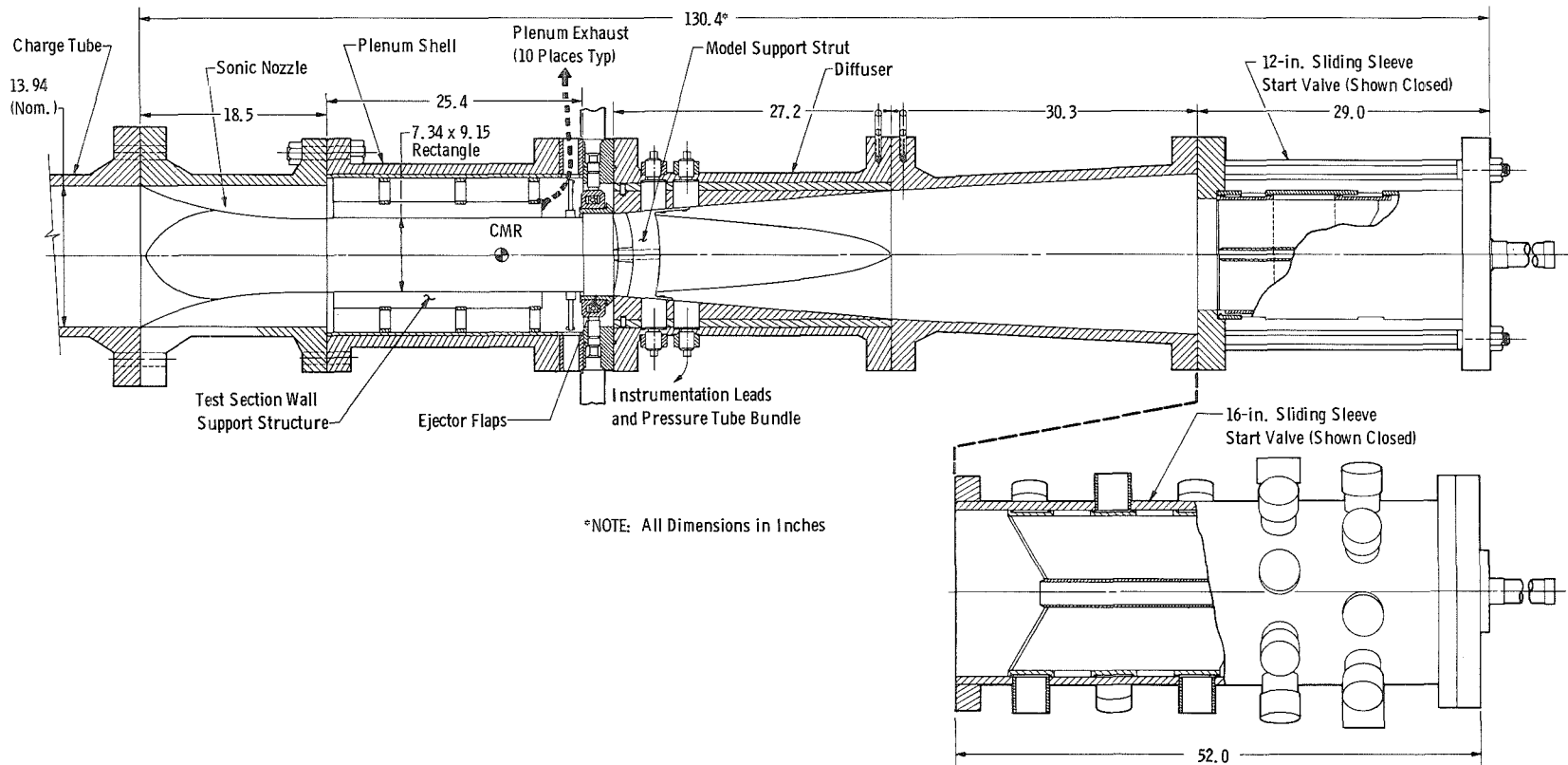
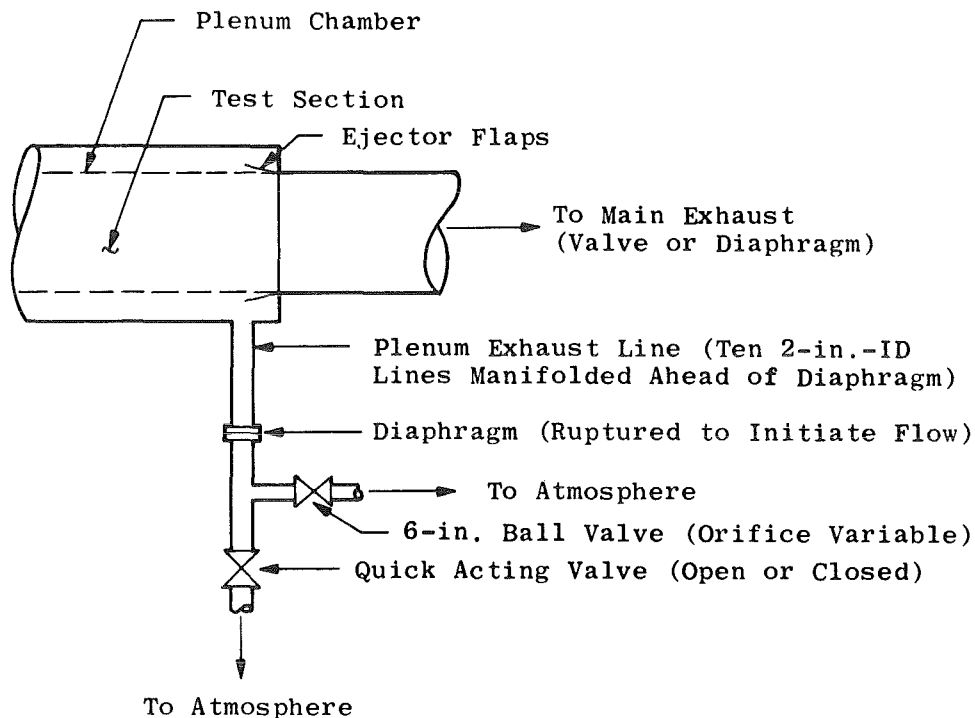


Figure 13. Cross-sectional view of nozzle, test section, diffuser, and main valve system.

Table 2. Geometric Data for Pilot HIRT Required by Mathematical Model

Charge Tube Diameter	1.162 ft
Charge Tube Flow Area	1.060 ft ²
Ratio of Charge Tube Area to Test Section Area	2.271
Test Section Length	2.114 ft
Test Section Width	0.7633 ft
Test Section Height	0.6117 ft
Test Section Flow Area	0.4669 ft ²
Test Section Wall Surface Area	5.813 ft ²
Test Section Porosity	3.5 to 10%
Test Section Volume	0.9870 ft ³
Flap Flow Area	0 to 0.2062 ft ²
Ratio of Plenum Volume to Test Section Volume	1.75 to 4.0
Nominally	2.8

**Figure 14. Plenum exhaust system.**

3.2 COMPARISON OF MATH MODEL AND EXPERIMENT

Data for nine different tunnel settings were studied with the mathematical model. Some basic data for runs typical of these nine conditions are summarized in Table 3. The data of primary interest in this table include the plenum-to-test section volume ratio, porosity, the opening times of the main valve and plenum exhaust valve, the maximum plenum exhaust area, and the experimental test section Mach number. The conditions listed for Run 2258 may be considered nominal values from which variations in plenum volume, porosity, flap setting, and test section Mach number were examined.

Figure 15 compares the experimental plenum pressure as a function of time with the present mathematical model for the nominal conditions (Run 2258). The data illustrated is for a plenum volume 2.8 times the test section volume, a porosity of 4-1/2 percent, and a flap setting of 0.4 in. (the gap between the flap and the test section wall where the flap flow empties into the diffuser). The main starting device was a Mylar diaphragm; and the exit flow area, the primary factor determining the asymptotic test section Mach number (0.921), was obtained by capping off the proper number of exit ports on the main exhaust manifold (Fig. 13, 16-in. valve). Since the desired Mach number was subsonic, the plenum exhaust system was not used. The resulting data for these tunnel settings are plotted in Fig. 15 as circles, and the solid line represents the output of the computer program. The program was run for the indicated tunnel settings (Table 3), but several not readily apparent inputs were assumed. The starting device (diaphragm) was treated as a linear area-time curve reaching its maximum area in 2 msec. The maximum area shown in Table 3 is approximately 99.46 percent of the test section flow area, which is based on the ideal, one-dimensional flow area ratio needed to produce a test section Mach number of 0.921. The resulting theoretical plenum pressure-time history shown in Fig. 15 agrees well with the experimental data. The greatest discrepancy occurs at 25 msec and reaches a peak there of 6.5 percent. This difference, due to a temporary leveling of the experimental data between 10 and 25 msec, results from the finite time required for the initial expansion wave to traverse the plenum volume, which includes the plenum exhaust lines shown in Fig. 14. These lines extend to a distance of about 4 ft from the major portion of the plenum. Since the model assumes a uniform plenum, it cannot account for this factor. Figure 15 also illustrates another deficiency of the model, which in this case produces the 3.1-percent error at a time of about 100 msec. Part of this error is due to error accumulation in the small perturbation solution, to which the program reverted entirely beyond 45 msec because of nonconvergence of the exact iterational solution. Another part of the error, in this case the smaller part, is due to neglect of the axial momentum of the test section flow by the crossflow model, which results in the smaller slope of the theoretical curve in the region of 60 to 90 msec. Since this discrepancy has been found to be generally small for subsonic runs, the coefficient in the crossflow model (A_{15}) has been left equal to one.

Table 3. Summary of Run Conditions for Experimental Data to be Compared with Theory

Run Number	Charge Pressure, P_c , psia	Plenum Volume (-), V_p/V_{ts}	Porosity, τ , %	Maximum Flow Area			Total Opening Times			Plenum Delay, sec	Asymptotic Plenum Pressure, psia	Test Section Mach Number (-), M_∞
				Main Valve, A_e , ft ²	Plenum Ex, A_{pe} , ft ²	Flaps, A_f , ft ²	Main Valve, sec	Plenum Ex, sec	Flaps, sec			
2226	60.11	2.8	4.5	0.466886	0	0.045835 ^a	0.002	---	---	---	25.00	0.992
2236	62.37	2.8	4.5	0.466331	0	0.2062 ^b	0.002	---	---	---	25.55	0.962
2241	61.84	2.8	1.5	0.465911	0	0.09167 ^c	0.002	---	---	---	23.83	1.013
2251	81.47	2.5	4.5	0.465911	0.1090	0.09167	0.002	0.040	---	0.005	30.56	1.039
2255	81.27	2.5	4.5	0.465911	0.1090	0.09167	0.002	0.040	---	0.004	24.04	1.228
2258	70.51	2.8	4.5	0.464351	0	0.09167	0.002	---	---	---	30.47	0.921 ^d
2260	70.90	4.0	4.5	0.466290	0	0.09167	0.002	---	---	---	29.36	0.960
2263	74.10	1.75	4.5	0.466662	0	0.09167	0.002	---	---	---	29.64	0.975
2742	152.15	2.5	4.0	0.465911	0.1090	0.09167	0.030	0.040	---	0.005	53.25	1.100

(a) $\dot{V}_f = 0.2 \text{ in.}^2$

(b) $\dot{V}_f = 0.9 \text{ in.}^2$

(c) $\dot{V}_f = 0.4 \text{ in.}$

(d) Nominal Conditions

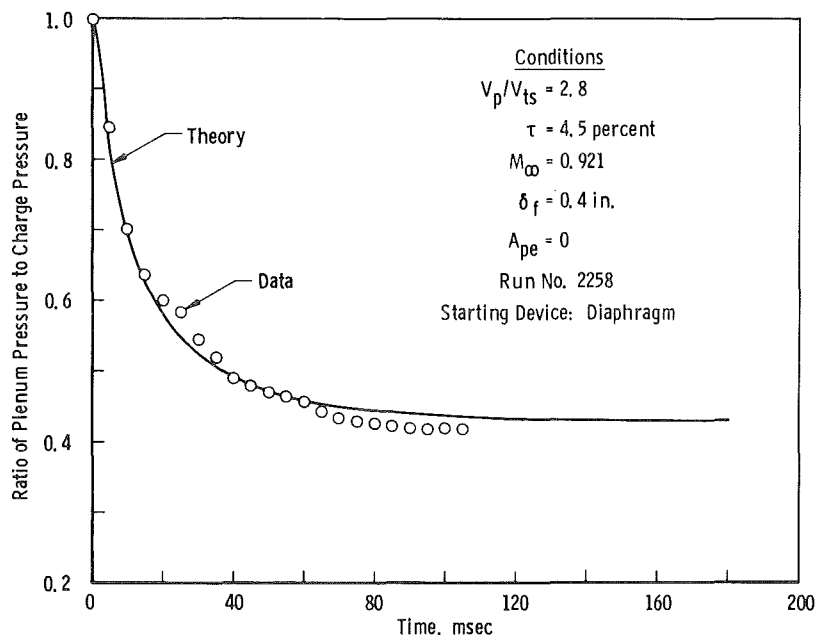


Figure 15. Plenum pressure versus time for subsonic run with medium plenum volume.

Since the amount of plenum volume which must be drawn down to the asymptotic pressure may logically be expected to have a first-order impact on the starting time, the plenum volume was the first parameter varied from the nominal conditions for Run 2258 (Fig. 15). Figures 16 and 17 show the plenum pressure for a smaller plenum volume ratio (1.75) and a larger ratio (4.0), respectively. As expected, the smaller volume case flattens more quickly than the medium volume case, and the larger volume more slowly. As in Fig. 15, the accuracy of the model is generally good for both the smaller and larger plenum volumes, though the effect of the wave propagation time in the plenum is much more pronounced for the larger volume.

Now return to a medium plenum volume case but vary another parameter - plenum exhaust - for a slightly supersonic run. The theoretical analysis depends on an experimentally derived plenum exhaust area-time curve, shown in Fig. 18, in the nondimensional form used by the computer program. Unfortunately, the uncertainty in the shape of this curve is quite large, and only the steady area is known accurately. Illustrated in Figs. 19 and 20 are the data for two supersonic cases, Mach 1.039 and 1.228. Both the theory and experiment of Fig. 18 show a slight over-shoot bottoming out at 30 msec and then approaching the asymptote from below. In addition, the experimental data show a slight rebound peaking at 60 msec, a result not predicted by the model. The rebound probably results from the overshoot, which would tend to draw the test section below its asymptotic pressure while the plenum exhaust area was decreasing

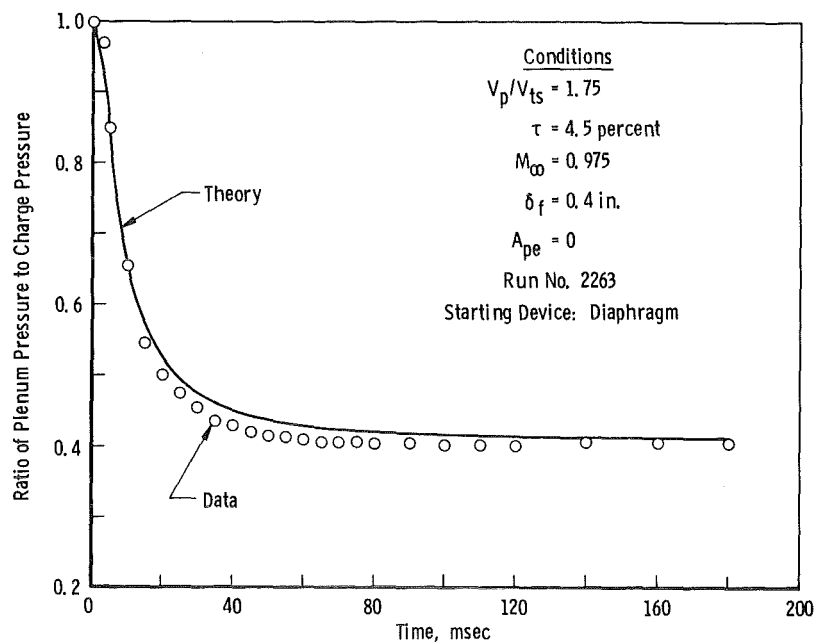


Figure 16. Plenum pressure versus time for subsonic run with small plenum volume.

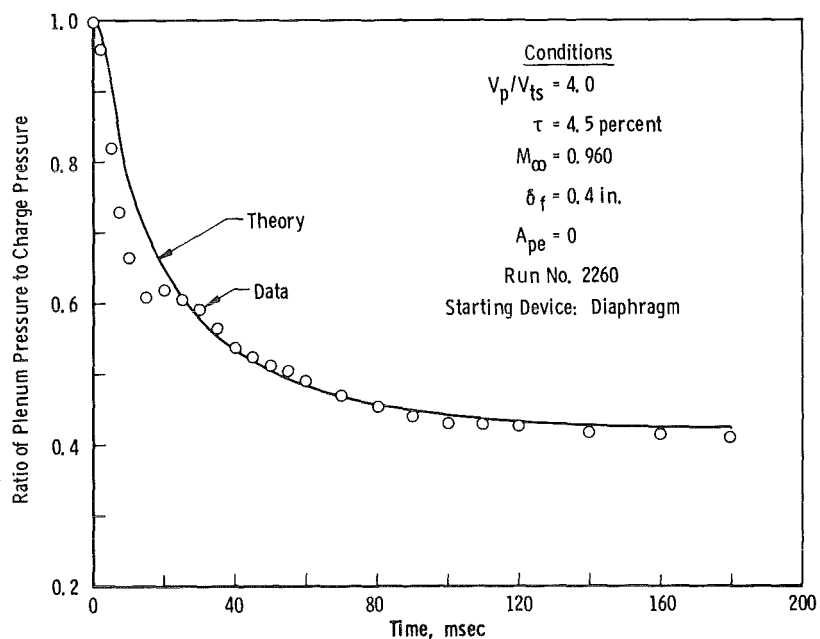


Figure 17. Plenum pressure versus time for subsonic run with large plenum volume.

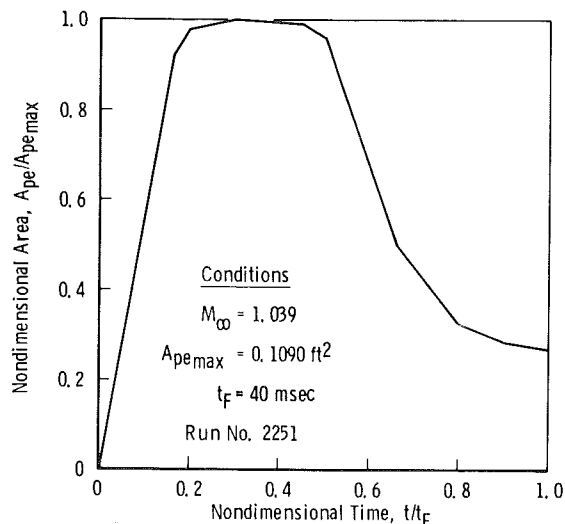


Figure 18. Plenum exhaust area-time curve for Mach 1.039 run.

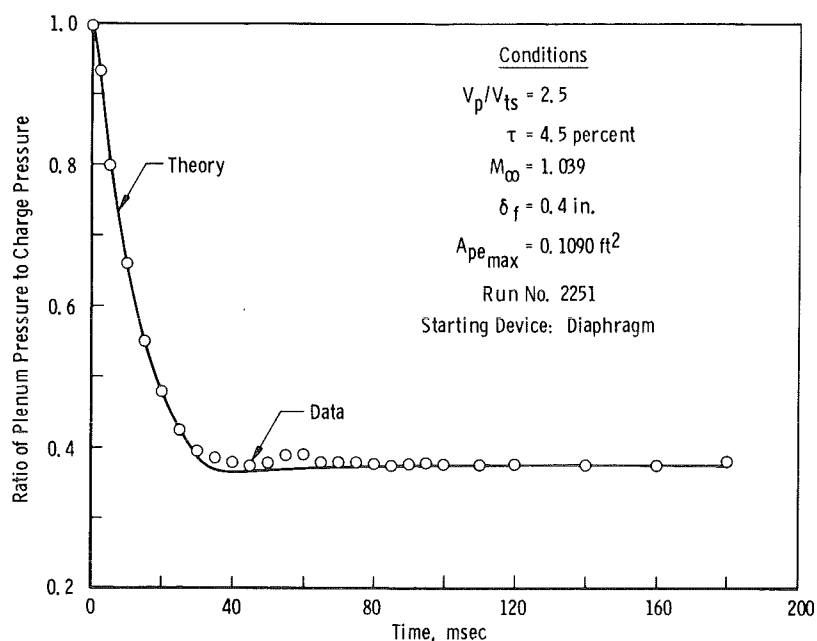


Figure 19. Plenum pressure versus time for supersonic run (Mach 1.039) with plenum exhaust.

to its steady value at 40 to 50 msec. This combination of occurrences would then produce a slight refilling of the plenum, manifesting itself in the observed rebound. For the higher Mach number (1.288) of Fig. 20, the plenum exhaust curve of the previous case was retained intact up to its peak but was linearly stretched beyond the peak to make it approach the steady area needed for the tunnel to reach the desired asymptotic Mach

number. The peak area and closing time were unchanged. The disagreement between theory and data at the knee of the curve may be charged to the uncertainty in the plenum exhaust area-time curve, which is known to vary somewhat from run to run since the plenum diaphragm rupture is not precisely repeatable.

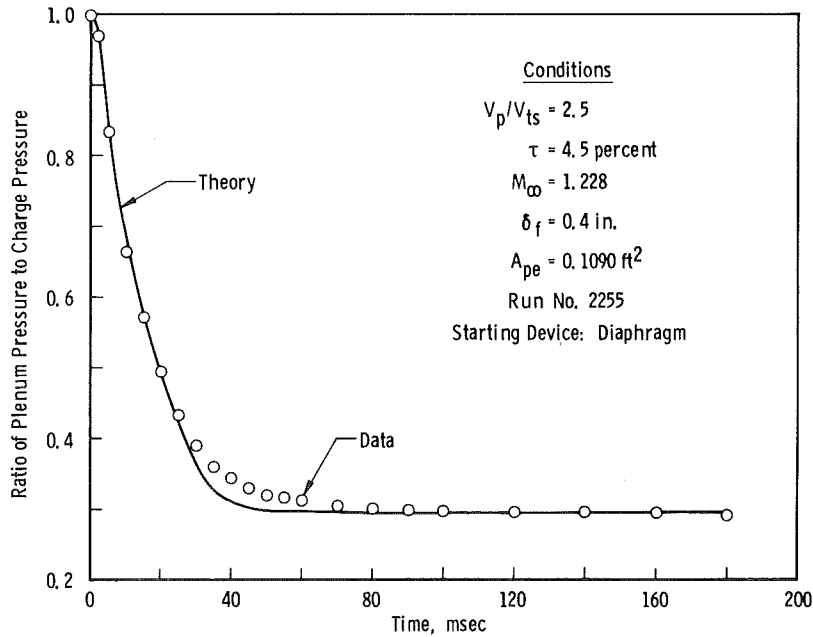


Figure 20. Plenum pressure versus time for supersonic run (Mach 1.228) with plenum exhaust.

The next parameter variation for which the model was tested was the opening time of the main starting device. Figure 21 shows the data and theory for a supersonic run made with a relatively slow opening 12-in. sliding sleeve valve instead of the diaphragm. Though not apparent from the excellent agreement for this case, there is also some uncertainty in the effective opening time of the main valve, assumed to be 30 msec for the theoretical calculation. This uncertainty results because the choke point of the tunnel changes position as the valve area increases, moving from the valve to the nozzle exit. Since the time at which this change occurs is not easily determined experimentally, the exact effective opening time is not known. In addition, the area-time curves are not precisely repeated from run to run.

To continue with the testing of the model for variations in other parameters, the program was run for a case of reduced porosity (1.5 percent), maintaining the nominal conditions of medium plenum volume and flap setting. Figure 22 shows that the model agrees well with the data. Cases were also run for which the flap flow area was halved

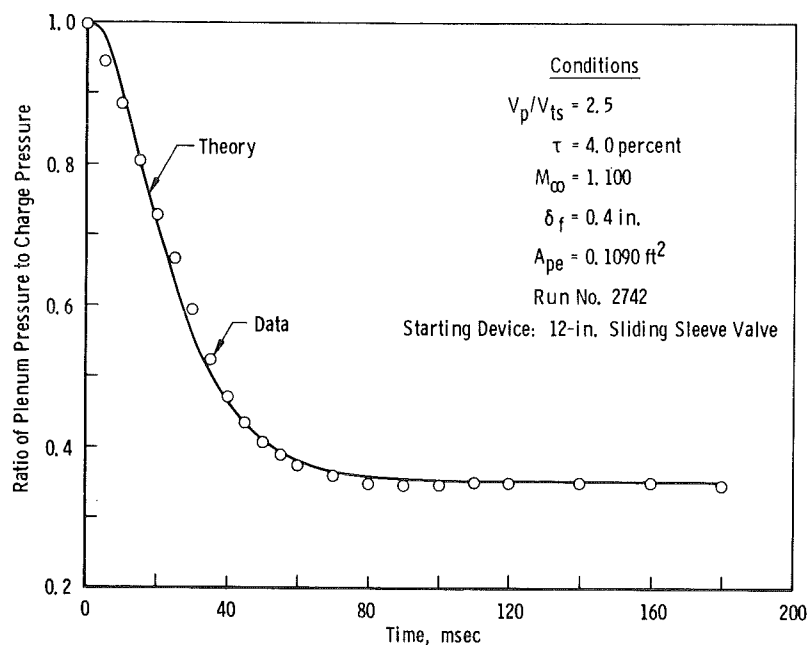


Figure 21. Plenum pressure versus time for supersonic run with sliding sleeve valve and plenum exhaust.

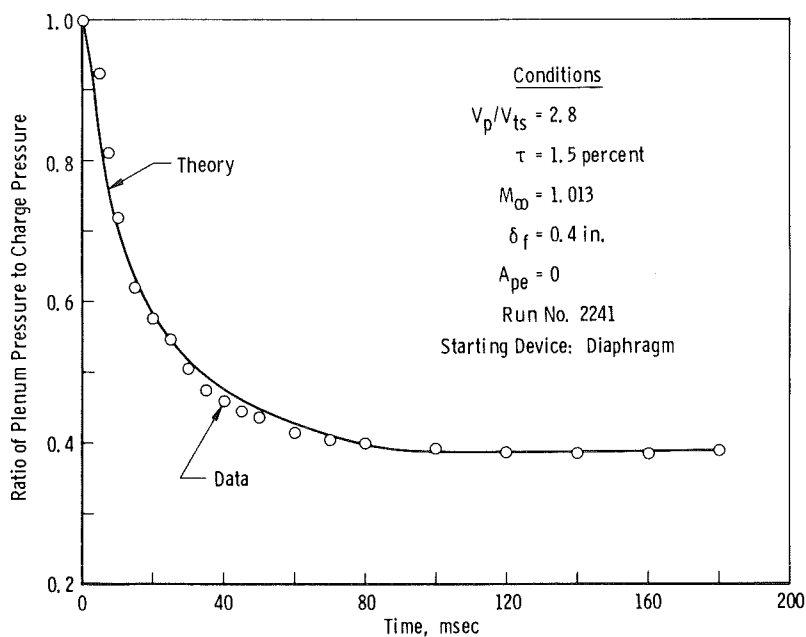


Figure 22. Plenum pressure versus time for supersonic run with 1-1/2-percent porosity and no plenum exhaust.

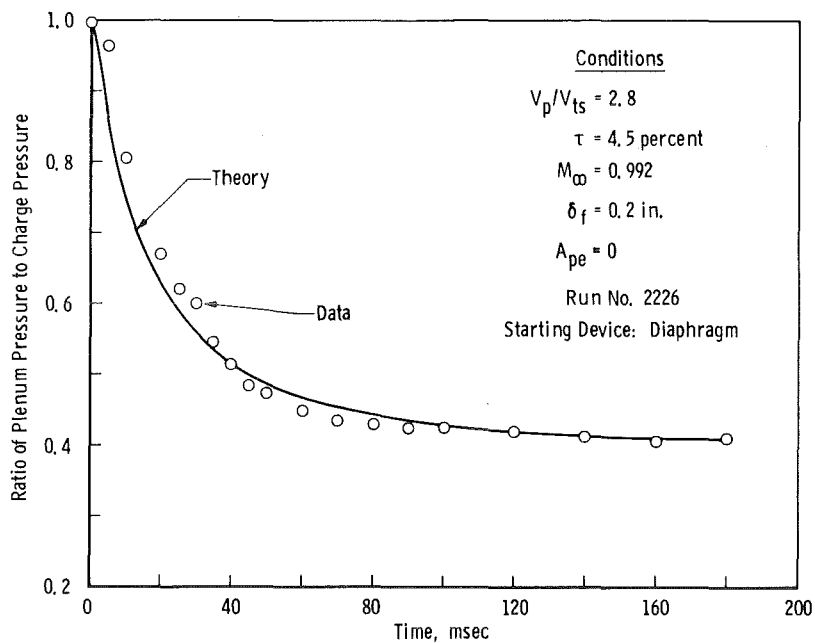


Figure 23. Plenum pressure versus time for subsonic run with small flap setting.

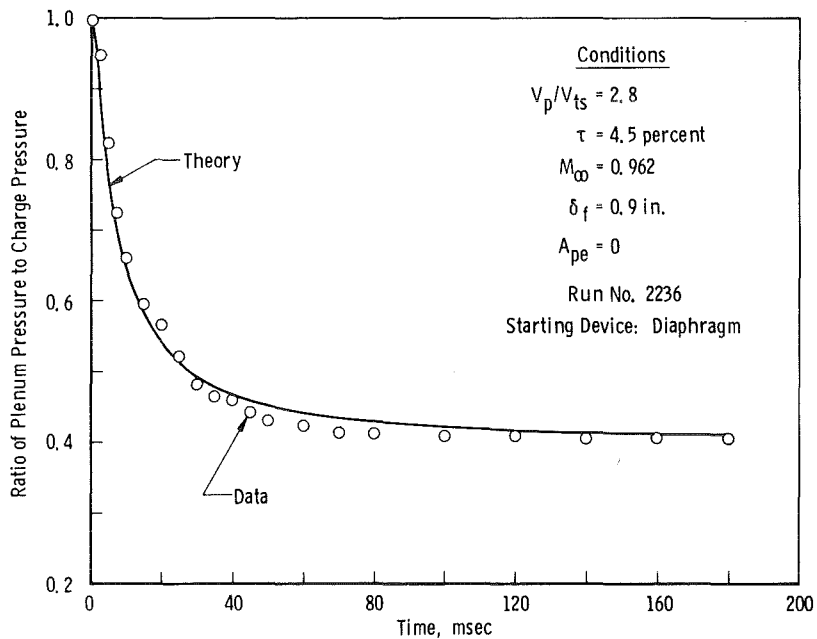


Figure 24. Plenum pressure versus time for subsonic run with large flap setting.

and doubled from the nominal settings. Illustrated in Figs. 23 and 24, both theoretical calculations are in acceptable agreement with experiment. As in previous cases, the disagreement just above the knees is due to the neglect of the finite wave propagation time across the plenum. The disagreement very early in the run (10 msec) is due to uncertainty in the rupture time of the diaphragm, and the slowness of the model in approaching the asymptote may be charged to inadequate handling of the momentum terms in the crossflow model.

3.3 OTHER RESULTS FROM THE MATH MODEL

To predict the data of primary interest, plenum pressure, the model must also calculate many other quantities including pressures and mass flow rates at various locations in the tunnel. Figure 25 shows the pressure-time histories for the case of nominal plenum volume (2.8) for a subsonic run with a diaphragm starting device. Besides plenum pressure, the stagnation pressure and static pressures at opposite ends of the test section are shown. This graph illustrates that the test section pressure initially drops much faster than the plenum, as expected since the rate of plenum depletion is limited by the porosity and flap area. Early in the run, the pressure at the exit of the test section leads the pressure at the entrance because the wall crossflow leaving the plenum increases the flow rate from the entrance to the exit. Eventually, of course, the test section and plenum pressures approach each other as the flap and wall crossflows become negligible and the steady conditions are reached. The stagnation pressure becomes nearly flat long before the static pressures in the test section and changes very slowly beyond 20 msec.

The subsonic case in Fig. 25 may be contrasted to the supersonic case in Fig. 26, which shows the same set of pressure curves. Besides the more rapid drop of all curves prior to 40 msec, due to the plenum exhaust, the most striking difference from the subsonic case is the approach of opposite ends of the test section to distinctly different asymptotes. The entrance to the test section levels rather suddenly at the choking pressure ratio, while the exit continues to drop to the lower pressure ratio corresponding to the supersonic Mach number. Another interesting feature is that the asymptotic pressure at the test section exit is lower than for the plenum even though the net wall crossflow must be into the plenum (to reduce the flow rate along the test section as needed for supercritical flow). Crossflow against the pressure gradient occurs because of the increasing momentum retained by the crossflow while separating off from the high-speed test section flow. Another feature of Fig. 26 due to this momentum is the crossing of the test section pressure curves at 12 msec, which signifies the reversing of the wall crossflow. To improve the crossflow model's representation of the effect of this momentum (which is neglected in modeling the crossflow rate as a function of pressure difference only), the momentum correction coefficient A_{15} in Eq. (7) was introduced. This quantity expediently models the small

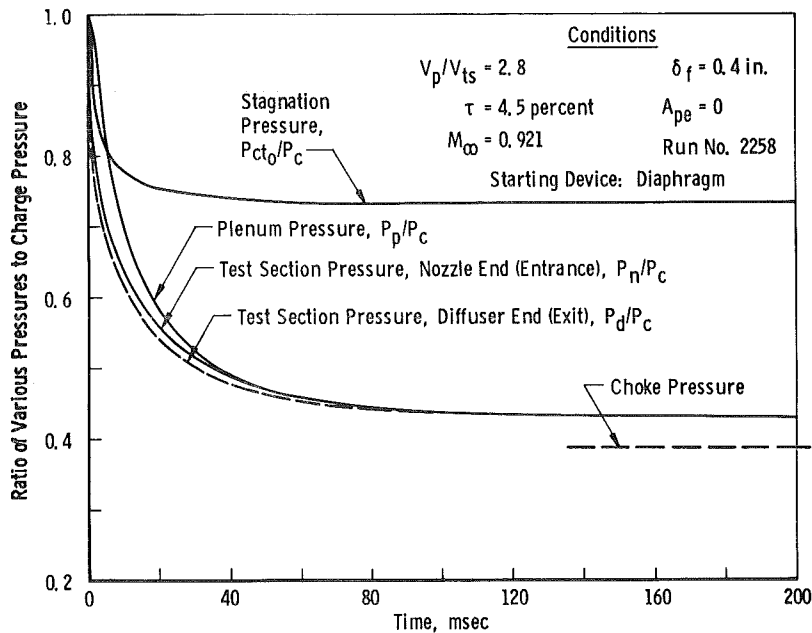


Figure 25. Various pressures versus time for nominal conditions.

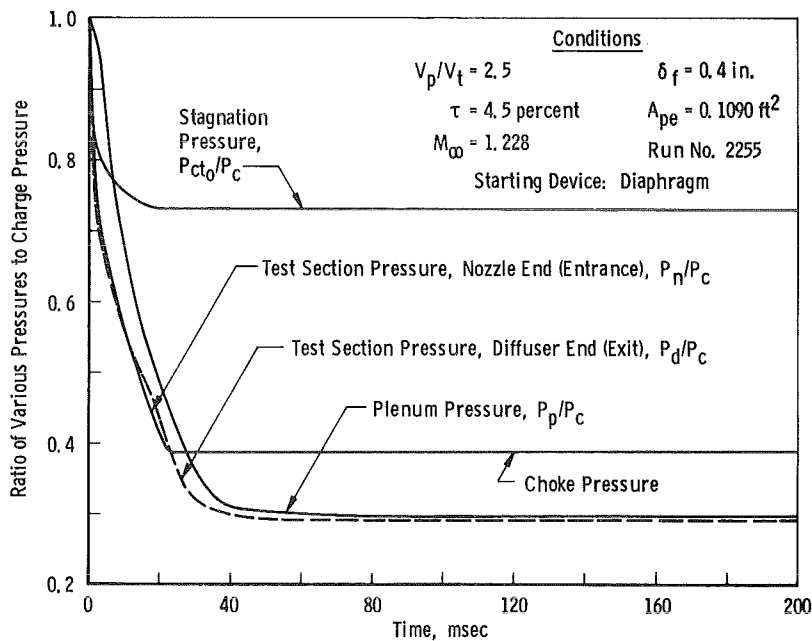
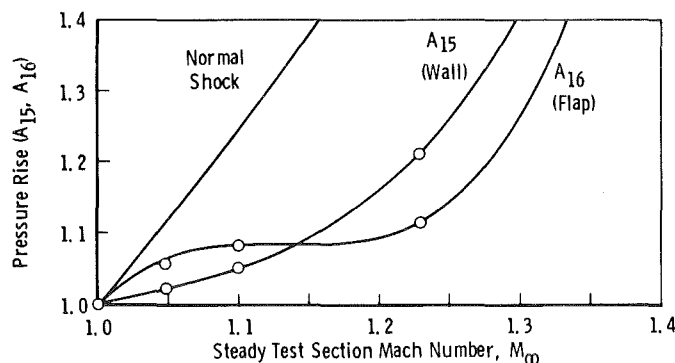
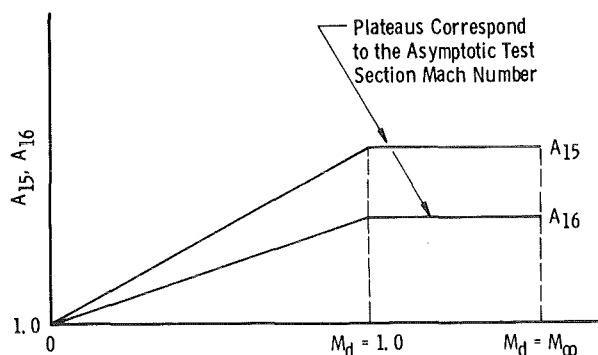


Figure 26. Various pressures versus time for supersonic run with plenum exhaust.

additional crossflow due to momentum in terms of a slightly elevated driving pressure. The steady-state value of A_{15} at a given steady test section Mach number was derived empirically for a given steady plenum pressure. These steady-state values of A_{15} are shown in Fig. 27a. During a run, however, A_{15} was assumed to vary according to the ramp function of Fig. 27b to simulate the increasing momentum.



a. Momentum correction coefficient (A_{15}) and flap correction coefficient (A_{16}) versus steady test section Mach number



b. Assumed variation with test section Mach number (M_d) of momentum (A_{15}) and flap (A_{16}) correction coefficients during starting process

Figure 27. Steady-state values of correction coefficients, A_{15} and A_{16} .

Looking at the mass flow rate-time curves corresponding to Figs. 25 and 26 provides further insight into the behavior of the mathematical model. Figure 28 shows the flow rate entering (from the charge tube) and leaving the test section, the flow rate through the flaps, and across the porous wall for the nominal conditions and subsonic flow. The flap and wall crossflows, though leaving the plenum in this run, are shown on the positive

axis for convenience. All data are expressed as ratios of the steady, asymptotic flow rate through the main valves. The flow in the test section is seen to rise very rapidly, in concert with the breaking diaphragm, and to approach the final flow rate only as the flap and crossflows approach zero. Both flows from the plenum reach peaks at about 3 msec, which results from the pressure differences between the plenum and test section reaching a maximum. The crossflow further manifests itself in the disparity between the flow entering and leaving the test section. Various experimentally derived flow rates are given in Ref. 4 for the pilot tunnel. These relatively well behaved results for the subsonic case may be contrasted to the tangle of curves resulting from a supersonic case with plenum exhaust (Fig. 29), which is based on the same conditions as Fig. 26. Initially similar to the subsonic case with peak flap and crossflows at 3 msec, the curves are considerably modified by the opening of the plenum exhaust at 4 msec (a programmed delay). The leveling of the flap and crossflow curves at 22 msec is associated with choking in the test section. Eventually, the plenum exhaust forces both the crossflow and flap flow to reverse and eventually to exactly balance the plenum exhaust flow rate when steady flow is reached. Reversal of the flap flow requires, in terms of the flow model (Eq. (8)), a driving pressure at the flap exit greater than the plenum pressure and in general greater than the computed pressure at the exit of the test section. Though the flap correction coefficient (A_{16}) is applied much like the wall crossflow coefficient, the physical explanation cannot be the same since the free-stream momentum is in the opposite direction of the reversed

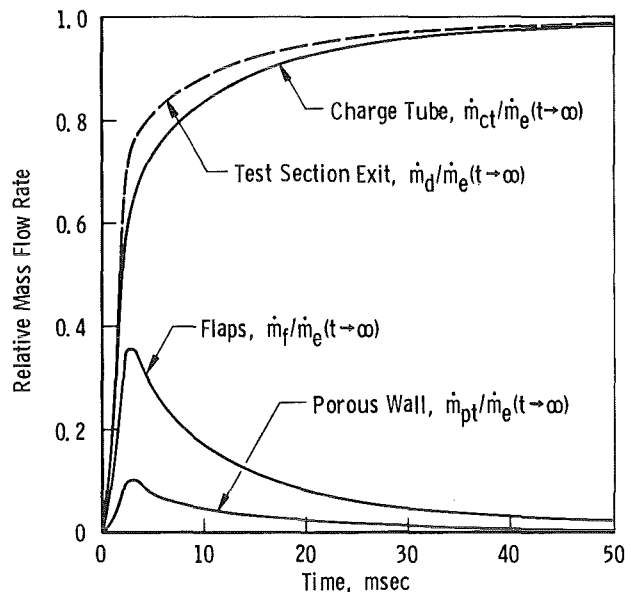


Figure 28. Relative theoretical mass flow rates for nominal conditions (Run 2258) of subsonic flow with no plenum exhaust.

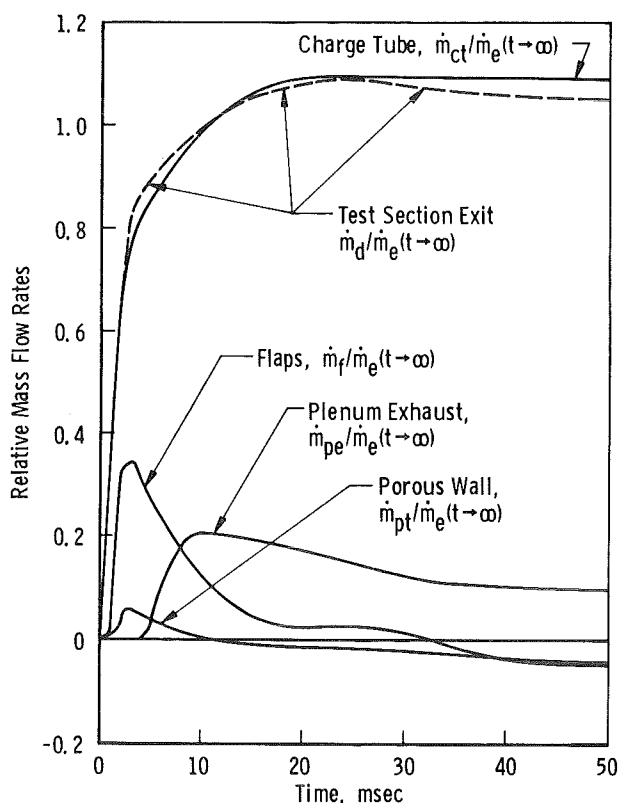


Figure 29. Relative mass flow rates for a supersonic run (Mach 1.228) with plenum exhaust (Run 2255).

flap flow. A more likely explanation is the shock structure and flow separation at the diffuser entrance. Since precise modeling of this complex flow is beyond the scope of the present work, the flap flow correction coefficient (A_{16}) was added to Eq. (8). Experimentally derived values of A_{16} as a function of steady test section Mach number are plotted in Fig. 27a along with the static pressure jump across a normal shock. The pressure rise during the reversed flap flow must be due to a flow more complex than a normal shock, since the pressure jump across the shock rises much more rapidly than experiment indicates. The lines through the circled points are cubic fits and are probably not accurate beyond Mach 1.25. As with the momentum correction, the flap correction was assumed to vary in time according to the ramp function in Fig. 27b.

3.4 APPLICATION OF THE MATH MODEL

Besides prediction of tunnel start time, there are several other ways the model can be applied in the design of a wind tunnel. Since the plenum exhaust area-time curves can be varied arbitrarily in the model, the number of plenum valves (or total valve area)

to achieve various start times can be determined. In addition, the sensitivity of the start time to the shape of the area-time curves can be predicted. This is important because it indicates how finely controllable and repeatable (and expensive) the valves must be. Another potential application is estimation of the structural loading of the test section wall due to transient pressure differences between the plenum and test section.

To illustrate some of these possibilities, the program was run for the three different plenum exhaust area time curves shown in Fig. 30. The solid line is a typical area-time curve from Pilot HIRT, and the two broken lines are variations having the same average open area. Processing the model with the triangular curve should indicate whether a curve with the same peak as the basic curve but having a different shape would significantly affect starting time. The trapezoidal curve should indicate whether a smaller number of valves kept at full open for a longer time could achieve the same start time as the more peaked curves. The plenum pressure-time histories for these three curves are shown in Fig. 31. It is clear that the triangular curve has little effect on the shape of the pressure curve and does not affect starting time. On the other hand, the trapezoidal curve has a larger effect but still does not lengthen the starting time. The logical conclusions for the tunnel configuration studied here is that very accurate controllability is not required of the plenum valves and that the tunnel could be started just as quickly with about

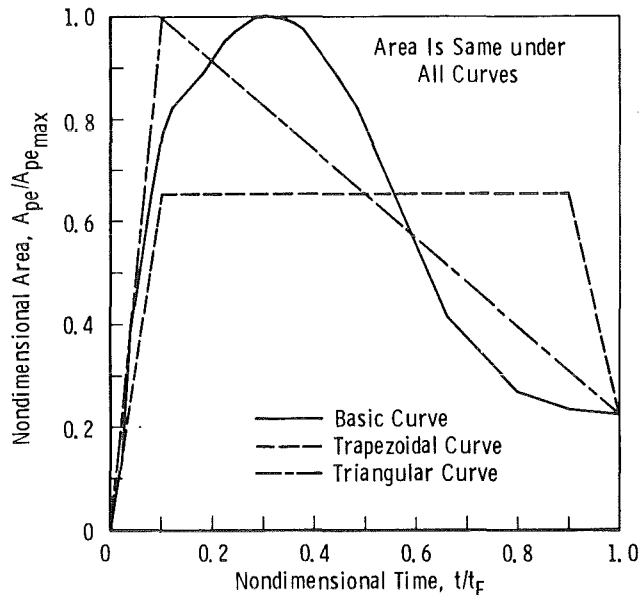


Figure 30. Nondimensional equal area plenum exhaust area-time curves.

2/3 of the available valve area if the valves were kept fully open for a longer duration. If these results were found to apply to a large scale facility, a considerable cost reduction could be realized.

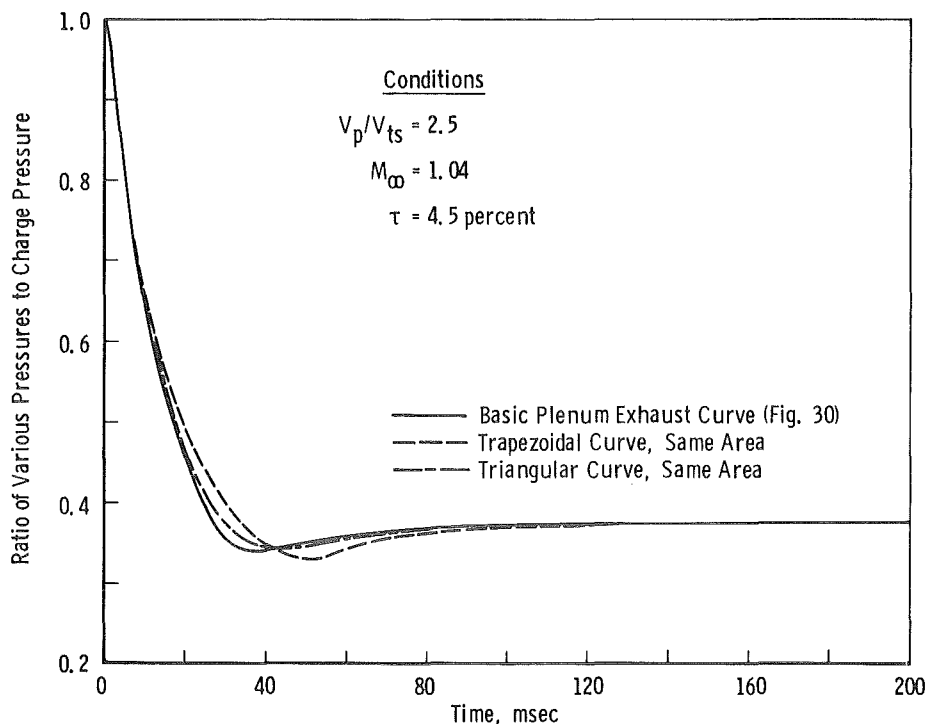


Figure 31. Plenum pressures versus time for three plenum exhaust area-time curves with same integrated area.

A second example of application of the model is illustrated by Fig. 32, which shows the pressure differential across the wall at the test section exit as a function of time for several conditions. From these results, it can be seen that reducing the porosity has little impact on wall loading, but raising the Mach number from 0.921 to 1.228 or reducing the flap gap by 1/2 significantly increases the loading by 25 and 50 percent, respectively. In contrast, lengthening the effective valve opening time from 2 to 30 msec reduces the peak load to about 1/3 of the nominal case. The peaks of the curves for the diaphragm runs occur just as the diaphragm reaches its full open area. The curve for the valve run, however, peaks first when the plenum exhaust area peaks and later when the valve reaches its steady area around 30 msec. Two data points for the peak pressure differential from Ref. 4 are shown in Fig. 32 and agree with the model.

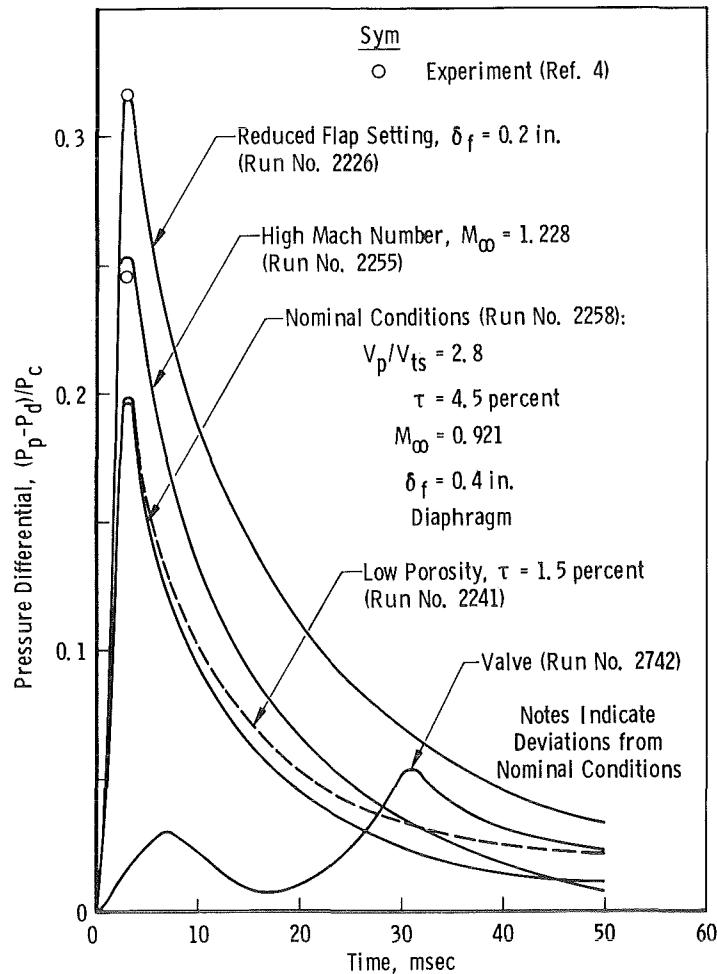


Figure 32. Transient loading of test section wall at exit for nominal conditions and selected deviations.

4.0 SUMMARY AND CONCLUSIONS

A mathematical flow model for the process of starting a transonic Ludwig tube wind tunnel has been developed. The present model uses the integral continuity equation for three specific control volumes, the steady form for the diffuser and test section control volumes, and the unsteady form for the plenum. The solution in the two former control volumes also uses the steady, isentropic energy equation, assumed applicable throughout the diffuser and test section control volumes for a given set of stagnation conditions. However, the stagnation conditions are allowed to vary in time according to the well-known exact solution for an unsteady, one-dimensional expansion wave. Application of this model takes the form of a numerical solution of 19 simultaneous algebraic equations to be solved at successive time points until the flow becomes steady. The iterative solution procedure

for these exact equations becomes nonconvergent in the vicinity of choking and is replaced with an analytical solution to a set of small perturbation equations until the choke point is passed. The numerical procedure is programmed for computer solution.

The mathematical model was evaluated by comparison with experimental plenum pressure-time histories from a small Ludwig tube wind tunnel. Agreement between the model and experiment was found to be good. Other numerical results from the computer model were also presented to illustrate application of the model to design of a large facility. Specific conclusions drawn from the present study include (1) verification of the model's ability to predict accurately plenum pressure-time histories and, therefore, tunnel starting time; (2) prediction that starting time is insensitive to the precise shape of area-time curve of the plenum exhaust and, therefore, that very precise controllability is not required of the plenum valves; (3) prediction that starting time is not significantly lengthened by even large changes in the shape of the plenum exhaust area-time curve if the area under the curve and open time are maintained, thus permitting considerable reduction in the number of start valves suggested by data from the pilot facility; and (4) verification that aerodynamic loading of the test section walls (and, therefore, the support structure) can be reduced by lengthening the opening time of the main starting valves, within limitations of the required starting time.

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APPENDIX A

SMALL PERTURBATION SOLUTION

This section presents the essential details of the small perturbation solution, the knowledge of which may be important to a user of the computer program HIRTSM1. Table A-1 shows the small perturbation variables and the exact variables they represent. Use of the expansions (Eq. (22)) in the exact equations listed in Table 1 produces the approximate small perturbation equations listed in Table A-2. Definitions of the coefficients A_1 , B_1 , C_1 ..., if needed, should be extracted directly from the computer program (subroutine SMPERT) where they are coded as CA1, CB1, CC1, ..., respectively. The equations of Table A-2 can be solved analytically without recourse to numerical iterative procedures. To accomplish this task, the linear equations were solved algebraically to eliminate all variables except those contained in the quadratic equations, Eqs. (12) and (13). After eliminating all variables but ϵ_{12} and ϵ_{13} from the two quadratics, Eqs. (12) and (13) were converted to a single quartic (subroutine QSIMUL), which was solved analytically for its four roots. If necessary, the reader can extract the algebraic details of this procedure from the computer program. The correctness of the algebra has been inferred from computation of residuals from the equations of Table A-2 (replacing the zeros on the right-hand side with residuals). For all cases tested, the residuals were found to be on the order of the computer's accuracy ($\sim 10^{-16}$). Similarly, the accuracy of the expansions in representing the exact equations was tested by computing residuals from the exact equations using perturbed values for the variables. The largest residuals (percentage basis) were generally less than 10^{-4} .

Table A-1. Perturbation Variables

Original Variable	Perturbation Variable
$\dot{m}_e (t^*)$	ϵ_1
$\dot{m}_{pe} (t^*)$	ϵ_2
$\dot{m}_f (t^*)$	ϵ_3
$\dot{m}_{pt} (t^*)$	ϵ_4
$\rho_p (t)$	ϵ_5
$\dot{m}_d (t^*)$	ϵ_6
$\dot{m}_{ct} (t^*)$	ϵ_7
$M_{ct} (t^*)$	ϵ_8
$P_{e0} (t^*)$	ϵ_9
$T_{e0} (t^*)$	ϵ_{10}
$P_t (t^*)$	ϵ_{11}
$P_d (t^*)$	ϵ_{12}
$P_n (t^*)$	ϵ_{13}
$\dot{m}_o (t^*)$	ϵ_{14}
$P_p (t^*)$	ϵ_{17}^a
$\rho_p (t^*)$	ϵ_{18}
$T_p (t^*)$	ϵ_{19}
$A_e (t^*)$	ϵ_{A_e}
$A_{pe} (t^*)$	$\epsilon_{A_{pe}}$
$A_f (t^*)$	ϵ_{A_f}

(a) Variables 15 and 16 were eliminated.

Table A-2. Perturbation Equations

Program Equation Number ^a	Perturbation Equation ^b
1	$A_1 \epsilon_1 + B_1 \epsilon_9 + C_1 \epsilon_{Ae} + D_1 \epsilon_{10} = 0$
2	$A_2 \epsilon_2 + B_2 \epsilon_{17} + C_2 \epsilon_{Ape} + D_2 \epsilon_{19} = 0$
3	$A_3 \epsilon_3 + B_3 \epsilon_{Af} + C_3 \epsilon_{17} + D_3 \epsilon_{12} = 0$
4	$A_4 \epsilon_4 + B_4 \epsilon_{17} + C_4 \epsilon_{11} = 0$
5	$A_5 \epsilon_5 + B_5 \epsilon_2 + C_5 \epsilon_3 + D_5 \epsilon_4 + E_5 = 0$
6	$A_6 \epsilon_6 + B_6 \epsilon_4 + C_6 \epsilon_7 = 0$
7	$A_7 \epsilon_6 + B_7 \epsilon_3 + C_7 \epsilon_1 = 0$
8	$A_8 \epsilon_7 + B_8 \epsilon_8 = 0$
9	$A_9 \epsilon_9 + B_9 \epsilon_8 = 0$
10	$A_{10} \epsilon_{10} + B_{10} \epsilon_8 = 0$
11	$A_{11} \epsilon_{11} + B_{11} \epsilon_{13} + C_{11} \epsilon_{12} = 0$
12	$A_{12} \epsilon_6 + B_{12} \epsilon_{14} + C_{12} (P_{ct0} \epsilon_{12} - P_d \epsilon_9) + D_{12} (P_{ct0} \epsilon_{12} - P_d \epsilon_9)^2 = 0$
13	$A_{13} \epsilon_7 + B_{13} \epsilon_{14} + C_{13} (P_{ct0} \epsilon_{13} - P_n \epsilon_9) + D_{13} (P_{ct0} \epsilon_{13} - P_n \epsilon_9)^2 = 0^c$
14	$A_{14} \epsilon_{14} + B_{14} \epsilon_9 + C_{14} \epsilon_{10} = 0$
17 ^d	$A_{17} \epsilon_{17} + B_{17} \epsilon_{18} = 0$
18	$A_{18} \epsilon_{18} + B_{18} \epsilon_5 + C_{18} = 0$
19	$A_{19} \epsilon_{19} + B_{19} \epsilon_{18} + C_{19} \epsilon_{17} = 0$

(a) See Table 1 for Corresponding Exact Equations

(b) Refer to Listing of Computer Program, Subroutine SMPERT, for Definitions of A_i , B_i , ...

(c) Variables P_{ct0} , P_d , and P_n Are Evaluated at $t^* - \Delta t$ As Are All the Coefficients A_i , B_i , ...

(d) Equations 15 and 16 Were Eliminated

APPENDIX B APPROXIMATED EQUATIONS

Reversion of Eqs. (11), (13), and (17) requires a time-consuming numerical procedure which has a major impact on the run time of HIRTSM1. To reduce the number of iterations needed for the reversions, approximations to the original equations were used to provide accurate initial guesses to the numerical procedure. Since these approximations may be of general interest, they are listed below. A good approximation to the mass flux-Mach number wave equation was obtained by expanding

$$\hat{m} = M \left(1 + \frac{\gamma-1}{2} M \right)^{-\frac{\gamma+1}{\gamma-1}} \quad (\text{B-1})$$

in a series of powers of M using the binominal expansion. Reversion of this series for $\gamma = 1.4$ then produced

$$\begin{aligned} M = \hat{m} - 1.200 \hat{m}^2 + 2.0400 \hat{m}^3 + 4.0480 \hat{m}^4 + 8.7965 \hat{m}^5 \\ + 20.106 \hat{m}^6 + 47.960 \hat{m}^7 + \dots \end{aligned} \quad (\text{B-2})$$

where $\hat{m} \equiv \dot{m}/\dot{m}_c$. The approximation used for the energy equation is much simpler and was discovered quite by accident. It was found that the equation

$$\tilde{m}^2 = \tilde{P}^{2/\gamma} - \tilde{P}^{\frac{\gamma+1}{\gamma}} \quad (\text{B-3})$$

could be very reasonably approximated over the interval $0 \leq M \leq 1.4$ by the ellipse

$$\left(\frac{\tilde{m}}{\tilde{m}^*} \right)^2 + \left(\frac{\tilde{P} - \tilde{P}^*}{1 - \tilde{P}^*} \right)^2 = 1 \quad (\text{B-4})$$

where

$$\tilde{m} \equiv \sqrt{\frac{\gamma-1}{2}} \frac{\dot{m}}{\dot{m}_o}$$

$$\tilde{P} \equiv \frac{P}{P_o}$$

and

$$\tilde{P} = \tilde{P}^*, \tilde{m} = \tilde{m}^* \text{ for } M = 1$$

APPENDIX C

DESCRIPTION OF THE COMPUTER PROGRAM HIRTSM1

Because of the complexity of the numerical calculations, potential users of the model must have access to the computer program (a manual calculation on a scientific calculator took about six hours to step through five time increments). For this reason, a listing of the source deck is given in this section along with a brief description of its content and use. Table C-1 lists the 15 subroutines comprising HIRTSM1. Of primary interest are the routines MAIN and SMPERT, which house the exact model equations and the small perturbation equations, respectively. Table C-2 defines some of the more important variables used in the program, information which is potentially useful if a program modification is necessary.

Of primary interest to the potential user, however, is the input, instructions for which are listed in Table C-3. The first card (NCTL) allows the user to retain manual control over some of the superficial program logic. While intended primarily for debugging purposes, the NCTL variable may be used to restart a run previously written onto a data file. To make a normal run and relinquish all control to the program, a blank card may be used. The second card (INSTR) provides the means to invoke certain program options via integer instructions. Table C-4 gives a set of values which have been used successfully to date, though occasional adjustments are necessary for some cases. Of particular importance for supersonic cases is INSTR(26). As the program approaches the choke point in the calculation (timewise, speaking), the number of iterations (ITER) for convergence always becomes inordinately large (~ 100); and the program must switch to the small perturbation solution entirely by automatically setting $\text{INSTR}(23) = 2$ when $\text{ITER} \geq \text{INSTR}(22)$. However, for supersonic cases, the solution is often not close to its asymptote, and significant error can accumulate from the small perturbation solution. To reduce this error, INSTR(26) may be used to direct the program to attempt to revert back to the exact solution a certain number of time increments (the input value of INSTR(26)) beyond the choke point. Sometimes the attempted reversion will be unsuccessful because the solution is either still too close to the choke point or is already too close to its asymptote; in which case $\text{ITER} \geq \text{INSTR}(22)$ will occur, and the program will continue with the small perturbation solution. When this situation occurs, the exact solution is not given a chance to correct the accumulated error, which may affect the asymptote by as much as 10 percent. If this result is encountered, different values of INSTR(26) should be tried, since even a temporary successful reversion to the exact solution can improve the accuracy of the solution considerably.

The remaining data cards constitute primarily a description of the tunnel and its geometry. While most of the table entries are self explanatory, some of them deserve more emphasis. On card number 4, the values of A15 and A16, if used, should be entered

as negative to invoke the use of ramps. On card number 5, the weight used in computation of the test section pressure for subsonic flow is programmed as 0.5. The input value is used only in supersonic flow. On card number 6, the variable A14 is used to sort the roots from the quartic. A value of -0.2 has been found more effective than -0.1. If the root sorting logic finds more than one value of ϵ_{13} acceptable, the program will halt in bewilderment, requiring some trial and error adjustment of A14 by the user. On card number 7, it has been found best to keep $\$EMAX \leq PERR/10$. The quantity A10 is used to obtain debugging information when $T > A10$. Following card number 10, three separate decks for the nondimensional area-time curves for the main valves, plenum exhaust valves, and flaps must be provided. Each deck must contain the number of cards entered on card number one. The times and areas must be nondimensionalized by the values entered on card numbers 9 and 8, respectively, and, therefore, will vary only between zero and one. The times must proceed in ascending order. Table C-5 gives recommended values for some of these entries. The remaining input instructions (I1, I2, ...) may be ignored unless NCTL has been entered as other than zero, in which case the user is invited to decipher the program logic in order to determine the endless uses to which this option may be put.

Table C-6 presents a sample job stream and data deck. The first four cards are peculiar to the computer facility. The first "GO" card designates data set 03 a dummy in order to suppress debugging printouts sent to DSRN* IDEBUG. The remaining data cards may be understood via Table A-3.

A portion of the output from this run is shown in Table A-7. The first four pages show the input data along with the initial values of most program variables. In addition, an interpretation of the INSTR(I) options selected is printed. The flow area-time curves are the redimensionalized form in units of seconds and square feet (or whatever units are used in the input data). The form of the remaining output is that due to the selection of $INSTR(5) = 2$ and generally displays all computed properties at the midpoint or end of each time interval. Each five lines of data separated by a space corresponds to a single time interval, and each block of five numbers corresponds to the similarly positioned block of five variable names in the page heading. Interpretation of these names may be accomplished via Table C-2. The illustrated run went to 180 msec, generated about 1,700 records (lines of print), and required 42.6 sec of central processor (CPU) time on an IBM 370/165. This run may be used as a check case by potential users.

Table C-8 presents a machine listing of the final source deck. All necessary subprograms are included except those available from the IBM subroutine library, from which HIRTSM1 uses DABS, DSQRT, DSIN, DCOS, DATAN2, CDSQRT, and CDABS.

*Data Set Reference Number

Table C-1. Description of Subroutines

PRIMARY MODEL SUBROUTINES

<u>Subroutine Name</u>	<u>Function</u>
MAIN	1. Overall program control 2. Exact model equations 3. Convergence control
SMPERT	Small perturbation equations

SPECIALIZED UTILITY SUBROUTINES

INPUT	Obtains initial data from DSRN IIN
CONST	Defines certain program constants
INIT	Initializes certain program variables
DUMP	Prints out all program variables at beginning and end of run and as needed for debugging
PRINT	Prints numerical solution and controls paging

GENERAL UTILITY SUBROUTINES

SOLVER	Provides logic for numerical reversion of a function (see Fig. 10)
BINOM	Expands a binomial to seven terms
REVERT	Reverts a series to seven terms
QSIMUL	Converts two conics to a quartic
QANDC	Computes the exact roots of a quartic
CUBRT	Computes the exact roots of a cubic
DREAL	Returns the real part of a double precision complex number
DIMAG	Returns the imaginary part of a double precision complex number

Table C-2. Definition of Major Program Variables

REAL ARRAYS

<u>Variable Name</u>	<u>Definition</u>
AREA	Input nondimensional area-time curves for main valves, flaps, and plenum exhaust valves
AREATS	Interpolated areas for time t^*
AREAM	Peak of area-time curves (dimensional)
TV	Nondimensional times for area-time curves
E	Convergence criteria errors
TVF	Total time for main valves, flaps, and plenum exhaust valves (dimensional)
TDELAY	Delays times for first motion of valves and flaps
RW	Coefficients for the reverted expansion of the mass Flux-Mach number wave equation
V	Array equivalenced to major property values
RSTR	Array equivalenced to certain real commoned variables to simplify writing of solution onto a storage device for restarting a run
ISTR	Array equivalenced to certain integer variables for storage and restarting

REAL SCALARS

Pxi	Pressure
MDxi	Mass flow rate
Txi	Temperature
Rxi	Density
Mxi	Mach number
Axi	Flow areas

Table C-2. Continued.

x-codes:

x = N	Nozzle exit (test section entrance)
= P	Plenum
= PT	Plenum at time t (PPT) or wall crossflow (MDPT)
= D	Diffuser entrance (test section exit)
= T	Test section midpoint (PT)
= CTO	Stagnation condition, charge tube
= CT	Charge tube
= E	Main valve exit
= F	Flaps
= PE	Plenum exhaust
= C	Charge conditions

i - codes:

i = blank	Values at current time interval and current iteration
i = 1	Converged values from last time interval
i = 2	Values from last iteration, current interval
i = 3	Scratch area
G	Specific heat ratio (γ)
R	Ideal gas constant
PERR	Error limit on pressures
KF	Flap flow coefficient
KW	Wall crossflow coefficient
TSL	Test section length
TSH	Height

Table C-2. Continued.

TSW	Width
TSP	Perimeter
TSA	Flow area
TSWA	Wall surface area
TSV	Volume
CTD	Charge tube diameter
CTA	Charge tube flow area
PV	Plenum volume
PVOTSV	Plenum: test section volume ratio
TAUW	Porosity
T	Time at end of current interval (t)
Tl	Time at end of last interval (t - Δt)
DT	Time increment
TSTR	Midpoint of current interval (t*)
TSTOP	Time for termination of run
Ai	Miscellaneous program constants

INTEGER ARRAYS

INSTR	Program control instructions (see input)
NVT	Number of time points in each of three input area-time curves

INTEGER SCALARS

IDBUG	Data set reference number (DSRN) for debugging output, normally dummied
IIN	DSRN of input data (usually 05 for card reader)
IOUT	DSRN of primary output data (usually 06 for line printer)
ITER	Number of iterations

Table C-2. Concluded.

NP Printing time interval

IFLGi Miscellaneous program control flags

Table C-3. Description of Program Input
a. Main Program

Variable	Index	Value	Action	Default Value	Format
NCTL		0	Proceed through normal programmed solution procedure	0	I3
		1	Read INSTR(*)		
		2	Write heading		
		3	Read data file and print results		
		4	Proceed to normal calculation		
		5	Call INPUT		
		6	Call INIT		
		7	Call CONST		
		8	Call DUMP		
		9	Call SOLVER		
		10	Call PRINT		
		11	Call BINOM		
		12	Call REVERT		
		13	Stop		
INSTR	1	06	Print debugging data		26I3
		03	Skip debugging prints (DSRN 03 Is Dummy)	03	(One Card)
	2	05	Input DSRN	05	
	3	06	Output DSRN	06	
	4		Printing time interval	1	
	5	1	Pressures in psf		
		2	Pressures in psi	2	
	6	≠0	Call PRINT on every iteration {		
		0	Call PRINT on ON convergence } set to zero when IDEBUG = IOUT	0	
	7	1	Extrapolate to next time interval as an initial guess		
		2	Do not extrapolate	2	
	8	1	Use reverted series from mass flux - Mach number wave equation		
		2	Use second-degree approximation	1	
	9	1	Use iterative solution to energy and wave equations		
		2	Use approximate expansions for energy and wave equations	1	
	10	1	Average current value with previous average value		
		2	Average current value with previous unaveraged value		
		0	Do not invoke option	0	

Table 3. Continued
a. Concluded

Variable	Index	Value	Action	Default Value	Format
INSTR	11	>0	Iteration limit beyond which current weight is halved	0	
	12	1	Do not invoke option		
		>1	Divide error limits PERR and \$EMAX by INSTR(12) if the fractional difference between successive time intervals is less than (errors) x (INSTR(12))	1	
	13	≠0	Print only time and pressure data		
		0	Print everything		
	14	>1	Set DT = DT*INSTR(14) based on INSTR(12) criteria, do not cut error limits		
		1	Do not invoke option	1	
	15	≠03	Read solution from DSRN = INSTR(15), skip other input		
		03	Do not read solution	03	
	16	[1,1000] ^a	First record number to be read	0	
	17	[1,1000]	Last record number to be read	0	
	18	≠03	Write solution on DSRN = INSTR(18)		
		03	Do not write solution	03	
	19	[1,1000]	First record number to be written	0	
	20	0	Do not invoke option		
		>0	When weight is halved, increment INSTR(11) by INSTR(20)	0	
	21	0	Do not invoke option		
		≠0	Set INSTR(7) = 2 to extrapolate next time interval when weight is halved	0	
	22	0	Do not invoke option, set INSTR(22) = 2 ³¹ -1		
		>0	Set INSTR(23) = 2 when number of iterations > INSTR(22)	9999999	
	23	0	Do not use small perturbation expansion		
		1	Use small perturbation initial guess for next time interval		
		2	Use small perturbation expansions as solution	0	
	24	≠0	SMPERT prints small perturbation results		
		0	Does not print without error	0	
	25	1	Use isentropic solution in plenum		
		2	Use anisentropic solution in plenum	2	
	26	0	Do not invoke option		
		>0	Revert to exact equation after the input number of time increments beyond choking	9999999	

^aSquare brackets [] indicate the range of the variable.

Table 3. Concluded
b. Subroutine INPUT

Variable, units	Card Number ^a	Value	Meaning	Default Value	Format
NVT(1) NVT(2) NVT(3)	1	[2,50] [2,50] [2,50]	Number of area-time points for main valve Number of area-time points for plenum exhaust valve Number of area-time points for flaps		2613
PC, psia TC, °R	2		Charge pressure Charge temperature		5E16.8
TSL, ft TSH, ft TSW, ft CTD, ft PVOTSV	3		Test section length Test section height Test section width Charge tube diameter Ratio of plenum volume to test section volume		5E16.8
TAUW KW, ft/sec KF, ft/sec A15 ^b A16 ^b	4		Porosity (fraction, not percent) Wall crossflow coefficient { Flap flow coefficient { from Dr. Varner's flow model Crossflow constant MDPT = -AWOKW x (PP - A15 x PT) Flap flow constant MDF = -AF/KF x (PP - A16 x PD)	1.0 1.0	5E16.8
A17	5	>0	Test section pressure weight, PT = A17 x PD + (1.00 - A17) x PN	1.0	5E16.8
R, ft ² /sec ² -°R G A11 A13, sec A14	6	(0,1) ^c	Perfect gas constant Ratio of specific heats (γ) Fraction of new values to be accepted Set INSTR(23) = 2 When T > A13 ε12 and ε13 limits	0.5 1.070 -0.1	5E16.8
DT, sec TSTOP, sec \$EMAX PERR A10, sec	7	(0,1) (0,1)	Time increment for numerical calculation Time to halt calculation Maximum allowable error - used in SOLVER { Maximum allowable error - used in MAIN { fractions, not percent Time at which INSTR(6) is set different from zero	1.070	5E16.8
AREAM(1), ft ² AREAM(2), ft ² AREAM(3), ft ²	8		Maximum main valve flow area Maximum plenum exhaust flow area Maximum flap flow area		5E16.8
TVF(1), sec TVF(2), sec TVF(3), sec	9		Final time in main valve area-time curve Final time in plenum exhaust area-time curve Final time in flap area-time curve		5E16.8
TDELAY(1), sec TDELAY(2), sec TDELAY(3), sec	10		Time delay for main valve Time delay for plenum exhaust Time delay for flaps		5E16.8
TV(1,I) AREA(1,I) TV(2,I) AREA(2,I) TV(3,I) AREA(3,I)		{ [0.,1.] }	Nondimensional time (final = 1.0) and nondimensional area (maximum = 1.0) } for { { main valve } { plenum } { exhaust } { flaps } }		2E16.8
I1d I2 I3 ISTR RSTR V	1	0 1 2 3 >1 >1 >1	Return 1 Read ISTR(I2) Read RSTR(I2) Read V(I2,I3) Indices of array elements to be read } one card } Enter one per card each preceded by a no. 1 card above - see common and equivalence statements to determine indices		2613 13 E16.8 E16.8

^aCard Order in Input Deck^bIf Less than Zero, Ramps of Fig. 27b Will Be Used^cRound Brackets Exclude End Points^dThese Cards Omitted Unless NCTL ≠ 0Note: If INSTR(5) = 1, any set of units for which $G_c = 1$ in $F = 1/G_c$ ma will work properly.

Table C-4. Suggested Values for INSTR(I)

I	Suggested Value of INSTR (I)
1	03
2	05
3	06
4	01
5	02
6	00
7	02
8	01
9	01
10	01
11	40 ^a
12	10
13	0 or 1
14	01
15	03
16	00
17	00
18	03
19	00
20	10 ^a
21	00
22	01
23	01
24	00
25	02
26	09 ^a

(a) Adjustment May Be Necessary for Specific Cases

Table C-5. Recommended Values for Certain Variables

Variable Names	Recommended Value
KW, KF	See Fig. 7
A15, A16	See Fig. 27a, Enter Negative
A11	0.5 or Leave Blank
A13	Leave Blank
A14	-0.2
A17	0.9
PERR	0.49999999E-04
\$EMAX	0.49999999E-05

Table C-6. Sample Jobstream and Input Data Deck

```

/*PRIORITY      2
//VKF05145      JOB          (ARO,
//  VRV00090,01,V37A=31A)009452SHOPE,MSGLEVEL=(2,0),CLASS=A,TIME=3
// EXEC FORTEPDS,PGMNO=VRV00090
//GO.FT03F001 DD DUMMY
//GO.FT05F001 DD *
000
                                02          01 20 10                                10 00 20 01 00          09
02 10 02
0.15215000E+03 0.53000000E+03
2.11400000E 00 0.61170000E 00 0.76330000E 00 1.16200000E 00 2.50000000E 00
0.04000000E 00 0.31000000E+03 0.20000000E+03 -1.04988410E 00 -1.08312800E 00
0.90000000E 00
0.17176000E+04 1.40000000E 00
0.00100000E 00 0.10000000E 00 0.49999999E-05 0.49999999E-04 -0.20000000E 00
0.46591116E 00 0.90371714E -1 0.09167000E 00
0.03000000E 00 0.04000000E 00 0.00000000E 00
0.00000000E 00 0.00500000E 00 0.00000000E 00
0.00000000E 00 0.00000000E 00
1.00000000E 00 1.00000000E 00
0.00000000E 00 0.00000000E 00
0.16400000E 00 0.92300000E 00
0.20000000E 00 0.98000000E 00
0.30000000E 00 1.00000000E 00
0.45000000E 00 0.99298055E 00
0.50000000E 00 0.97332608E 00
0.66000000E 00 0.64902737E 00
0.80000000E 00 0.52478305E 00
0.90000000E 00 0.49810913E 00
1.00000000E 00 0.48687801E 00
0.00000000E 00 1.00000000E 00
1.00000000E 00 1.00000000E 00
/*

```


TABLE C-7
SAMPLE OUTPUT FROM HIRTSM1 FOR RUN 2742

G	GM1	GP1	GM102	GP102	00G	GM10G	GP10G
0.14000000 01	0.40000000 00	0.24000000 01	0.20000000 00	0.12000000 01	0.71428571 00	0.28571429 00	0.17142857 01
G0GM1	G0GP1	SGM102	T0G	T0GP1	00GM1	GPOGM1	GPGM12
0.35000000 01	0.58333333 00	0.44721360 00	0.14285714 01	0.83333333 00	0.25000000 01	0.60000000 01	0.30000000 01
T0GM1	MGP0GM2	MGP0GM	00GP1	R	00R	GR	DT02
0.50000000 01	-0.30000000 01	-0.60000000 01	0.41666667 00	0.17176080 04	0.58220502 03	0.24046512 04	0.50000000 03
DT0PY	00KF	INFIN	00A1	00DT	M00GM1	TM00GS	GP02G5
0.40524036 03	0.50000000 02	0.99999945 70	0.60525968 02	0.10000000 04	-0.35000000 01	0.30612245 00	0.61224490 00
SGOR							
0.28549729 01							

V EQUIVALENC ARRAY

0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.0
0.0	0.0	0.0	0.0	0.0	0.12686220 02	0.28813801 02	0.53000000 03
0.53000000 03	0.53000000 03	0.53000000 03	0.24067696 01	0.24067696 01	0.24067696 01	0.24067696 01	0.11289221 04
0.0	0.0	0.0	0.0	0.0	0.0	0.15215000 03	0.15215000 03
0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.0	0.0	0.0
0.0	0.0	0.0	0.12686220 02	0.28813801 02	0.53000000 03	0.53000000 03	0.53000000 03
0.53000000 03	0.24067696 01	0.24067696 01	0.24067696 01	0.24067696 01	0.11289221 04	0.0	0.0
0.0	0.91670000 01	0.0	0.0	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03
0.15215000 03	0.15215000 03	0.15215000 03	0.0	0.0	0.0	0.0	0.0
0.0	0.12686220 02	0.28813801 02	0.53000000 03	0.53000000 03	0.53000000 03	0.53000000 03	0.24067696 01
0.24067696 01	0.24067696 01	0.24067696 01	0.11289221 04	0.0	0.0	0.0	0.0
0.0	0.0	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03	0.15215000 03
0.15215000 03	0.0	0.0	0.0	0.0	0.0	0.0	0.12686220 02
0.28813801 02	0.53000000 03	0.53000000 03	0.53000000 03	0.53000000 03	0.24067696 01	0.24067696 01	0.24067696 01
0.24067696 01	0.11289221 04	0.0	0.0	0.0	0.0	0.0	0.0

RW(1)	RW(2)	RW(3)	RW(4)	RW(5)	RW(6)	RW(7)
0.10000000 01	0.12000000 01	0.20400000 01	0.40480000 01	0.87696000 01	0.20106240 02	0.47961472 02

INSTR(1)= 3 SEND DEBUGGING OUTPUT TO DSRN 3

INSTR(2)= 5 OBTAIN INPUT FROM DSRN 5

INSTR(3)= 6 SEND REGULAR OUTPUT TO DSRN 6

INSTR(4)= 1 PRINTING TIME INTERVAL: 1

INSTR(5)= 2 INPUT AND OUTPUT PRESSURES IN PSIA

INSTR(6)= 0 PRINT DATA ONLY WHEN CONVERGED

INSTR(7)= 2 DO NOT EXTRAPOLATE TO NEXT TIME INTERVAL

INSTR(8)= 1 USE SEVENTH DEGREE REVERTED SERIES AS INITIAL GUESS TO MASS FLUX-MACH NUMBER WAVE EQUATION

INSTR(9)= 1 USE ITERATIVE SOLUTION TO ENERGY AND WAVE EQUATIONS

INSTR(10)= 1 AVERAGE VALUES OF CURRENT ITERATION WITH AVERAGE VALUES OF PREVIOUS ITERATION

INSTR(11)= 20 CURRENT WEIGHT IS HALVED BEYOND 20 ITERATIONS

INSTR(12)= 10 DIVIDE ERRORS BY 10 WHEN TIME-DIFFERENCES ARE LESS THAN 10 TIMES THE ERRORS
 INSTR(13)= 0 PRINT ALL DATA
 INSTR(14)= 1 DO NOT INVOKE DT-RAISING OPTION
 INSTR(15)= 3 DO NOT READ SOLUTION FROM PERMANENT DATA SET
 INSTR(16)= 0 FIRST RECORD TO BE READ: 0
 INSTR(17)= 0 LAST RECORD TO BE READ: 0
 INSTR(18)= 3 DO NOT WRITE SOLUTION ON PERMANENT DATA SET
 INSTR(19)= 0 FIRST RECORD TO BE WRITTEN: 0
 INSTR(20)= 10 INCREMENT INSTR(11) BY 10 WHENEVER WEIGHT IS HALVED
 INSTR(21)= 0 DO NOT CHANGE EXTRAPOLATION OPTION (INSTR(7))
 INSTR(22)= 20 SET INSTR(23)=2 WHEN ITER >= 20
 INSTR(23)= 1 USE SMALL PERTURBATION EXPANSIONS AS INITIAL GUESS FOR NEXT TIME INTERVAL
 INSTR(24)= 0 RESULTS FROM SMPRT NOT PRINTED
 INSTR(25)= 2 SET TP AND IPT = MAX(ISEN TP, TCTU)
 INSTR(26)= 9 REVERT TO EXACT SUPERSONIC SOLUTION 9 TIME INCREMENTS AFTER CHOKE

J 1	J 2	J 3	J 4	J 5	J 6	J 7	J 8	J 9	J 10	J 11	J 12	J 13	J 14	J 15	J 16	J 17	J 18	J 19	J 20	J 21	J 22	J 23	J 24	J 25	J 26
3	5	6	1	2	0	2	1	1	1	20	10	0	1	3	0	0	3	0	10	0	20	1	0	2	9

FLOW AREAS VERSUS TIME

I	TV(1,I)	AREA(1,I)	TV(2,I)	AREA(2,I)	TV(3,I)	AREA(3,I)
1	0.0	0.0	0.0	0.0	0.0	0.91670000D-01
2	0.30000000D-01	0.46591116D 00	0.50000000D-02	0.0	0.0	0.91670000D-01
3	0.0	0.0	0.11560000D-01	0.83*13092D-01	0.0	0.0
4	0.0	0.0	0.13000000D-01	0.88564280D-01	0.0	0.0
5	0.0	0.0	0.17000000D-01	0.90371714D-01	0.0	0.0
6	0.0	0.0	0.23000000D-01	0.89/37354D-01	0.0	0.0
7	0.0	0.0	0.25000000D-01	0.87901146D-01	0.0	0.0
8	0.0	0.0	0.31400000D-01	0.58653716D-01	0.0	0.0
9	0.0	0.0	0.37000000D-01	0.47425544D-01	0.0	0.0
10	0.0	0.0	0.41000000D-01	0.45014976D-01	0.0	0.0
11	0.0	0.0	0.45000000D-01	0.44000000D-01	0.0	0.0

T PCT0 MDF TCT0 E(1)	TSTR PE0 MDE RP E(2)	T1 MCT MDE RPT E(3)	PT MD MDT50 REU E(4)	PP MN MOCT0 RCT0 E(5)	PD MDCT TP AE E(6)	PN MDPT TPT APE E(7)	PPT MOD TE0 DT	ND NM NCT ITER J16
0.0	0.0	0.0	0.152150D 03	0.152150D 03	0.152150D 03	0.152150D 03	0.152150D 03	0
0.152150D 03	0.152150D 03	0.0	0.0	0.0	0.0	0.0	0.0	0
0.0	0.0	0.0	0.126862D 02	0.288138D 02	0.530000D 03	0.530000D 03	0.530000D 03	0
0.530000D 03	0.240677D-01	0.240677D-01	0.240677D-01	0.240677D-01	0.0	0.0	0.0	0
0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.100000D-02	-1
0.100000D-02	0.500000D-03	0.0	0.151568D 03	0.152074D 03	0.151570D 03	0.151565D 03	0.151998D 03	21
0.151569D 03	0.151569D 03	0.273812D-02	0.694732D-02	0.621917D-02	0.786372D-01	-0.927192D-02	0.878475D-01	19
-0.331411D-01	0.121702D 00	0.0	0.126447D 02	0.287194D 02	0.529924D 03	0.529849D 03	0.529421D 03	19
0.529421D 03	0.240591D-01	0.240505D-01	0.240020D-01	0.240020D-01	0.776519D-02	0.0	0.916700D-01	3
0.442799D-04	0.989949D-05	0.198061D-04	0.782624D-05	0.182267D-04	0.436644D-04	0.436644D-04	0.100000D-02	-1
PERR CUT TO 0.4999999D-05 AND SEMAX CUT TO 0.4999999D-06								
0.200000D-02	0.150000D-02	0.100000D-02	0.150304D 03	0.151785D 03	0.150299D 03	0.150308D 03	0.151573D 03	18
0.150348D 03	0.150348D 03	0.855315D-02	0.215218D-01	0.194300D-01	0.243934D 00	-0.262455D-01	0.270181D 00	25
-0.923444D-01	0.362571D 00	0.0	0.125573D 02	0.285211D 02	0.529647D 03	0.529425D 03	0.528199D 03	19
0.528199D 03	0.240265D-01	0.240025D-01	0.238638D-01	0.238638D-01	0.232956D-01	0.0	0.916700D-01	10
0.891248D-06	0.663618D-06	0.132856D-05	0.226175D-05	0.157648D-05	0.852219D-06	0.852219D-06	0.100000D-02	-1
0.300000D-02	0.250000D-02	0.200000D-02	0.148907D 03	0.151254D 03	0.148895D 03	0.148919D 03	0.150935D 03	21
0.149039D 03	0.149039D 03	0.148910D-01	0.370687D-01	0.338401D-01	0.421478D 00	-0.401483D-01	0.461626D 00	21
-0.138061D 00	0.599771D 00	0.0	0.124635D 02	0.283080D 02	0.529106D 03	0.528787D 03	0.526881D 03	19
0.526881D 03	0.239663D-01	0.239302D-01	0.237151D-01	0.237151D-01	0.388259D-01	0.0	0.916700D-01	10
0.198541D-08	0.119035D-05	0.238428D-05	0.458196D-05	0.229179D-05	0.258359D-08	0.258359D-08	0.100000D-02	-1
0.400000D-02	0.350000D-02	0.300000D-02	0.147420D 03	0.150526D 03	0.147398D 03	0.147442D 03	0.150118D 03	24
0.147690D 03	0.147690D 03	0.215295D-01	0.531828D-01	0.489560D-01	0.604559D 00	-0.520260D-01	0.656585D 00	21
-0.176481D 00	0.833162D 00	0.0	0.123668D 02	0.280883D 02	0.528378D 03	0.527968D 03	0.525514D 03	18
0.525514D 03	0.238839D-01	0.238376D-01	0.235616D-01	0.235616D-01	0.543563D-01	0.0	0.916700D-01	10
0.355785D-06	0.138066D-05	0.276666D-05	0.430807D-05	0.233162D-05	0.313059D-06	0.313059D-06	0.100000D-02	-1
0.500000D-02	0.450000D-02	0.400000D-02	0.145855D 03	0.149633D 03	0.145819D 03	0.145890D 03	0.149149D 03	18
0.146317D 03	0.146317D 03	0.284065D-01	0.697672D-01	0.646495D-01	0.791146D 00	-0.622755D-01	0.853422D 00	21
-0.209144D 00	0.106267D 01	0.0	0.122682D 02	0.278643D 02	0.527480D 03	0.526992D 03	0.524113D 03	18
0.524113D 03	0.237826D-01	0.237276D-01	0.234050D-01	0.234050D-01	0.698867D-01	0.0	0.916700D-01	10
0.451639D-06	0.148317D-05	0.297319D-05	0.371255D-05	0.208170D-05	0.385848D-06	0.385848D-06	0.100000D-02	-1
0.600000D-02	0.550000D-02	0.500000D-02	0.144154D 03	0.148439D 03	0.144105D 03	0.144203D 03	0.147731D 03	9
0.144875D 03	0.144875D 03	0.357640D-01	0.872758D-01	0.814911D-01	0.987357D 00	-0.694685D-01	0.105683D 01	17
-0.230891D 00	0.128783D 01	-0.978732D-01	0.121845D 02	0.276288D 02	0.526274D 03	0.525555D 03	0.522632D 03	17
0.522632D 03	0.236469D-01	0.235663D-01	0.232400D-01	0.232400D-01	0.854170D-01	0.635771D-02	0.916700D-01	10
0.924620D-06	0.178422D-05	0.358066D-05	0.339026D-05	0.215702D-05	0.765156D-06	0.765156D-06	0.100000D-02	-1
0.700000D-02	0.650000D-02	0.600000D-02	0.142203D 03	0.146678D 03	0.142141D 03	0.142264D 03	0.145627D 03	19
0.143275D 03	0.143275D 03	0.440922D-01	0.106629D 00	0.100638D 00	0.120527D 01	-0.707997D-01	0.127607D 01	18
-0.231349D 00	0.150757D 01	-0.290630D 00	0.120493D 02	0.273671D 02	0.524482D 03	0.523406D 03	0.520977D 03	17
0.520977D 03	0.234461D-01	0.233260D-01	0.230564D-01	0.230564D-01	0.100947D 00	0.190731D-01	0.916700D-01	10
0.173806D-05	0.241179D-05	0.484827D-05	0.324127D-05	0.248934D-05	0.138259D-05	0.138259D-05	0.100000D-02	-1

T	ISTR	T1	PT	PP	PD	PN	PPT	ND
PCTO	PEU	MCT	MD	MN	MDCT	MUPT	MDD	NM
MUF	MDE	MOPE	MDT50	MDCTO	TP	TPT	TEO	NCT
ICTO	RP	RPT	RLG	RCIO	AE	APE	AE	JIER
E(1)	E(2)	E(3)	E(4)	E(5)	E(6)	E(7)	DT	J16
0.800000-02	0.750000-02	0.700000-02	0.1399640 03	0.1442960 03	0.139940 03	0.1400330 03	0.1429590 03	20
0.1415080 03	0.1415080 03	0.5951150-01	0.1281090 00	0.1224250 00	0.1446460 01	-0.6588740-01	0.1512340 01	22
-0.2086850 00	0.1721100 01	-0.4776350 00	0.1128110 02	0.2707740 02	0.5220350 03	0.520590 03	0.5191320 03	16
0.5191320 03	0.2317460-01	0.2302120-01	0.2285280-01	0.2285280-01	0.118780 00	0.3178850-01	0.9167000-01	11
0.4968930-05	0.1057430-05	0.2128720-05	0.3510060-06	0.2310280-05	0.3829350-05	0.3829350-05	0.1000000-02	-1
0.9000000-02	0.8500000-02	0.8000000-02	0.1372490 03	0.1408870 03	0.1371840 03	0.1373130 03	0.1398190 03	20
0.1394500 03	0.1394500 03	0.6476820-01	0.1531790 00	0.1486910 00	0.1727530 01	-0.5070140-01	0.1778230 01	22
-0.1479720 00	0.1926240 01	-0.6560840 00	0.1177300 02	0.2673970 02	0.516950 03	0.516950 03	0.516950 03	16
0.516950 03	0.2284800-01	0.2267480-01	0.2261500-01	0.2261500-01	0.1320080 00	0.4450390-01	0.9167000-01	12
0.1364310-05	0.2734000-06	0.8438120-06	0.2777420-05	0.2070630-05	0.1014920-05	0.1014920-05	0.1000000-02	-1
0.1000000-01	0.9500000-02	0.9000000-02	0.1387400 03	0.1380760 03	0.1348770 03	0.1348030 03	0.1366540 03	17
0.1376390 03	0.1376390 03	0.7495780-01	0.1755770 00	0.1727320 00	0.1975350 01	-0.4235090-01	0.2017700 01	16
-0.1111340 00	0.2128870 01	-0.8282500 00	0.1156180 02	0.2644170 02	0.5150370 03	0.5150370 03	0.5150370 03	15
0.5150370 03	0.2247590-01	0.2227700-01	0.2240480-01	0.2240480-01	0.1475390 00	0.5721930-01	0.9167000-01	11
0.3136960-06	0.4882170-06	0.9232930-06	0.26650500-05	0.1167710-05	0.2111240-06	0.2111240-06	0.1000000-02	-1
0.1100000-01	0.1050000-01	0.1000000-01	0.1320880 03	0.1349770 03	0.1320250 03	0.1321340 03	0.1336250 03	23
0.1358010 03	0.1358010 03	0.8558320-01	0.2010990 00	0.1981250 00	0.2227230 01	-0.3183720-01	0.2259060 01	19
-0.6690610-01	0.2325990 01	-0.9914900 00	0.1150880 02	0.2613870 02	0.5130620 03	0.5130620 03	0.5130620 03	16
0.5130620 03	0.2205610-01	0.2183510-01	0.2215070-01	0.2219070-01	0.1630690 00	0.6993480-01	0.9167000-01	11
0.1595360-06	0.4231440-06	0.8254130-06	0.3346420-05	0.1594220-05	0.9954570-07	0.9954570-07	0.1000000-02	-1
0.1200000-01	0.1150000-01	0.1100000-01	0.1242720 03	0.1316420 03	0.1292340 03	0.1293110 03	0.1301830 03	20
0.1339460 03	0.1339460 03	0.9660340-01	0.2287770 00	0.2248750 00	0.2481580 01	-0.1970630-01	0.2501280 01	22
-0.1637000-01	0.2517670 01	-0.1145090 01	0.1137360 02	0.2533240 02	0.5110510 03	0.5110510 03	0.5110510 03	12
0.5110510 03	0.2159580-01	0.2133550-01	0.2197380-01	0.2197380-01	0.1785990 00	0.6265020-01	0.9167000-01	11
0.1100930-06	0.3276260-06	0.6839040-06	0.4024260-05	0.2066690-05	0.7428430-07	0.7428430-07	0.1000000-02	-1
0.1300000-01	0.1250000-01	0.1200000-01	0.1284010 03	0.1282660 03	0.1243430 03	0.1264180 03	0.1268520 03	17
0.1321380 03	0.1321380 03	0.1076610 00	0.2530210 00	0.2522220 00	0.2729910 01	-0.7998930-02	0.2737890 01	14
0.3599670-01	0.2704900 01	-0.1173650 01	0.1142180 02	0.2553320 02	0.5090700 03	0.5090700 03	0.5090700 03	11
0.5090700 03	0.2112370-01	0.2089100-01	0.2176150-01	0.2176150-01	0.1941300 00	0.8877900-01	0.9167000-01	12
0.6719840-06	0.3077570-07	0.5240270-07	0.2849080-05	0.1756270-05	0.3930860-06	0.3930860-06	0.1000000-02	-1
0.1400000-01	0.1350000-01	0.1300000-01	0.1235550 03	0.1250370 03	0.1235590 03	0.1235520 03	0.1236620 03	20
0.1304320 03	0.1304320 03	0.1183840 00	0.2791710 00	0.2793200 00	0.2964340 01	0.1430110-02	0.2962910 01	19
0.3997440-01	0.2888940 01	-0.1172850 01	0.1111730 02	0.2525040 02	0.5071840 03	0.5071840 03	0.5071840 03	10
0.5071840 03	0.2066660-01	0.2046910-01	0.2156040-01	0.2156040-01	0.2098600 00	0.8877900-01	0.9167000-01	12
0.5398320-06	0.2453560-07	0.3704320-07	0.2989280-05	0.1760900-05	0.2973000-06	0.2973000-06	0.1000000-02	-1
0.1500000-01	0.1450000-01	0.1400000-01	0.1207580 03	0.1219890 03	0.1207810 03	0.1207350 03	0.1207200 03	16
0.1288360 03	0.1288360 03	0.1287420 00	0.3051030 00	0.3059970 00	0.4143910 01	0.6282350-02	0.3175630 01	20
0.1052720 00	0.3070360 01	-0.1152100 01	0.1100050 02	0.2486540 02	0.5054020 03	0.5054020 03	0.5054020 03	16
0.5054020 03	0.2023580-01	0.2002540-01	0.2137170-01	0.2137170-01	0.2251900 00	0.8928210-01	0.9167000-01	12
0.4000520-06	0.1318680-07	0.3420770-07	0.2993650-05	0.1698090-05	0.2090110-06	0.2090110-06	0.1000000-02	-1

[illegible]

T PCTU MDF TCTU E(1)	TSTR PEU MDE RPE E(2)	T1 MGT MDPE RPT E(3)	PT MU MDTSO REU E(4)	PP MN MDCTO RCTU E(5)	PD MDCI TP AE E(6)	PM MDPT TPT APE E(7)	PPT MDD TEO AF DT	NM NCT ITER J16
0.2400000-01	0.2350000-01	0.2300000-01	0.9630180 02	0.9926770 02	0.9624790 02	0.9635560 02	0.9821410 02	21
0.1179810 03	0.1179810 03	0.2072310 00	0.5472500 00	0.5457020 00	0.4679880 01	-0.8782530-02	0.4688660 01	16
0.7414950-01	0.4614510 01	-0.9499250 00	0.1020130 02	0.2316980 02	0.4928520 03	0.4928520 03	0.4928520 03	16
0.4928520 03	0.1688610-01	0.1670690-01	0.2006950-01	0.2006950-01	0.3649640 00	0.8929330-01	0.9167000-01	11
0.1504470-06	0.7129180-07	0.1322010-06	0.4064640-05	0.1955920-05	0.4075270-07	0.4075270-07	0.1000000-02	-1
0.2500000-01	0.2450000-01	0.2400000-01	0.9347310 02	0.9693530 02	0.9336010 02	0.9358600 02	0.9587540 02	17
0.1170640 03	0.1170640 03	0.2146640 00	0.5778380 00	0.5746430 00	0.4806410 01	-0.1671970-01	0.4823130 01	17
0.4434860-01	0.4778780 01	-0.9194050 00	0.1013320 02	0.2301520 02	0.4917540 03	0.4917540 03	0.4917540 03	16
0.4917540 03	0.1652620-01	0.1634550-01	0.1995780-01	0.1995780-01	0.3804940 00	0.8840520-01	0.9167000-01	11
0.1448180-06	0.6518980-07	0.1212910-06	0.3877490-05	0.1863910-05	0.3601700-07	0.3601700-07	0.1000000-02	-1
0.2600000-01	0.2550000-01	0.2500000-01	0.9060070 02	0.9462030 02	0.9040690 02	0.9079440 02	0.9356720 02	17
0.1162030 03	0.1162030 03	0.2217710 00	0.6097260 00	0.6043210 00	0.4925130 01	-0.2585720-01	0.4950990 01	17
0.4925270-02	0.4942460 01	-0.8786140 00	0.1006930 02	0.2287000 02	0.4907170 03	0.4907170 03	0.4907170 03	17
0.4907170 03	0.1616560-01	0.1598570-01	0.1985290-01	0.1985290-01	0.3960240 00	0.8567150-01	0.9167000-01	11
0.1139590-06	0.4992980-07	0.9357150-07	0.3766770-05	0.1812260-05	0.2534750-07	0.2534750-07	0.1000000-02	-1
0.2700000-01	0.2650000-01	0.2600000-01	0.8769650 02	0.9234630 02	0.8739160 02	0.8800140 02	0.9130930 02	17
0.1154060 03	0.1154060 03	0.2284680 00	0.6429800 00	0.6345710 00	0.5035000 01	-0.3640510-01	0.5071400 01	11
-0.3466770-01	0.5106970 01	-0.8050650 00	0.1001010 02	0.2273550 02	0.4897530 03	0.4897530 03	0.4897530 03	16
0.4897530 03	0.1580810-01	0.1563060-01	0.1975550-01	0.1975550-01	0.4115550 00	0.8109220-01	0.9167000-01	11
0.1042410-06	0.3332320-07	0.6361590-07	0.3599960-05	0.1741420-05	0.1295070-07	0.1295070-07	0.1000000-02	-1
0.2800000-01	0.2750000-01	0.2700000-01	0.8475700 02	0.9010850 02	0.8430020 02	0.8521380 02	0.8907330 02	18
0.1146750 03	0.1146750 03	0.2347110 00	0.6778650 00	0.6653670 00	0.5125670 01	-0.4834940-01	0.5184020 01	13
-0.8596600-01	0.5269990 01	-0.7818680 00	0.9955720 01	0.2261210 02	0.4888650 03	0.4888650 03	0.4888650 03	14
0.4888650 03	0.1545310-01	0.1527550-01	0.1966610-01	0.1966610-01	0.4270850 00	0.7651290-01	0.9167000-01	11
0.6332840-07	0.2721640-07	0.5183590-07	0.3338860-05	0.1628600-05	0.1104320-07	0.1104320-07	0.1000000-02	-1
0.2900000-01	0.2850000-01	0.2800000-01	0.8179400 02	0.8787800 02	0.8109940 02	0.8242860 02	0.8683050 02	14
0.1140110 03	0.1140110 03	0.2404780 00	0.7168880 00	0.6967930 00	0.5227210 01	-0.6150160-01	0.5288710 01	11
-0.1457780 00	0.5434490 01	-0.6807680 00	0.9906260 01	0.2249980 02	0.4880550 03	0.4880550 03	0.4880550 03	14
0.4880550 03	0.1509560 01	0.1491570-01	0.1958460-01	0.1958460-01	0.4426160 00	0.7193360-01	0.9167000-01	11
0.4239890-07	0.2062460-07	0.3954830-07	0.3092150-05	0.1512140-05	0.7512800-08	0.7512800-08	0.1000000-02	-1
0.3000000-01	0.2950000-01	0.2900000-01	0.7870050 02	0.8562620 02	0.7774760 02	0.7965350 02	0.8455170 02	14
0.1134150 03	0.1134150 03	0.2457240 00	0.7547040 00	0.7288160 00	0.5309280 01	-0.7563930-01	0.5384920 01	15
-0.2150540 00	0.5599970 01	-0.6215610 00	0.9861870 01	0.2239890 02	0.4873240 03	0.4873240 03	0.4873240 03	14
0.4873240 03	0.1473080-01	0.1454600-01	0.1951150-01	0.1951150-01	0.4581460 00	0.6735440-01	0.9167000-01	11
0.1866950-06	0.9082520-08	0.1103180-07	0.2969260-05	0.1561130-05	0.2513910-07	0.2513910-07	0.1000000-02	-1
0.3100000-01	0.3050000-01	0.3000000-01	0.7633100 02	0.8343330 02	0.7521410 02	0.7744800 02	0.8240560 02	14
0.1129890 03	0.1129890 03	0.2495180 00	0.7851820 00	0.7588110 00	0.5367880 01	-0.7765040-01	0.5445530 01	14
-0.2310430 00	0.5676560 01	-0.5647770 00	0.9830140 01	0.2232690 02	0.4868010 03	0.4868010 03	0.4868010 03	18
0.4868010 03	0.1436900-01	0.1419200-01	0.1945920-01	0.1945920-01	0.4659110 00	0.6277510-01	0.9167000-01	14
0.3216940-05	0.3737040-06	0.6385690-06	0.2708810-05	0.2966590-05	0.4035020-06	0.4035020-06	0.1000000-02	-1

T	ISTK	IL	PI	PP	PO	PN	PPT	ND
PC10	PEU	MCI	MDTSU	MN	MDCT	MDPT	MDD	NM
MUF	MDE	MOPE	MDTSU	MN	MDCT	MDPT	TEU	NCT
TC10	RP	RPT	REU	RCTU	AE	APE	AF	ITER
E(1)	E(2)	E(3)	E(4)	E(5)	E(6)	E(7)	DT	J16
0.320000-01	0.315000-01	0.310000-01	0.7480920 02	0.8144010 02	0.7373650 02	0.7588160 02	0.8053220 02	11
0.1127130 03	0.1127130 03	0.2519960 00	0.8026020 00	0.7735840 00	0.5405990 01	-0.6795810-01	0.5473850 01	16
-0.1907930 00	0.5684670 01	-0.5135040 00	0.9009550 01	0.2228010 02	0.4864610 03	0.4864610 03	0.4864610 03	15
0.4864610 03	0.1403550-01	0.1387900-01	0.1442520-01	0.1942520-01	0.4659110 00	0.5845320-01	0.9167000-01	19
0.4748090-05	0.5680200-06	0.9915150-06	0.3990720-05	0.4374940-05	0.5385210-06	0.5385210-06	0.1000000-02	-1
SWITCHING TO SMALL PERTURBATION SOLUTION ENTIRELY								
0.3300000-01	0.3250000-01	0.3200000-01	0.7358740 02	0.7970360 02	0.7256980 02	0.7460500 02	0.7883420 02	0
0.1125040 03	0.1125040 03	0.2538820 00	0.8168610 00	0.7890840 00	0.5434710 01	-0.5944390-01	0.5494160 01	0
-0.1150660 00	0.5652650 01	-0.4849540 00	0.9793930 01	0.2224460 02	0.4864610 03	0.4864610 03	0.4862030 03	0
0.4862030 03	0.1373660-01	0.1359420-01	0.1373940-01	0.1939940-01	0.4659110 00	0.5644820-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3400000-01	0.3350000-01	0.3300000-01	0.7244450 02	0.7812590 02	0.7148720 02	0.7362180 02	0.7733170 02	0
0.1123230 03	0.1123230 03	0.2555180 00	0.8300650 00	0.8036180 00	0.5459570 01	-0.5210240-01	0.5511680 01	0
-0.1305670 00	0.5642270 01	-0.4580990 00	0.9780460 01	0.2221400 02	0.4864610 03	0.4862030 03	0.4859800 03	0
0.4859800 03	0.1346440-01	0.1333450-01	0.1937710-01	0.1937710-01	0.4659110 00	0.5444320-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3500000-01	0.3450000-01	0.3400000-01	0.7140330 02	0.7664720 02	0.7048170 02	0.7232490 02	0.7592200 02	0
0.1121670 03	0.1121670 03	0.2569360 00	0.8422620 00	0.8172430 00	0.5481930 01	-0.4576770-01	0.5528500 01	0
-0.1064900 00	0.5633310 01	-0.4327420 00	0.9768620 01	0.2218760 02	0.4864750 03	0.4859800 03	0.4857870 03	0
0.4857870 03	0.1321600-01	0.1309750-01	0.1935790-01	0.1935790-01	0.4659110 00	0.5253810-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3600000-01	0.3550000-01	0.3500000-01	0.7042750 02	0.7537250 02	0.6954710 02	0.7130780 02	0.7463520 02	0
0.1120330 03	0.1120330 03	0.2581630 00	0.8542250 00	0.8300100 00	0.5499540 01	-0.4402980-01	0.5539840 01	0
-0.8572280-01	0.5655580 01	-0.4087150 00	0.9758790 01	0.2216480 02	0.4864870 03	0.4857870 03	0.4856210 03	0
0.4856210 03	0.1298910-01	0.1288080-01	0.1934140-01	0.1934140-01	0.4659110 00	0.5083310-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3700000-01	0.3650000-01	0.3600000-01	0.5952150 02	0.7417070 02	0.6867810 02	0.7036490 02	0.7386220 02	0
0.1119170 03	0.1119170 03	0.2592250 00	0.8652290 00	0.8419650 00	0.5515490 01	-0.3557350-01	0.5551070 01	0
-0.6782620-01	0.5618920 01	-0.3858700 00	0.9750130 01	0.2214310 02	0.4865010 03	0.4856210 03	0.4854770 03	0
0.4854770 03	0.1278160-01	0.1268250-01	0.1932710-01	0.1932710-01	0.4659110 00	0.4882810-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3800000-01	0.3750000-01	0.3700000-01	0.5864750 02	0.7306450 02	0.6786510 02	0.6948600 02	0.7237630 02	0
0.1118170 03	0.1118170 03	0.2601480 00	0.8756430 00	0.8532150 00	0.5529320 01	-0.3166640-01	0.5526070 01	0
-0.5233470-01	0.5613140 01	-0.3696500 00	0.9742630 01	0.2212810 02	0.4865160 03	0.4854770 03	0.4853530 03	0
0.4853530 03	0.1259060-01	0.1244870-01	0.1934140-01	0.1934140-01	0.4659110 00	0.4471240-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0
0.3900000-01	0.3850000-01	0.3800000-01	0.5787550 02	0.7203260 02	0.6709460 02	0.6865630 02	0.7135770 02	0
0.1117290 03	0.1117290 03	0.2609580 00	0.8855940 00	0.8639290 00	0.5541420 01	-0.2784990-01	0.5559280 01	0
-0.3878200-01	0.5604080 01	-0.3596540 00	0.9730060 01	0.2211320 02	0.4865430 03	0.4853530 03	0.4852440 03	0
0.4852440 03	0.1261240-01	0.1232600-01	0.1930380-01	0.1930380-01	0.4659110 00	0.4652160-01	0.9167000-01	1
0.4428960-04	0.5331820-05	0.9479610-05	0.3964650-04	0.4200070-04	0.4442130-05	0.4442130-05	0.1000000-02	0

[illegible]

T PCTO MDF TCTO E(1)	TSTR PEO MDE HP E(2)	T1 MCT MUPE HPT E(3)	PT MD MDT50 REU E(4)	PP MN MDCTO RCTO E(5)	PD MDC TP AE E(6)	PN MUPT TPT APE E(7)	PPT MDD TEU AF DT	NU NM NCT ITER J16
0.4800000-01 0.1112940 03 0.2649000-01 0.4847040 03 0 0.4428960-04	0.4750000-01 0.1112940 03 0.5583110 01 0.1116140-01 0.5331820-05	0.4700000-01 0.2649820 00 -0.3054250 00 0.1110320-01 0.9479610-05	0.6190880 02 0.9639200 00 0.9703620 01 0.1925020-01 0.3964650-04	0.6479620 02 0.9445530 00 0.2203950 02 0.1925020-01 0.4200070-04	0.6129100 02 0.5601160 01 0.4867060 03 0.4659110 00 0.4442130-05	0.6264660 02 -0.8408780-02 0.4847360 03 0.4400000-01 0.4442130-05	0.6419590 02 0.5609570 01 0.4847040 03 0.9167000-01 0.1000000-02	0 0 0 1 0
0.4900000-01 0.1112730 03 0.2929050-01 0.4846780 03 0 0.4428960-04	0.4850000-01 0.1112730 03 0.5581880 01 0.1104640-01 0.5331820-05	0.4800000-01 0.2651810 00 -0.3022660 00 0.1098970-01 0.9479610-05	0.6140570 02 0.9720260 00 0.9702020 01 0.1924750-01 0.3964650-04	0.6413160 02 0.9522300 00 0.2203590 02 0.1924750-01 0.4200070-04	0.6071490 02 0.5604100 01 0.4867260 03 0.4659110 00 0.4442130-05	0.6209640 02 -0.7041180-02 0.4847040 03 0.4400000-01 0.4442130-05	0.6353660 02 0.5611160 01 0.4846780 03 0.9167000-01 0.1000000-02	0 0 0 1 0
0.5000000-01 0.1112550 03 0.3138880-01 0.4846560 03 0 0.4428960-04	0.4950000-01 0.1112550 03 0.5580870 01 0.1093430-01 0.5331820-05	0.4900000-01 0.2653450 00 -0.2991830 00 0.1087890-01 0.9479610-05	0.6085600 02 0.9801140 00 0.9700710 01 0.1924540-01 0.3964650-04	0.6348320 02 0.9596490 00 0.2203290 02 0.1924540-01 0.4200070-04	0.6014440 02 0.5606510 01 0.4867470 03 0.4659110 00 0.4442130-05	0.6156760 02 -0.5717340-02 0.4846780 03 0.4400000-01 0.4442130-05	0.6289240 02 0.5612230 01 0.4846560 03 0.9167000-01 0.1000000-02	0 0 0 1 0
0.5100000-01 0.1112410 03 0.3280400-01 0.4846380 03 0 0.4428960-04	0.5050000-01 0.1112410 03 0.5580060 01 0.1082460-01 0.5331820-05	0.5000000-01 0.2654760 00 -0.2961680 00 0.1077030-01 0.9479610-05	0.6032000 02 0.9882030 00 0.9699660 01 0.1924360-01 0.3964650-04	0.6284920 02 0.9667890 00 0.2203050 02 0.1924360-01 0.4200070-04	0.5957800 02 0.5608440 01 0.4867680 03 0.4659110 00 0.4442130-05	0.6106200 02 -0.4399680-02 0.4846560 03 0.4400000-01 0.4442130-05	0.6226220 02 0.5612840 01 0.4846380 03 0.9167000-01 0.1000000-02	0 0 0 1 0
0.5200000-01 0.1112310 03 0.3354930-01 0.4846250 03 0 0.4428960-04	0.5150000-01 0.1112310 03 0.5579440 01 0.1071710-01 0.5331820-05	0.5100000-01 0.2655760 00 -0.2932130 00 0.1066390-01 0.9479610-05	0.5979820 02 0.9963120 00 0.9698860 01 0.1924230-01 0.3964650-04	0.6222780 02 0.9736050 00 0.2202870 02 0.1924230-01 0.4200070-04	0.5901420 02 0.5609920 01 0.4867890 03 0.4659110 00 0.4442130-05	0.6058220 02 -0.3043240-02 0.4846380 03 0.4400000-01 0.4442130-05	0.6164460 02 0.5612970 01 0.4846250 03 0.9167000-01 0.1000000-02	0 0 0 1 0
NOZZLE HAS CHOKED								
0.5300000-01 0.1112150 03 0.2792220-01 0.4846050 03 Q 0.4428960-04	0.5250000-01 0.1109150 03 0.5584960 01 0.1061010-01 0.5331820-05	0.5200000-01 0.2657420 00 -0.2902720 00 0.1055640-01 0.9479610-05	0.5840320 02 0.1005680 01 0.9697670 01 0.1918850-01 0.3964650-04	0.6160940 02 0.1000000 01 0.2202500 02 0.1924040-01 0.4200070-04	0.5836440 02 0.5612080 01 0.4868110 03 0.4659110 00 0.4442130-05	0.5875280 02 -0.2960900-02 0.4846050 03 0.4400000-01 0.4442130-05	0.6101880 02 0.5612890 01 0.4846050 03 0.9167000-01 0.1000000-02	0 0 18 1 0
0.5400000-01 0.1112150 03 0.2538520-01 0.4846050 03 Q 0.4428960-04	0.5350000-01 0.1109490 03 0.5586660 01 0.1050290-01 0.5331820-05	0.5300000-01 0.2657420 00 -0.2873240 00 0.1044940-01 0.9479610-05	0.5785610 02 0.1014590 01 0.9697670 01 0.1919430-01 0.3964650-04	0.6098940 02 0.1000000 01 0.2202600 02 0.1924040-01 0.4200070-04	0.5775650 02 0.5612080 01 0.4868330 03 0.4659110 00 0.4442130-05	0.5875280 02 -0.2123610-02 0.4846050 03 0.4400000-01 0.4442130-05	0.6040030 02 0.5612050 01 0.4846050 03 0.9167000-01 0.1000000-02	0 0 0 1 0
0.5500000-01 0.1112150 03 0.2539530-01 0.4846050 03 Q 0.4428960-04	0.5450000-01 0.1109200 03 0.5585210 01 0.1039670-01 0.5331820-05	0.5400000-01 0.2657420 00 -0.2844070 00 0.1034410-01 0.9479610-05	0.5734650 02 0.1022430 01 0.9697670 01 0.1918940-01 0.3964650-04	0.6037590 02 0.1000000 01 0.2202600 02 0.1924040-01 0.4200070-04	0.5719020 02 0.5612080 01 0.4868560 03 0.4659110 00 0.4442130-05	0.5875280 02 -0.6835320-03 0.4846050 03 0.4400000-01 0.4442130-05	0.5979190 02 0.5610610 01 0.4846050 03 0.9167000-01 0.1000000-02	0 0 0 1 0

[illegible]

REVERTING TO EXACT SUPERSONIC SOLUTION

T PCT0 MOF ICT0 E(1)	TSTR PE0 MOE RP E(2)	T1 MCT MUPE RPT E(3)	PT MU MOTS0 RE0 E(4)	PP MN MDCT0 RCT0 E(5)	PD MDCT TP AE E(6)	PN MDPT TPT APE E(7)	PPT MDO TE0 AF DT	NU NM NCT ITER J16
0.6400000-01	0.6350000-01	0.6300000-01	0.5466440 02	0.5627520 02	0.5421010 02	0.5875280 02	0.5617450 02	11
0.1112150 03	0.1078460 03	0.2657420 00	0.1067500 01	0.1000000 01	0.5612080 01	0.2049460-01	0.5591580 01	0
0.1611310 00	0.5430430 01	-0.2676060 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9735710-02	0.9718290-02	0.1865760-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	3
0.0	0.4265060-06	0.8545420-06	0.3978110-05	0.3550550-05	0.0	0.2238810-05	0.1000000-02	-1
0.6500000-01	0.6450000-01	0.6400000-01	0.5454290 02	0.5608010 02	0.5407510 02	0.5875280 02	0.5598570 02	12
0.1112150 03	0.1077580 03	0.2657420 00	0.1069540 01	0.1000000 01	0.5612080 01	0.2173480-01	0.5590340 01	0
0.1643570 00	0.5425980 01	-0.2666790 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9701960-02	0.9685630-02	0.1864230-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	4
0.0	0.1672130-06	0.3349890-06	0.1565000-05	0.1396420-05	0.0	0.8438320-06	0.1000000-02	-1
0.6600000-01	0.6550000-01	0.6500000-01	0.5443010 02	0.5589730 02	0.5394980 02	0.5875280 02	0.5580900 02	8
0.1112150 03	0.1076730 03	0.2657420 00	0.1071450 01	0.1000000 01	0.5612080 01	0.2291740-01	0.5589160 01	0
0.1674650 00	0.5421710 01	-0.2658100 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9670340-02	0.9655060-02	0.1862760-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	3
0.0	0.1712640-06	0.3430700-06	0.1588540-05	0.1417070-05	0.0	0.9403770-06	0.1000000-02	-1
0.6700000-01	0.6650000-01	0.6600000-01	0.5432540 02	0.5572640 02	0.5383350 02	0.5875280 02	0.5564380 02	9
0.1112150 03	0.1075910 03	0.2657420 00	0.1073220 01	0.1000000 01	0.5612080 01	0.2403710-01	0.5588040 01	0
0.1704280 00	0.5417610 01	-0.2649970 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9640770-02	0.9626480-02	0.1861350-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	3
0.0	0.1787030-06	0.3579370-06	0.1651500-05	0.1472890-05	0.0	0.1012350-05	0.1000000-02	-1
0.6800000-01	0.6750000-01	0.6700000-01	0.5422850 02	0.5556660 02	0.5372580 02	0.5875280 02	0.5548950 02	9
0.1112150 03	0.1075140 03	0.2657420 00	0.1074850 01	0.1000000 01	0.5612080 01	0.2510230-01	0.5586980 01	0
0.1732740 00	0.5413710 01	-0.2642370 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9613130-02	0.9599790-02	0.1860010-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	2
0.0	0.2382980-06	0.4772590-06	0.2235760-05	0.1993530-05	0.0	0.1132110-05	0.1000000-02	-1
0.6900000-01	0.6850000-01	0.6800000-01	0.5413870 02	0.5541750 02	0.5362600 02	0.5875280 02	0.5534560 02	11
0.1112150 03	0.1074400 03	0.2657420 00	0.1070380 01	0.1000000 01	0.5612080 01	0.2610840-01	0.5585970 01	0
0.1759790 00	0.5409990 01	-0.2635280 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9587340-02	0.9574890-02	0.1858730-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	2
0.0	0.7187900-08	0.1439450-07	0.4101090-07	0.3655990-07	0.0	0.2058240-06	0.1000000-02	-1
0.7000000-01	0.6950000-01	0.6900000-01	0.5405550 02	0.5527850 02	0.5353350 02	0.5875280 02	0.5521140 02	5
0.1112150 03	0.1073700 03	0.2657420 00	0.1077780 01	0.1000000 01	0.5612080 01	0.2705730-01	0.5585020 01	0
0.1785470 00	0.5406670 01	-0.2628670 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9553290-02	0.9551680-02	0.1857520-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	2
0.0	0.1750020-06	0.3504300-06	0.1621850-05	0.1445570-05	0.0	0.16484120-06	0.1000000-02	-1
0.7100000-01	0.7050000-01	0.7000000-01	0.5397850 02	0.5514900 02	0.5344800 02	0.5875280 02	0.5508650 02	11
0.1112150 03	0.1073040 03	0.2657420 00	0.1079090 01	0.1000000 01	0.5612080 01	0.2795090-01	0.5584130 01	0
0.1809780 00	0.5403140 01	-0.2622510 00	0.9697670 01	0.2202600 02	0.4846050 03	0.4846050 03	0.4846050 03	0
0.4846050 03	0.9540880-02	0.9530070-02	0.1850380-01	0.1924040-01	0.4659110 00	0.4400000-01	0.9167000-01	2
0.0	0.2260830-06	0.4526780-06	0.2071580-05	0.1846100-05	0.0	0.1370530-05	0.1000000-02	-1

TABLE 8
LISTING OF THE COMPUTER PROGRAM HIRTSM1

MAIN

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C. HIRTSMI - HIRT STARTING MODEL

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IMPLICIT REAL*8 (A-H,M,O-Z,$)
COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),E(7),B(30),
1 TVF(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDCTC,FA(3) ,MDTSTR,INFIN,TMGOGS,GP02GS,
1 SGOR
COMMON G,GM1,GP1,OGG,GP102,GM102,GP10G,GM10G,GOGP1,GOGM1,
1 OGGM1,OGGP1,GPOGM1,SGM102,TGGM1,TGG,MGPGM2,TGGP1,GPGM12,
2 MGPOGM,MSOGM1,R,GR,00R,PI,PERR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TSH,TSW,TSP,ISA,ISWA,ISV,CTD,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DT02,TSTOP,00NT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN , PP , PPT , PD , PT , PCT0 , PE0 , MDE ,
- MDD , MDE , MDPT , MDCT , MDPE , MDT0 , MDCT0 , TE0 , TP ,
- TPT , TCT0 , RP , RPT , RE0 , RCT0 , ACT0 , MCT , AE ,
- APE , AF , MN , MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDD1, MDE1, MDPT1, MDCT1, MDPE1,MDT01,MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDD2, MDE2, MDPT2, MDCT2, MDPE2,MDT02,MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDD3, MDE3, MDPT3, MDCT3, MDPE3,MDT03,MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PSOP0,TSOT0,RSOR0,MSOM0
COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$E2,$EMAX,$EP,$DE
COMMON INSTR(26),IDERUG,TIN,IOU,NP,IP,ITER,NVT(3),I,NT,IPAGE,
1 NPAGE,$N,II(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2 IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
DIMENSION V(30,4),RSTR(582),ISTR(35),IEXTP(7),JV(26)
DIMENSION C(4,2)
EQUIVALENCE (RSTR(1),AREA(1)),(ISTR(1),NP),(PN,V(1)),(JV(1),J1)
INTEGER $N
REAL*8 INFIN,KF,KW
DATA IEXTP/1,2,6,7,10,14,16/
DATA C/-,14885381D2,.45347501D2,-,43530673D2,.14068554D2,
1 -,55230057D2,.14939843D3,-,13208170D3,.38913333D2/
DEFINE FILE 01(300,1200,U,J16)
POP0(D1)=(1+GM102*D1*2)**MGOGM1
MDOTPT(D1,D2)=-AWOKW*(D1-D2*A15)*A2
ITIME=0
IFLG1=1
IFLG2=1
IFLG6=1
IFLG9=1
IFLG10=1
IFLG11=1
IFLG12=1
IDERUG=03

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      IIN=05
      IOUT=06
      NP=1
      IS=0
10  READ(IIN,120)NCTL
C-----
C  MANUAL PROGRAM CONTROL
C-----
      IF(NCTL.EQ.0)GO TO 100
      WRITE(IOUT,15)NCTL
15  FORMAT('0NCTL=',I3)
C  1  2  3  4  5  6  7  8  9 10 11 12 13 14
      GO TO(100,125,129,1116,20,30,40,50,60,70,80,90,151,95),NCTL
20  CALL INPUT(&I0)
30  CALL INIT(&I0)
40  CALL CONST(&I0)
50  CALL DUMP(&I0)
60  CALL SOLVER(&I0)
70  CALL PRINT(&I0)
80  CALL BINOM(&I0)
90  CALL RFVERT(&I0)
95  CALL SMPERT(&I0)
C-----
C  READ AND DEFINE DEFAULTED RUN CONTROL INSTRUCTIONS
C-----
100 READ(IIN,120)INSTR
120 FORMAT(26I3)
      IF(INSTR(1).NE.0)IDBUG=INSTR(1)
      IF(INSTR(1).EQ.0)INSTR(1)=IDBUG
      IF(INSTR(2).NE.0)IIN=INSTR(2)
      IF(INSTR(2).EQ.0)INSTR(2)=IIN
      IF(INSTR(3).NE.0)IOUT=INSTR(3)
      IF(INSTR(3).EQ.0)INSTR(3)=IOUT
      IF(INSTR(4).NE.0)NP=INSTR(4)
      IF(INSTR(4).EQ.0)INSTR(4)=1
      IF(INSTR(5).EQ.0)INSTR(5)=2
      IF(IDBUG.EQ.IOUT)INSTR(6)=0
      IF(INSTR(7).EQ.0)INSTR(7)=2
      IF(INSTR(8).EQ.0)INSTR(8)=1
      IF(INSTR(9).EQ.0)INSTR(9)=1
      IF(INSTR(12).EQ.0)INSTR(12)=1
      IF(INSTR(14).EQ.0)INSTR(14)=1
      IF(INSTR(15).EQ.0)INSTR(15)=03
      IF(INSTR(18).EQ.0)INSTR(18)=03
      IF(INSTR(22).EQ.0)INSTR(22)=9999999
      IF(INSTR(25).EQ.0)INSTR(25)=2
      IF(INSTR(26).EQ.0)INSTR(26)=9999999
      DO 121 I=1,26
121  JV(I)=INSTR(I)
      IF(NCTL.NE.0)GO TO 10
C-----
C  PRINT HEADING
C-----
125 WRITE(IOUT,130)
130 FORMAT('1',29X,74('S')/30X,'S',72X,'S'/30X,'S' HIRTSM1 = MATHEMATI
      CAL STARTING MODEL FOR A LUDWIG TURE WIND TUNNEL 'S'/30X,'S',

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MAIN

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272X,'S'/30X,'S' ARNOLD RESEARCH ORGANIZATION, ARNOLD AIR FORCE STA
TION, TN, 12X,'S'/30X,'S', 72X,'S'/30X, 74('S')
IF(NCTL.NE.0)GO TO 10
129 IF(J15.EQ.03)GO TO 135
C-----
C READ SOLUTION FROM DATA FILE AND PRINT
C-----
      K1=0
131 IF(J15.NE.07)GO TO 128
      READ(J15,END=132)RSTR,ISTR
      J16=J16+1
      GO TO 127
128 READ(J15,J16)RSTR,ISTR
      FIND(J15,J16)
127 IF(K1.EQ.0)CALL DUMP
      IPAGE=K1
      CALL PRINT
      K1=IPAGE
      IF(J16-1.EQ.J17)GO TO 132
      GO TO 131
132 IF(INSTR(6).EQ.0)GO TO 134
      WRITE(10UT,133)
133 FORMAT('1')
      CALL DUMP
      IPAGE=0
      CALL PRINT
134 J16=J19
      WRITE(10UT,136)
136 FORMAT('***')
      DMD1=MD-MD1
      DYN1=MN-MV1
      IF(NCTL.NE.0)GO TO 10
      GO TO 1115
C-----
C READ INPUT, INITIALIZE VARIABLES, AND PRINT RESULTS
C-----
135 J16=J19
      IFLG4=0
      CALL INPUT
      CALL CONST
      CALL INIT
      CALL DUMP
      IF(A16.GT.0.00)GO TO 140
      IFLG11=1
      IF(A16.GE.(-1.00))GO TO 139
      A15A=DARS(A15)
      A16A=DARS(A16)
      IFLG11=2
139 A15=1.00
      A16=1.00
140 CALL PRINT
      IF(J18.EQ.03)GO TO 150
      IF(J18.NE.07)GO TO 145
      WRITE(J18,141)
141 FORMAT(32(' '), 'SHOPE = VKF/ADP', 33(' '))
      WRITE(J18)RSTR,ISTR

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J16=J16+1		
GO TO 150		
145 WRITE(J18,J16)RSTR,ISTR		
C-----		
C START NEW TIME INTERVAL		
C-----		
150 T1=T		
T=I+DT		
IF (IFLG11.EQ.0)GO TO 280		
GO TO(269,282,276),IFLG11		
269 IF(MD1.GE.1.00)GO TO 270		
A15=1.00		
A16=1.00		
GO TO 276		
270 A(1)=0.00		
A(2)=0.00		
DO 274 I=1,4		
A(3)=MD1** (I-1)		
DO 272 J=1,2		
272 A(J)=A(J)+C(I,J)*A(3)		
274 CONTINUE		
A15=A(1)		
A16=A(2)		
GO TO 276		
282 IF(MD1.GE.1.00)GO TO 284		
A15=(A15A-1.00)*MD1+1.00		
A16=(A16A-1.00)*MD1+1.00		
GO TO 276		
284 A15=A15A		
A16=A16A		
IFLG11=3		
276 WRITE(IDEBUG,278)MD1,A15,A16		
278 FORMAT(' MD1=',F16.8,' A15=',E16.8,' A16=',E16.8)		
280 IF (IFLG2.LT.0) I5=I5+1		
IF (I5.NE.INSTR(26))GO TO 143		
INSTR(23)=1		
WRITE(IOUT,142)		
142 FORMAT('0 REVERTING TO EXACT SUPERSONIC SOLUTION')		
143 IF ((INSTR(10).EQ.0).OR.(ITER.LT.INSTR(11)))GO TO 152		
C WEIGHT CUTTING		
A11=.5*A11		
A12=1.-A11		
INSTR(11)=INSTR(11)+INSTR(20)		
WRITE(IOUT,1205)ITER,A11,INSTR(11)		
1205 FORMAT('0',5X,'ITER=',I3,' WT HALVED TO ',F5.3,' INSTR(11) RAISED		
10 TO ',I7)		
IPAGE=IPAGE+2		
152 IF (T.GE.A10)IDEBUG=IOUT		
TSTR=T1+DT02		
IIME=IIME+1		
IF (IFLG1.EQ.2)IFLG4=1		
C SET PRESSURES OF LAST ITERATION TO INFINITY FOR ERROR COMPUTATION		
DO 153 I=1,7		
153 V(I,3)=INFIN		
ITER=0		
IF (T.LF.TSTOP)GO TO(155,240),IFLG1		

4

	MAIN	DATE = 75157	11/58/40
	MCT=\$X1		
	MCT=\$N		
	A(1)=(1.+GM102*MCT**2)/(1.+GM102*MCT)**2		
	TCT0=TC*A(1)		
	TE0=TCT0		
	PCT0=PC*A(1)**GOGM1		
	RCT0=PCT0*00R/TCT0*A2		
	ACT0=DSQRT(GR*TCT0)		
	MDTS0=RCT0*ACT0*TSA		
	MDCT0=RCT0*ACT0*CTA		
	MN=1.		
	PN=PSOP0*PCT0		
	MDCT=RSOR0*RCT0*DSQRT(TS0T0)*ACT0*TSA		
	IFLG2=-1		
	WRITE(TOUT,249)		
249	FORMAT('0NOZZLE HAS CHOKED')		
	MCT1=MCT		
	TCT01=TCT0		
	PCT01=PCT0		
	RCT01=RCT0		
	ACT01=ACT0		
	MN1=MN		
	PN1=PN		
	MDCT1=MDCT		
	MDTS01=MDTS0		
	MDCT01=MDCT0		
	PT=A17*PD+(1.D0-A17)*PN		
	PT1=PT		
250	IF(INSTR(23).EQ.0)GO TO 255		
	IF(I.GT.A13)INSTR(23)=2		
	IF((INSTR(23).EQ.3).AND.(ITER.GE.INSTR(11)))GO TO 253		
252	IF(ITER.NE.1)GO TO 255		
C	-----		
C	CALL SMALL PERTURBATION PACKAGE		
C	-----		
253	DO 260 I=1,3		
	I1=I+25		
	V(I1,1)=ARFATS(I)		
	EA(I)= V(I1,1)-V(I1,2)		
260	CONTINUE		
	IF((ARFATS(3).EQ.0.D0).AND.(EA(3).EQ.0.D0))MDF1=0.D0		
	IF((ARFATS(2).EQ.0.D0).AND.(EA(2).EQ.0.D0))MDPE1=0.D0		
	CALL SMPERT		
	K1=DSIGN(1.5D0,PT-PCT0*PSOP0)		
	IF(K1.EQ.IFLG2)GO TO 256		
	IF(IFLG10.EQ.2)GO TO 258		
	IFLG10=2		
	IF(IFLG2.EQ.1)IFLG2=K1		
	IF(IFLG2.EQ.(-1))GO TO 241		
	GO TO 256		
258	IFLG10=1		
256	IF((INSTR(23).EQ.3).AND.(ITER.GE.INSTR(11)))GO TO 254		
	IF(INSTR(23).EQ.2)GO TO 254		
	GO TO 255		
254	IF(J18.EQ.03)GO TO 1190		
	WRITE(J18,J16)RSTR,ISTR		

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                                MAIN                                DATE = 75157                                11/58/40
FJND(J18*J16)
GO TO 1190
255 IF(IFLG2)500,9999,251
C-----
C SURSONIC BRANCH
C-----
251 MDE=A1*PE0*AREATS(1)/DSORT(TF0)*A2
IFLG2=-1
IFLG6=1
IFLG12=1
MDD=MDF+MDF
C DIFFUSER MACH NUMBER AND PRESSURE
SY=MDD/MOTS0
IFLG=2
CALL SOLVER
MD=$X1
ND=$N
PD=PCT0*POP0(MD)
PT=.5*(PD+PN)
MDPT=MDDTPT(PD,PT)
MDCT=MDD+MDPT
C CHARGE TUBE MACH NUMBER
SY=MDCT/MOTS0
IFLG=4
CALL SOLVER
MCT=$X1
NCT=$N
A(1)=(1.+GM102*MCT**2)/(1.+GM102*MCT)**2
TCT0=TC*A(1)
PCT0=PC*A(1)**GOGM1
RCT0=PCT0*QOR/TCT0*A2
PE0=PCT0
TEN=TCT0
REN=RCT0
ACT0=DSORT(GR*TCT0)
A(1)=RCT0*ACT0
MDCT0=A(1)*CTA
MOTS0=A(1)*ISA
C NOZZLE MACH NUMBER AND PRESSURE
SY=MDCT/MOTS0
IFLG=2
CALL SOLVER
MN=$X1
NN=$N
PN=PCT0*POP0(MN)
IF(PT.LE.PCT0*PSOP0)60 TO 241
GO TO 1000
C-----
C SJPERSONIC BRANCH
C-----
500 PT=A17*PD*(1.70-A17)*PN
IFLG2=-1
IFLG12=1
MDPT=MDDTPT(PD,PT)
MDD=MDDT-MDDPT
MDE=-MDF+MDD

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                                MAIN                                DATE = 75157          11/58/40
TE0=TCT0
PE0=MDE*DSORT(TE0)*00A1/AREATS(1)*A3
RE0=PE0*00R/TE0 *A2
C DIFFUSER PRESSURE AND MACH NUMBER
SY=MDD/MDTS0
IFLG=2
CALL SOLVER
MD=$X1
ND=$N
PD=PCT0*POP0(MD)
C-----
C UPDATE PLENUM CONDITIONS
C-----
1000 IF(IFLG2)1001,9999,1002
1001 PT=(1.00-A17)*PN+A17*PD
GO TO 1003
1002 PT=0.500*(PN+PD)
1003 MDPT=MDDTPT(PD,PT)
MDF=-AREATS(3)*00KF*(PP-PD*A16)*A2
RPT=RPT1+(MDPT+MDF+MDPE)*DTOPV
RP=.5*(RPT1+RPT)
IF(IFLG9.EQ.2)GO TO 1010
PPT=PPT1*(RPT/RPT1)**6
TPT=PPT*00R/RPT*A2
PP=PPT1*(RP/RPT1)**6
TP=PP*00R/RP *A2
IF(INSTR(25).EQ.1)GO TO 1020
IF(TPT.GE.TCT0)GO TO 1020
IFLG9=2
1010 TP=TCT0
TPT=TCT0
PP=3P*R*TP/A2
PPT=RPT*RP/TPT/A2
1020 MDPE=-A1*PP*AREATS(2)/DSORT(TP)*A2
C-----
C CONVERGENCE CHECK
C-----
IFLG3=1
DO 1050 I=1,7
1050 E(I)=2.*DABS(V(I,1)-V(I,3))/(V(I,1)+V(I,3))
DO 1100 I=1,7
IF(E(I).GT.PERR)GO TO 1200
1100 CONTINUE
C WRITE DATA ON FILE AND PRINT CONVERGED DATA
IFLG3=2
IF(J18.EQ.03)GO TO 1115
WRITE(J18*J16)RSTR,ISTR
FIND(J18*J16)
1115 CALL PRINT
1116 IF(INSTR(7).EQ.2)GO TO 1180
IF(INSTR(6).NE.0)WRITE(1OUT,1120)V
1120 FORMAT(15(/' *RE16.8))
C-----
C PERFORM EXTRAPOLATION TO NEXT TIME INTERVAL
C-----
DO 1170 I=1,7

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                                MAIN                                DATE = 75157                                11/58/40
J=IEXTP(I)
C SAVE DATA FOR CURRENT INTERVAL
A(I)=V(J,1)
IF(I TIME.EQ.1)GO TO 1160
IF(IFLG4.EQ.1)GO TO 1160
C EXTRAPOLATE
V(J,1)=2.*V(J,1)-V(J,2)
IF(INSTR(6).NE.0)CALL PRINT
C RESET DATA TO BEGINNING OF TIME INTERVAL
1160 V(J,2)=A(I)
IF(IFLG4.EQ.1)IFLG4=2
1170 CONTINUE
C-----
C DETERMINE IF ERROR CUTTING OR DT DOUBLING IS REQUIRED
C-----
1180 IF((INSTR(12).EQ.1).AND.(INSTR(14).EQ.1))GO TO 1190
DO 1185 I=1,7
IF((IFLG2.NE.1).AND.((I.EQ.1).OR.(I.EQ.6)))GO TO 1185
E(I)=2.*DABS(V(I,1)-V(I,2))/(V(I,1)+V(I,2))/INSTR(12)
IF(E(I).GT.PERR)GO TO 1185
IF(INSTR(14).GT.1)GO TO 1184
C ERROR CUTTING
PERR=PERR/INSTR(12)
$EMAX=$EMAX/INSTR(12)
WRITE(10UT,1183)PERR,$EMAX
1183 FORMAT('0 PERR CUT TO',F16.8,' AND $EMAX CUT TO',F16.8)
IPAGE=IPAGE+2
GO TO 1190
C DT DOUBLING
1184 DT=DT*INSTR(14)
DTPV=DT/PV
ODDT=1./DT
DT02=DT*.5
WRITE(10UT,1187)DT
1187 FORMAT('0 DT RAISED TO',F16.8)
IPAGE=IPAGE+2
GO TO 1190
1185 CONTINUE
1190 IF(I TIME.LE.2)GO TO 1191
C-----
C DETERMINE IF NEXT INTERVAL IS PREDICTED TO CHOKE
C-----
DMD=MD-MD1
DMN=MN-MN1
IF(DMD.LT.DMD1)DMD=DMD1
IF(DMN.LT.DMN1)DMN=DMN1
IFLG2=DSIGN(1.5D0,PT-PCT0*PSOP0)
IF((MD.GE.(1.D0-DMD)).OR.(MN.GE.(1.D0-DMN)))IFLG2=-1
1191 DMD1=MD-MD1
DMN1=MN-MN1
IF(IFLG2.LT.0)IFLG10=2
WRITE(1DEBUG,2000)IFLG2,DMD,DMN,MD,MN,MD1,MN1
2000 FORMAT('0IFLG2=',I3,' DMD,DMN=',2E13.5,' MD,MN=',2E13.5,
1 ' MD1,MN1=',2E13.5)
C RESET DATA TO BEGINNING OF TIME INTERVAL
IF(INSTR(23).EQ.0)GO TO 1189

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	MAIN	DATE = 75157	11/58/40
	DO 118A I=1,30		
118A	V(I,2)=V(I,1)		
	GO TO 150		
1189	RPI1=RPT		
	PPT1=PPT		
	IF(INSTR(7).EQ.1)GO TO 150		
	DO 1186 I=1,7		
	J=IFXTP(I)		
1186	V(J,2)=V(J,1)		
	GO TO 150		
C	-----		
C	RESET CONVERGENCE CONTROL DATA		
C	-----		
1200	IF(ITERUG.EQ.1000)CALL PRINT		
	IF(INSTR(6).NE.0)CALL PRINT		
	I1=INSTR(10)		
1210	DO 1260 I=1,30		
	GO TO(1220,1240),I1		
	V(I,3)=V(I,1)		
	GO TO 1260		
1220	IF(ITER.NE.1) V(I,1)=A11*V(I,1)+A12 *V(I,3)		
	V(I,3)=V(I,1)		
	GO TO 1260		
1240	V(I,4)=V(I,3)		
	V(I,3)=V(I,1)		
	IF(ITER.NE.1)V(I,1)=A11*V(I,1)+A12*V(I,4)		
1260	CONTINUE		
	GO TO 240		
	END		

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                                INPUT                                DATE = 75157                11/58/40
SUBROUTINE INPUT(*)
IMPLICIT REAL*8 (A-H,M,O-Z,$)
COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),E(7),B(30),
1 TVF(3),IDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDCTC,EAE,EAPE,EAF,MDTSTR,INFIN,TMGOGS,GP02GS,
1 SGOR
COMMON G,GM1,GP1,00G,GP102,GM102,GP10G,GM10G,G0GP1,G0GM1,
1 00GM1,00GP1,GPOGM1,SGM102,T0GM1,T0G,MGP0GM2,T0GP1,GPGM12,
2 MGPOGM,MGOGM1,R,GR,00R,PI,PERR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TSH,TSW,TSP,TSa,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DTU2,TSTOP,00DT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MDD, MDF, MDPT, MDCT, MDPE, MDT0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDD1, MDF1, MDPT1, MDCT1, MDPE1,MDT01,MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDD2, MDF2, MDPT2, MDCT2, MDPE2,MDT02,MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDD3, MDF3, MDPT3, MDCT3, MDPE3,MDT03,MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PSOP0,TSOT0,RSOR0,MSOMO
COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$E2,$EMAX,$EP,$OF
COMMON INSTR(26),IDERUG,TIN,IOUT,NP,IP,ITER,NVT(3),I,NT,IPAGE,
1 NPA3E,$N,$T(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2 IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
DIMENSION V(30,4),RSTR(579),ISTR(35)
EQUIVALENCE (V(1),PN),(RSTR(1),AREA(1)),(ISTR(1),NP)
INTEGER $V
REAL*8 INFIN,KF,KW
IF(NCTL.NE.0)GO TO 200
READ(IIN,50)NVT,NT
50 FORMAT(26I3)
READ(IIN,100)PC,TC
100 FORMAT(5E16.8)
READ(IIN,100)TSL,TSH,TSW,CTD,PVOTSV,TAUW,KW,KF,A15,A16,
1 A17,A18,A19,A20,A21
READ(IIN,100)R,G,A11,A13,A14
READ(IIN,100)DT,TSTOP,$EMAX,PERR,A10
READ(IIN,100)AREAM
READ(IIN,100)TVF
READ(IIN,100)IDELAY
DO 110 J=1,3
DO 105 I=1,50
TV(J,I)=0.
105 AREA(J,I)=0.

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	INPUT	DATE = 75157	11/58/40
	I1=VVT(J)		
110	READ(IYN,120)(TV(J,I),AREA(J,I),I=1,I1)		
120	FORMAT(2E16.8)		
	RETURN		
200	READ(IYN,50)I1,I2,I3		
	IF(I1.EQ.0)RETURN 1		
	GO TO(220,240,260),I1		
220	READ(IYN,50)ISTR(I2)		
	GO TO 200		
240	READ(IYN,100)RSTR(I2)		
	GO TO 200		
260	READ(IYN,100)V(I2,I3)		
	GO TO 200		
	END		

CONST

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SURROUTINE CONST(*)
IMPLICIT REAL*8 (A-H,M,O-Z,$)
COMMON ARFA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),E(7),R(30),
1 TVF(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDCIC,FAF,FAPE,FAF,MDTSTR,INFIN,TM30GS,GP02GS,
1 SGOR
COMMON G,GM1,GP1,00G,GP102,GM102,GP10G,GM10G,GP0GP1,GP0GM1,
1 00GM1,00GP1,GP0GM1,SGM102,T0GM1,T0G,MGP0GM2,T0GP1,GP0GM12,
2 MGP0GM,M30GM1,R,GR,00H,PI,PERR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TSL,TST,TSP,TSI,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DT02,TSTOP,00DT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PF0, MDF,
- MD0, MDF, MDPT, MDCT, MDPE, MDTS0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PF01, MDF1,
- MD01, MDF1, MDPT1, MDCT1, MDPE1, MDTS01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PF02, MDF2,
- MD02, MDF2, MDPT2, MDCT2, MDPE2, MDTS02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PF03, MDF3,
- MD03, MDF3, MDPT3, MDCT3, MDPE3, MDTS03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PS0P0,TS0T0,PS0R0,MS0M0
COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$F2,$EMAX,$EP,$OF
COMMON INSTR(26),IDFHUG,TIN,TOUT,NP,JP,ITER,NVT(3),I,VT,IPAGE,
1 NPAGE,$N,IT(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2 IFL55,IFL56,IFL67,IFL68,IFL69,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
INTEGER $N
REAL*8 INFIN,KF,KW
PI=3.141592653589793
GM1=G-1
GP1=G+1
GM102=.5*GM1
GP102=GP1*.5
00G=1./G
GM10G=GM1*00G
GP10G=GP1*00G
GP02GS=.5*GP10G*00G
GP0GM1=G/GM1
GP0GP1=G/GP1
SGM102=DSRT(GM102)
T0G=2.*00G
T0GP1=2./GP1
00GM1=1./GM1
GP0GM1=GP1*00GM1
GP0GP1=.5*GP0GM1
00GP1=1./GP1

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DOR=1./R
GR=G*R
MGP0GM=-GPOGM1
MGP0GM2=-GPOGM12
MGO0GM1=-G0GM1
TMG0GS=(2.-G)*006**2
T0GM1=2./GM1
IF (INSTR(5).EQ.1) GO TO 100
A2=144.
A3=1./A2
GO TO 200
100 A2=1.
A3=A2
200 CONTINUE
C SERIES FOR UNSTEADY MASS FLUX FROM MACH NUMBER
A(R)=MGP0GM
A(9)=GM102
CALL BINOM
CALL REVERT
DO 220 I=1,7
220 RW(I)=A(I)
SGOR=DSORT(G*DOR)
IF (NCTI.EQ.7) RETURN 1
RTURN
END

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TNIT

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SURROUTINE INIT(*)
IMPLICIT REAL*8 (A-H,M,O-Z,*)
COMMON AREA(3,50),AREATS(3),ARFAM(3),TV(3,50),A(10),E(7),R(30),
1 TVF(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDC TC,FAF,FAPE,FAF,MNTSTR,INFIN,TM30GS,GPO2GS,
1 SGOR
COMMON G,GM1,GP1,GOG,GP102,GM102,GP10G,GM10G,GOGP1,GOGM1,
1 OOGM1,OGSP1,GPOGM1,SGM102,T0GM1,T0G,MGPGM2,T0GP1,GPGM12,
2 MGPOGM,M30GM1,P,GR,00R,PI,PFRR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TS4,TSW,TSP,TSa,TSWA,TSV,CT0,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DT02,TSTOP,00NT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,AR,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MDO, MDE, MDE, MDE, MDE, MDE, MDE, MDE, MDE, MDE,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDO1, MDE1, MDE1, MDE1, MDE1, MDE1, MDE1, MDE1, MDE1, MDE1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDO2, MDE2, MDE2, MDE2, MDE2, MDE2, MDE2, MDE2, MDE2, MDE2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDO3, MDE3, MDE3, MDE3, MDE3, MDE3, MDE3, MDE3, MDE3, MDE3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PSOP0,TSOT0,RSOR0,MSOM0
COMMON $Y,$Y1,$Y2,$X1,$X2,$OX,$E1,$E2,$EMAX,$EP,$DE
COMMON INSTR(26),TDERUG,TIN,TOUT,NP,IP,ITER,NVT(3),I,NT,IPAGE,
INPAGE,$V,IT(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
DIMENSION V(30,4)
EQUIVALENCE (PN,V(1,1))
INTEGER $V
REAL*8 INFIN,AF,KW
DO 5 I=1,3
IT(I)=2
5 AREATS(I)=0.
TSA=TSW*TSW
TSP=2.*(TSW*TSW)
TSWA=TSL*TSW
CTA=PI*CT0**2*.25
TSV=TSa*TSL
PV=TSV*PVOTSV
DTOPV=DT/PV
RC=PC*00R/TC*A2
AC=JSGRT(AR*TC)
MDC TC=RC*AC*CTA
MDSO=RC*AC*TSA
DO 50 J=1,4
DO 10 I=1,7

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10 V(I,J)=PC
   DO 20 I=8,13
20 V(I,J)=0.
   V(14,J)=MDTS0
   V(15,J)=MDCTC
   DO 30 I=16,19
30 V(I,J)=TC
   DO 40 I=20,23
40 V(I,J)=RC
   V(24,J)=AC
   DO 45 I=25,30
45 V(I,J)=0.
50 CONTINUE
   T=0.
   T1=0.
   MD=0.
   TSTR=0
   DT02=.5*DT
   QDOT=1./DT
   MN=0.
   MCT=0.
   DO 160 J=1,3
     I1=NVT(J)
     DO 140 I=1,I1
       TV(J,I)=TV(J,I)*TVF(J)
140 AREA(J,I)=AREA(J,I)*AREAM(J)
       IF(TDELAY(J).EQ.0.)GO TO 160
     DO 150 I1=1,49
       I=51-I1
       AREA(J,I)=AREA(J,I-1)
150 TV(J,I)=TV(J,I-1)+TDELAY(J)
       NVT(J)=NVT(J)+1
160 CONTINUE
     DO 170 I=1,3
170 V(I+25,2)=AREA(I,1)
       PSOP0=TGP1**GPGM1
       TSOT0=TGP1
       PSOR0=TGP1**GPGM1
       MSOM0=PSOR0*DSQRT(TSOT0)
       MDTSR=MDTS0*MSOM0
       TINFN=1.E+70
       A1=DSQRT(TGP1**GPGM1*G*QOR)
       QOA1=1./A1
       A4=TGP1**GPGM12
       A5=1./MSOM0
       A6=1.-PSOP0
       A7=2.*GPI/MOCTC
       A8=-GM1OG
       A9=2.*GPI
       IF(A10.EQ.0.)A10=TINFN
       IF(A11.EQ.0.)A11=.5
       IF(A13.EQ.0.)A13=TINFN
       IF(A14.EQ.0.)A14=-0.1
       IF(A14.GT.0.)A14=-A14
       A12=1.-A11
       IF(A15.EQ.0.D0)A15=1.00

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IF (A16.FQ.0.00) A16=1.00
IF (A17.FQ.0.00) A17=1.00
IPAGE=0
NPAGE=50
IP=0
OOKF=1./KF
AWOKW=.17*TAUW*TSWA/KW
ITER=0
IFLG=0
IFLG2=1
IFLG3=0
IFLG4=0
IFLG5=0
IFLG6=0
IFLG7=0
IFLG8=0
IFLG9=0
ND=0
NV=0
NCT=0
I1=0
I2=0
I3=0
I4=0
I5=0
DO 200 I=1,7
200 E(I)=0.
IF (NCTL.EQ.6) RETURN 1
RETURN
END

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SUBROUTINE DUMP (*)
  IMPLICIT REAL*8 (A-H,M,O-Z,%)
  COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),F(7),B(30),
1 TVF(3),TDELAY(3),RW(7)
  COMMON PC,RC,TC,AC,MDCIC,FAF,EAPF,FAF,MDTSTP,INFIN,TM30GS,GP02GS,
1 SGOR
  COMMON G,GM1,GP1,DOG,GP102,GM102,GP10G,GM10G,G0GP1,G0GM1,
1 G0GM1,G0GP1,GP0GM1,SGM102,T0GM1,T0G,M0PGM2,T0GP1,GP0GM12,
2 MG0GM,M0GM1,P,GR,00R,PT,PERR,AWOKW,00A1,00KF,KF,KW
  COMMON TSL,TSH,TSW,TSP,TSA,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
  COMMON T,T1,DT,TSTR,DT02,TSTOP,00DT,DTOPV
  COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
  COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MD0, MDE, MDPT, MDCI, MDPE, MDT0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, PE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
  COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MD01, MDE1, MDPT1, MDCI1, MDPE1, MDT01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, PE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
  COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MD02, MDE2, MDPT2, MDCI2, MDPE2, MDT02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, PE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
  COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MD03, MDE3, MDPT3, MDCI3, MDPE3, MDT03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, PE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
  COMMON PS0P0,TS0T0,PS0R0,MS0M0
  COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$F1,$F2,$FMAX,$EP,$DE
  COMMON INSTR(26),IDFRUG,TIN,TOUT,NP,IP,ITER,NVT(3),I,VT,IPAGE,
1 NPAGE,$V,IT(3),J,TM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2 IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
  COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
  COMMON NCTL
  DIMENSION JV(24),V(30,4)
  EQUIVALENCE (JV(1),J1),(V(1),PN)
  INTEGER $V
  REAL*8 INFIN,KF,KW
  LOGICAL*4 CHAR1(2),CHAR2(2,2)
  DATA CHAR1/'PSFA','PSIA',CHAR2/'SFC0','ND ','SEVE','INTH'/'
  WRITE(TOUT,100)(I,I=1,26),INSTR
100 FORMAT('0',16(' INSTR',I2)/' ',10(' INSTR',I2),2(/' ',16(I8))
  WRITE(TOUT,120)IDFRUG,TIN,TOUT,I1,IPAGE,NPAGE,NP,IP,ITER,I2,I,
1 IFLG,IFLG1,IFLG2,IFLG3,IFLG4,IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,ND,NN,
2 NCT,ITIME,NVT,I3,I4,I5,IT,NCTL
120 FORMAT('0 IDFRUG TIN TOUT I1 IPAGE NPAGE NP IP
1, ITER I2 I IFLG IFLG1 IFLG2 IFLG3 IFLG4 IFLG5
2, IFLG6,/,',18I7/'0 IFLG7 IFLG8 IFLG9 ND NN NCT
3, ITIME NVT(1) NVT(2) NVT(3) I3 I4 I5 IT(1)',
4, IT(2) IT(3) NCTL'
-/' ',18I7)
  WRITE(TOUT,140)TSL,TSH,TSW,CTD,PVOTSV,TAUW,KW,KF,AREAM,TSA,
1 TSP,TSWA,CTA,TSV,PV,AWOKW
140 FORMAT('0',7X,'TSL',13X,'TSH',13X,'TSW',13X,'CTD',12X,'PVOTSV',

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1 11X,'TAUW',13X,'KW',14X,'KF'/'',.8E16.8/'0',7X,'AEM',13X,'APM',
213X,'AFM',13X,'TSA',13X,'ISP',12X,'TSA',13X,'CTA',13X,'TSV'/'
3' ,.8E16.8/'0',7X,'PV',13X,'AWOKW',10X,'',10X,'',
4 10X,'',10X,'',10X,'',10X,'',
5 /' ,.8E16.8)
WRITE(IOUT,180)PC,RC,TC,PP,RP,TP,PP1,RP1,TP1,PPT,RPT,TPT,PCTO,
1 RCTO,TC TO,PEO,REO,TF0,AC,ACTO,MDE,MDCT,MDE,MDF,T,T1,MN,TSTR,
2MDPT,MDD,PD,PV,PT,MD,MCT,PP2,MDE2,MDF2,MDF2,MDCT2,$EMAX,
3TSTOP,DT,PSOP0,TSOT0,RSOR0,MSOM0,MDTSTR,MDTS0,MDCTC,MDCT0,PERR
4,PN3,PP3,MDF3,PE03,TE03,MDTS03,PC TO3,TDELAY
190 FORMAT('0',7X,'PC',14X,'RC',14X,'TC',14X,'PP',14X,'RP',14X,
1'TP',14X,'PP1',13X,'RP1'/'',.8E16.8/'0',7X,'TP1',13X,'PPT',13X,
2'RPT',13X,'TPT',12X,'PC TO',12X,'RCTO',12X,'TC TO',13X,'PEO'/'',
3.8E16.8/'0',7X,'REO',13X,'TF0',13X,'AC',12X,'ACTO',13X,'MDE',11X,
4'MDCT',13X,'MDE',13X,'MDF'/'',.8E16.8/'0',8X,'T',14X,'T1',14X,'MN',
5,13X,'TSTR',12X,'MDPT',13X,'MDD',13X,'PD',14X,'PN'/'',.8E16.8/
6'0',7X,'PT',14X,'MD',14X,'MCT',13X,'PP2',12X,'MDE2',12X,'MDF2',
7_11X,'MDF2',12X,'MDCT2'/'',.8E16.8/'0',6X,'$EMAX',11X,'TSTOP',
811X,'DT',12X,'PSOP0',11X,'TSOT0',11X,'RSOR0',11X,'MSOM0',10X,
9'MDTSTR'/'',.8E16.8/'0',6X,'MDTS0',11X,'MDCTC',11X,'MDCT0',
A 10X,'PERR',10X,'PN3',10X,'PP3',10X,'MDF3',10X,
B 'PE03'/'',.8E16.8/'0',6X,'TE03',11X,'MDTS03',11X,'PC TO3',10X,
C TDELAY',10X,'TDELAY',10X,'TDELAY',10X,'',10X,'',/'',
D.8E16.8)
WRITE(IOUT,200)A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,
1 A15,A17,A18,A19,A20,A21,A22,A23,A24
200 FORMAT('0',7X,'A1',14X,'A2',14X,'A3',14X,'A4',14X,'A5',14X,
1'A6',14X,'A7',14X,'A8'/'',.8E16.8/'0',7X,'A9',14X,'A10',
313X,'A11',13X,'A12',13X,'A13',13X,'A14',13X,'A15',13X,'A16'/'',
4 .8E16.8/'0',7X,'A17',13X,'A18',13X,'A19',13X,'A20',13X,'A21',
5 13X,'A22',13X,'A23',13X,'A24'
-/' ,.8E16.8)
WRITE(IOUT,220)G,GMI,GPI,GM102,GP102,NOG,GM106,GP106,GOGMI,
1GOGPI,SGM102,TOG,TGPI,OGMI,GP0GM1,GP0GM12,TGMI,MGP0GM2,MGP0GM,
200GPI,P,OCR,GR,DT02,DTOPV,OKF,INFIN,OOA1,ODDT,MGOGMI,TMGOGS,
3 GP02GS,S3OR
220 FORMAT('1',8X,'G',14X,'GMI',13X,'GPI',12X,'GM102',11X,'GP102',
1 12X,'NOG',12X,'GM106',11X,'GP106'/'',.8E16.8/'0',6X,'GOGMI',
211X,'GOGPI',10X,'SGM102',12X,'TOG',12X,'TGPI',11X,'OGMI',
310X,'GP0GM1',10X,'GP0GM12'/'',.8E16.8/'0',6X,'TGMI',10X,'MGP0GM2',
4 10X,'MGP0GM',11X,'OGPI',13X,'R',14X,'OCR',13X,'GR',13X,'DT02'/'
5' ,.8E16.8/'0',6X,'DTOPV',11X,'OKF',12X,'INFIN',11X,'OOA1',
6 12X,'ODDT',11X,'MGOGMI',10X,'TMGOGS',10X,'GP02GS'/'',.8E16.8
7 /'0',6X,'S3OR'
- /' ,.8E16.8)
WRITE(IOUT,260)V
260 FORMAT('0V EQUIVALENCE ARRAY',15(' ',.8E16.8))
WRITE(IOUT,240)(I,I=1,7),RW
240 FORMAT('0',7(5X,'RW('',I1,''),5X)/'',.7E16.8)
NINSTR=26
DO 1000 I=1,NINSTR
IF(I.EQ.12)WRITE(IOUT,490)
490 FORMAT('1')
WRITE(IOUT,500)I,INSTR(I)
500 FORMAT('0INSTR('',I2,'')='',I7)
C 1 2 3 4 5 6 7 8 9 10 11 12 13 14 15

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C	15 17 18 19 20 21 22 23 24 25 26 27 28 29 30 31		
	GO TO(510,520,530,540,550,560,570,580,590,600,610,620,630,640,650		
	1,660,670,680,690,700,710,720,730,740,750,760		
	=),I		
510	WRITE(IOUT,511)INSTR(1)		
511	FORMAT('++',20X,'SEND DEBUGGING OUTPUT TO DSRN',I3)		
	GO TO 1000		
520	WRITE(IOUT,521)INSTR(2)		
521	FORMAT('++',20X,'OBTAIN INPUT FROM DSRN',I3)		
	GO TO 1000		
530	WRITE(IOUT,531)INSTR(3)		
531	FORMAT('++',20X,'SEND REGULAR OUTPUT TO DSRN',I3)		
	GO TO 1000		
540	WRITE(IOUT,541)INSTR(4)		
541	FORMAT('++',20X,'PRINTING TIME INTERVAL:',I3)		
	GO TO 1000		
550	WRITE(IOUT,551)CHAR1(INSTR(5))		
551	FORMAT('++',20X,'INPUT AND OUTPUT PRESSURES IN ',A4)		
	GO TO 1000		
560	IF(INSTR(5).EQ.0)WRITE(IOUT,562)		
	IF(INSTR(5).NE.0)WRITE(IOUT,561)		
561	FORMAT('++',20X,'PRINT DATA AFTER EVERY ITERATION')		
562	FORMAT('++',20X,'PRINT DATA ONLY WHEN CONVERGED')		
	GO TO 1000		
570	IF(INSTR(7).EQ.1)WRITE(IOUT,571)		
	IF(INSTR(7).EQ.2)WRITE(IOUT,572)		
571	FORMAT('++',20X,'LINEARLY EXTRAPOLATE TO NEXT TIME INTERVAL AS AN I		
	INITIAL GUESS')		
572	FORMAT('++',20X,'DO NOT EXTRAPOLATE TO NEXT TIME INTERVAL')		
	GO TO 1000		
580	WRITE(IOUT,581)(CHAR2(J,3-INSTR(8)),J=1,2)		
591	FORMAT('++',20X,'USE ',2A4,' DEGREE REVERTED SERIES AS INITIAL GUESS		
	IS TO MASS FLUX-MACH NUMBER WAVE EQUATION')		
	GO TO 1000		
590	IF(INSTR(9).EQ.1)WRITE(IOUT,591)		
	IF(INSTR(9).EQ.2)WRITE(IOUT,592)		
591	FORMAT('++',20X,'USE ITERATIVE SOLUTION TO ENERGY AND WAVE EQUATION		
	IS')		
592	FORMAT('++',20X,'USE APPROXIMATE EXPANSIONS FOR ENERGY AND WAVE EQU		
	ATIONS')		
	GO TO 1000		
600	IF(INSTR(10).EQ.0)WRITE(IOUT,601)		
	IF(INSTR(10).EQ.1)WRITE(IOUT,602)		
	IF(INSTR(10).EQ.2)WRITE(IOUT,603)		
601	FORMAT('++',20X,'DO NOT INVOKE AVERAGING OPTION')		
602	FORMAT('++',20X,'AVERAGE VALUES OF CURRENT ITERATION WITH AVERAGE V		
	VALUES OF PREVIOUS ITERATION')		
603	FORMAT('++',20X,'AVERAGE VALUES OF CURRENT ITERATION WITH UNAVERAGE		
	ED VALUES OF PREVIOUS ITERATION')		
	GO TO 1000		
610	WRITE(IOUT,611)INSTR(11)		
611	FORMAT('++',20X,'CURRENT WEIGHT IS HALVED BEYOND ',I7,' ITERATIONS'		
	1)		
	GO TO 1000		
620	IF(INSTR(12).EQ.1)WRITE(IOUT,621)		
	IF(INSTR(12).NE.1)WRITE(IOUT,622)(INSTR(12),J=1,2)		

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621 FORMAT('++',20X,'DO NOT INVOKE ERROR CUTTING OPTION')
622 FORMAT('++',20X,'DIVIDE ERRORS BY',I3,' WHEN TIME-DIFFERENCES ARE L
    LESS THAN',I3,' TIMES THE ERRORS')
    GO TO 1000
630 IF (INSTR(13).EQ.0) WRITE(IOUT,631)
    IF (INSTR(13).NE.0) WRITE(IOUT,632)
631 FORMAT('++',20X,'PRINT ALL DATA')
632 FORMAT('++',20X,'PRINT ONLY TIMES AND PRESSURES')
    GO TO 1000
640 IF (INSTR(14).EQ.1) WRITE(IOUT,641)
    IF (INSTR(14).GT.1) WRITE(IOUT,642) INSTR(14),INSTR(12)
641 FORMAT('++',20X,'DO NOT INVOKE DT-RAISING OPTION')
642 FORMAT('++',20X,'SET DT =',I3,'*DT IF TIME-DIFFERENCES ARE LESS THA
    N',I3,' TIMES THE ERRORS')
    GO TO 1000
650 IF (INSTR(15).EQ.03) WRITE(IOUT,651)
    IF (INSTR(15).NE.03) WRITE(IOUT,652) INSTR(15)
651 FORMAT('++',20X,'DO NOT READ SOLUTION FROM PERMANENT DATA SET')
652 FORMAT('++',20X,'READ SOLUTION FROM PERMANENT DATA SET',I3)
    GO TO 1000
660 WRITE(IOUT,661) INSTR(16)
661 FORMAT('++',20X,'FIRST RECORD TO BE READ:',I4)
    GO TO 1000
670 WRITE(IOUT,671) INSTR(17)
671 FORMAT('++',20X,'LAST RECORD TO BE READ:',I4)
    GO TO 1000
680 IF (INSTR(18).EQ.03) WRITE(IOUT,681)
    IF (INSTR(18).NE.03) WRITE(IOUT,682) INSTR(18)
681 FORMAT('++',20X,'DO NOT WRITE SOLUTION ON PERMANENT DATA SET')
682 FORMAT('++',20X,'WRITE SOLUTION ON PERMANENT DATA SET',I3)
    GO TO 1000
690 WRITE(IOUT,691) INSTR(19)
691 FORMAT('++',20X,'FIRST RECORD TO BE WRITTEN:',I4)
    GO TO 1000
700 IF ((INSTR(11).NE.0).AND.(INSTR(20).NE.0)) WRITE(IOUT,701) INSTR(2
    0)
    IF ((INSTR(11).EQ.0).OR.(INSTR(20).EQ.0)) WRITE(IOUT,702)
701 FORMAT('++',20X,'INCREMENT INSTR(11) BY',I3,' WHENEVER WEIGHT IS H
    ALVED')
702 FORMAT('++',20X,'DO NOT MODIFY INSTR(11)')
    GO TO 1000
710 IF (INSTR(21).EQ.0) WRITE(IOUT,711)
    IF (INSTR(21).NE.0) WRITE(IOUT,712)
711 FORMAT('++',20X,'DO NOT CHANGE EXTRAPOLATION OPTION (INSTR(7))')
712 FORMAT('++',20X,'INVOKE EXTRAPOLATION OPTION (SET INSTR(7)=2) WHEN
    *WEIGHT IS HALVED')
    GO TO 1000
720 WRITE(IOUT,721) INSTR(22)
721 FORMAT('++',20X,'SET INSTR(23)=2 WHEN ITER >=',I7)
    GO TO 1000
730 IF (INSTR(23).EQ.0) WRITE(IOUT,731)
    IF (INSTR(23).EQ.1) WRITE(IOUT,732)
    IF (INSTR(23).EQ.2) WRITE(IOUT,733)
731 FORMAT('++',20X,'DO NOT USE SMALL PERTURBATION EXPANSIONS')
732 FORMAT('++',20X,'USE SMALL PERTURBATION EXPANSIONS AS INITIAL GUESS
    1 FOR NEXT TIME INTERVAL')

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733	FORMAT('++',20X,'USE SMALL PERTURBATION EXPANSION AS SOLUTION')		
	GO TO 1000		
740	IF(INSTR(24).EQ.0)WRITE(IOUT,741)		
	IF(INSTR(24).NE.0)WRITE(IOUT,742)		
741	FORMAT('++',20X,'RESULTS FROM SMPERT NOT PRINTED')		
742	FORMAT('++',20X,'RESULTS FROM SMPERT ARE PRINTED')		
	GO TO 1000		
750	IF(INSTR(25).EQ.1)WRITE(IOUT,751)		
	IF(INSTR(25).EQ.2)WRITE(IOUT,752)		
751	FORMAT('++',20X,'ASSUME PLENUM IS ISENTROPIC')		
752	FORMAT('++',20X,'SET TP AND TPT = MAX(TSEN TP,TCT0)')		
	GO TO 1000		
760	WRITE(IOUT,761)INSTR(26)		
751	FORMAT('++',20X,'REVERT TO EXACT SUPERSONIC SOLUTION ',I7,' TIME IN		
	CREMENTS AFTER CHOKE')		
	GO TO 1000		
1000	CONTINUE		
	WRITE(IOUT,1110)(J,J=1,26),JV		
1110	FORMAT('1',26(' J',I2)/' ',I32('=')/' ',I26I5)		
	I1=0		
	DO 1020 I=1,3		
	IF(NVT(I).GT.I1)I1=NVT(I)		
1020	CONTINUE		
	WRITE(IOUT,1040)((J,I=1,2),J=1,3),(I,(TV(J,I),ARFA(J,I),J=1,3),		
	I=1,I1)		
1040	FORMAT('0FLOW AREAS VERSUS TIME',I0,I,3(5X,'TV(',I1,' ',I),',		
	1 AX,'ARFA(',I1,' ',I),',3X)/' ',I99('='),50('/' ',I3,6E16,A))		
	IF(NCTI.EQ.A)RETURN 1		
	RETURN		
	END		

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SUBROUTINE SOLVER(*)
  IMPLICIT REAL*8 (A-H,M,O-Z,$)
  COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),F(7),B(30),
  1 TVE(3),TDELAY(3),RW(7)
  COMMON PC,RC,TC,AC,MDCTC,FAE,EAPE,FAF,MDTSTR,INFIN,TM30GS,GPO2GS,
  1 SGOR
  COMMON G,GM1,GPI,OGG,GPI02,GM102,GPI0G,GM10G,GOGPI,GOGM1,
  1 OOGM1,OGGPI,GPOGM1,SGM102,TGGM1,TGG,MGP6M2,TGGPI,GPGM12,
  2 MPOGM,MPOGM1,R,GR,ORR,PI,PERR,AWOKW,OOA1,OOKF,KF,KW
  COMMON TSL,TSH,TSW,TSP,TTA,TTWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
  COMMON T,T1,DT,TSTR,DT02,TSTOP,ODDT,OTOPV
  COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
  COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
  - MDD, MDF, MDPT, MDCT, MDPE, MDT0, MDCT0, TE0, TP,
  - TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
  - APE, AF, MN, MD
  COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
  - MDD1, MDF1, MDPT1, MDCT1, MDPE1, MDT01, MDCT01, TE01, TP1,
  - TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
  - APE1, AF1, MN1, MD1
  COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
  - MDD2, MDF2, MDPT2, MDCT2, MDPE2, MDT02, MDCT02, TE02, TP2,
  - TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
  - APE2, AF2, MN2, MD2
  COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
  - MDD3, MDF3, MDPT3, MDCT3, MDPE3, MDT03, MDCT03, TE03, TP3,
  - TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
  - APE3, AF3, MN3, MD3
  COMMON PSOP0,TSOT0,RSOR0,MSOM0
  COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$E2,$EMAX,$EP,$DE
  COMMON INSTR(26),IDERUG,IIN,IOUT,NP,IP,ITER,NVT(3),I,NT,IPAGE,
  INPAGE,$N,IT(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
  2IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
  COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
  1 J18,J19,J20,J21,J22,J23,J24,J25,J26
  COMMON NCTL
  INTEGER $N
  REAL*8 INFIN,KF,KW
  C ELLIPTIC ENERGY EQUATION GIVING PRESSURE FROM MASS FLUX
  GUESS1(D1)=PSOP0+IFLG2*A6*DSQRT(1.-(A5*D1)**2)
  C ELLIPTIC ENRGY EQUATION GIVING MACH NUMBER FROM MASS FLUX
  GUESS2(D1)=DSQRT((GUESS1(D1)**A8 -1.)*TOGM1)
  C ELLIPTIC ENRGY/CONTINUITY EQUATION GIVING MACH NUMBER FROM AREA RATIO
  GUESS3(D1)=GUESS2(1./D1)
  C APPROXIMATE UNSTEADY WAVE EQUATION GIVING MACH NUMBER FROM MASS FLUX
  GUESS4(D1)=OGGPI*(1.-DSQRT(1.-A9*D1))
  SOX=.001
  GO TO (10,20,30,40),IFLG
  10 $X1=GUESS1($Y)
  GO TO 50
  20 IF($Y.GE.4SOM0)GO TO 25
  $X1=GUESS2($Y)
  GO TO 50
  25 $X1=1.00
  $N=0
  RETURN

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30	\$X1=GUESS3(\$Y*45)		
	GO TO 50		
40	IF(INSTR(8).EQ.1)GO TO 45		
	\$X1=GUESS4(\$Y)		
	GO TO 50		
45	\$X1=0.		
	DO 46 I=1,7		
45	\$X1=\$X1+RW(I)*\$Y**I		
50	IF(INSTR(9).EQ.2)RETURN		
	WRITE(IDEBUG,100)\$Y,\$DX,\$EMAX		
100	FORMAT(' SOLVER/'0\$Y='F16.8,' \$DX='F16.8,' \$EMAX='F16.8/ 1'0 N'5X,'X(N)'10X,'Y(N)'9X,'X(N-1)'8X,'Y(N-1)'9X,'E(N)'2 2 9X,'E(N-1)'10X,'DE'12X,'EP'12X,'DX'/' '132(''-'))		
	\$N=0		
120	GO TO(1100,1200,1300,1400),IFLG		
130	\$E1=(\$Y-\$Y1)/\$Y		
	IF(\$N.NF.0)GO TO 180		
140	\$E2=\$E1		
	\$Y2=\$Y1		
	\$X2=\$X1		
	\$X1=\$X2+\$DX		
150	\$N=\$N+1		
	GO TO 120		
180	\$EP=\$E1*\$E2		
	\$DF=DABS(\$E1)-DABS(\$E2)		
	WRITE(IDEBUG,200)\$N,\$X1,\$Y1,\$X2,\$Y2,\$E1,\$E2,\$DF,\$EP,\$DX		
200	FORMAT(' ',I3,9F14.6)		
	IF(DABS(\$E1).LE.\$EMAX)RETURN		
	IF(\$EP.GT.0.)GO TO 220		
	\$DX=.5*\$DX		
210	\$X1=\$X2+\$DX		
	GO TO 150		
220	IF(\$DE.LT.0.)GO TO 140		
	\$DX=-\$DX		
	GO TO 210		
	C ENERGY EQUATION GIVING MASS FLUX FROM PRESSURE		
1100	\$Y1=\$X1**TOG-\$X1**GP1OG		
	\$Y1=USORT(\$Y1)		
	GO TO 130		
	C ENERGY EQUATION GIVING MASS FLUX FROM MACH NUMBER		
1200	\$Y1=\$X1*(1.+GM102*\$X1**2)**MGPGM2		
	GO TO 130		
	C AREA RATIO VERSUS MACH NUMBER		
1300	\$Y1=(TOGP1*(1.+GM102*\$X1**2)**GPGM12/\$X1		
	GO TO 130		
	C UNSTEADY MASS FLUX FROM MACH NUMBER		
1400	\$Y1=\$X1*(1.+GM102*\$X1)**MGPOGM		
	GO TO 130		
	END		

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SUBROUTINE PRINT(*)
IMPLICIT REAL*8 (A-H,M,O-Z,S)
COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),E(7),B(30),
1 TVE(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MNCTC,FAF,EAPE,FAF,MDTSTR,INFIN,TMGOGS,GP02GS,
1 SGOR
COMMON G,GM1,GP1,OGG,GP102,GM102,GP10G,GM10G,GOGP1,GOGM1,
1 OOGM1,OGGP1,GP0GM1,SGM102,T0GM1,T0G,MGP0GM2,T0GP1,GP0GM12,
2 MGPOGM,M0SGM1,R,GR,00R,PI,PERR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TSH,TSH,TSP,ISA,ISWA,TSV,CTD,CTA,PV,PV0TSV,TAUW
COMMON T,T1,DT,TSTR,DT02,TSTOP,00DT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MDD, MDE, MDPT, MDCT, MDPE, MDTS0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDD1, MDE1, MDPT1, MDCT1, MDPE1, MDTS01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDD2, MDE2, MDPT2, MDCT2, MDPE2, MDTS02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDD3, MDE3, MDPT3, MDCT3, MDPE3, MDTS03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON P5OP0,T5OT0,RSOR0,MSOM0
COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$E2,$EMAX,$EP,$DF
COMMON INSTR(26),TDERUG,ITN,IOUT,NP,IP,ITER,NVT(3),I,NT,TPAGE,
INPAGE,$N,IT(3),J,IM1,ITIME,ND,NN,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
DIMENSION V(24,4)
EQUIVALENCE (PN,V(1,1))
INTEGER $N
REAL*8 INFIN,KF,KW
IP=IP+1
IF(IP.NE.NP)RETURN
IP=0
I1=J16-1
IF(INSTR(13).NE.0)GO TO 300
IF(IPAGE.EQ.0)GO TO 100
IF(IPAGE.LT.NPAGE)GO TO 140
100 WRITE(IOUT,120)(I,I=1,7)
120 FORMAT(11,AX,'T',13X,'TSTR',13X,'T1',14X,'PT',14X,'PP',14X,
1'PD',14X,'PN',14X,'PPT',14X,'ND',14X,'PCT0',13X,'PE0',13X,'MCT',
-13X,'MD',
2'14X,'MN',13X,'MDCT',12X,'MDPT',13X,'MDD',8X,'NM',14X,'MDE',13X,
3'MDE',12X,'MDPE',12X,'MDTS0',11X,'MDCT0',12X,
4'TP',14X,'TPT',13X,'TE0',7X,'NCT',14X,'TCT0',
5 13X,'PP',14X,'RPT',13X,'RE0',12X,'RCT0',13X,'AE',14X,'APE',
613X,'AF',7X,'ITER',14X,'E(1,T1)',6X,'5X',DT,6X,'J16'

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SURROUTINE RINOM(*)
IMPLICIT REAL*8 (A-H,M,O-Z,S)
COMMON ARFA(3,50),ARFATS(3),AREAM(3),TV(3,50),A(10),E(7),B(30),
1 TVF(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDCTC,EAF,EAPF,EAF,MDTSTR,INFIN,IMGOGS,GPO2GS,
1 SGJR
COMMON G,SM1,GP1,OGG,GP102,GM102,GP10G,GM10G,GOGP1,GOGM1,
1 GOGM1,OGGP1,GPOGM1,SGM102,TGGM1,TGG,MGPGM2,TGGP1,GPGM12,
2 MGPOGM,MSGGM1,R,GR,OCR,PI,PERR,AWOKW,OOA1,OOKF,KF,KW
COMMON TSL,TSH,TSW,TSP,TSA,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DTU2,TSTOP,ODDT,DTOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MDD, MDF, MDPT, MDCT, MDPE, MDTS0, MDCT0, TE0, TP,
- TPI, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDD1, MDF1, MDPT1, MDCT1, MDPE1, MDTS01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDD2, MDF2, MDPT2, MDCT2, MDPE2, MDTS02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDD3, MDF3, MDPT3, MDCT3, MDPE3, MDTS03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PSOP0,TSOT0,RSOR0,MSOM0
COMMON $Y,$Y1,$Y2,$X1,$X2,$X,$E1,$E2,$EMAX,$EP,$DE
COMMON INSTR(26),IDEBUG,YIN,IOUT,NP,IP,ITER,NVT(3),I,NT,IPAGE,
1 NPAGE,$N,IT(3),J,IM1,ITIME,ND,NN,NCT,IFLG1,IFLG2,IFLG3,IFLG4,
2 IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
INTEGER $N
REAL*8 INFIN,KF,KW
J=0
A(1)=1.
10 I=I+1
IF(I.GT.6)GO TO 20
IM1=I-1
A(I+1)=A(I)*(A(8)-IM1)/I
WRITE(IDEBUG,15)I,I1,IM1
15 FORMAT(' I=',I7,' I1=',I7,' IM1=',I7)
GO TO 10
20 WRITE(IDEBUG,30)A
30 FORMAT('09RINOM/' ' ,10E13.5)
DO 40 I=1,7
40 A(I)=A(I)*A(9)**(I-1)
WRITE(IDEBUG,30)A
RETURN
END

```

REVERT

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SUBROUTINE REVERT(*)
IMPLICIT REAL*8 (A-H,M,O-Z,%)
COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),E(7),B(30),
1 TVF(3),TDELAY(3),RW(7)
COMMON PC,RC,TC,AC,MDCIC,EAF,EAPF,FAF,MDTSTR,INFIN,TMGOGS,GP02GS,
1 SGOR
COMMON G,GM1,GP1,00G,GP102,GM102,GP10G,GM10G,G0GP1,G0GM1,
1 00GM1,00GP1,G00GM1,5GM102,T0GM1,T0G,M0PGM2,T0GP1,GPGM12,
2 MGPOGM,M00GM1,R,GR,00R,PT,PERR,AWOKW,00A1,00KF,KF,KW
COMMON TSL,TSH,TSW,TSP,TTSA,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAUW
COMMON T,T1,DT,TSTR,DT02,TSTOP,000T,0TOPV
COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MDD, MDF, MDPT, MDCT, MDPE, MDTS0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MDD1, MDF1, MDPT1, MDCT1, MDPE1, MDTS01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MDD2, MDF2, MDPT2, MDCT2, MDPE2, MDTS02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MDD3, MDF3, MDPT3, MDCT3, MDPE3, MDTS03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
COMMON PSOP0,TSOT0,RSOR0,MSOMO
COMMON SY,$Y1,$Y2,$X1,$X2,$DX,$E1,$E2,$EMAX,$EP,$DE
COMMON INSTR(26),IDEBUG,ITN,IOUT,NP,IP,ITER,NVT(3),I,VT,IPAGE,
1 NPAGE,$N,I(3),J,IM1,ITIME,ND,NN,NCT,IPLG,IPLG1,IPLG2,IPLG3,IPLG4,
2 IPLG5,IPLG6,IPLG7,IPLG8,IPLG9,I1,I2,I3,I4,I5
COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
COMMON NCTL
INTEGER $N
REAL*8 INFIN,XF,KW
R(1)=1./A(1)
R(2)=-A(2)/A(1)**3
R(3)=(2.*A(2)**2-A(1)*A(3))/A(1)**5
R(4)=(5.*A(1)*A(2)*A(3)-A(1)**2*A(4)-5.*A(2)**3)/A(1)**7
R(5)=(4.*A(1)**2*A(2)*A(4)+3.*A(1)**2*A(3)**2+14.*A(2)**4-
A(1)**3*A(5)-21.*A(1)*A(2)**2*A(3))/A(1)**9
R(6)=(7.*A(1)**3*A(2)*A(5)+7.*A(1)**3*A(3)*A(4)+84.*A(1)*A(2)**3
+ A(3)-A(1)**4*A(6)-28.*A(1)**2*A(2)**2*A(4)-28.*A(1)**2*A(2)
+ A(3)**2-42.*A(2)**5)/A(1)**11
R(7)=(8.*A(1)**4*A(2)*A(6)+8.*A(1)**4*A(3)*A(5)+4.*A(1)**4*A(4)**2
+ 120.*A(1)**2*A(2)**3*A(4)+180.*A(1)**2*A(2)**2*A(3)**2+
R 132.*A(2)**6-A(1)**5*A(7)-36.*A(1)**3*A(2)**2*A(5)-
C 72.*A(1)**3*A(2)*A(3)*A(4)-12.*A(1)**3*A(3)**3-330.*A(1)*A(2)
D **4*A(3))/A(1)**13
DO 10 I=1,7
10 A(I)=R(I)
WRITE(IDEBUG,20)A
20 FORMAT('0REVERT'/,' ',10E13.5)
RETURN
END

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SMPERT

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SURROUTINE SMPERT(*)
  IMPLICIT REAL*8 (A-H,M,O-Z,%)
  COMMON AREA(3,50),AREATS(3),AREAM(3),TV(3,50),A(10),F(7),B(30),
1 TVF(3),TDELAY(3),RW(7)
  COMMON PC,RC,TC,AC,MDCIC,FAF,FAPE,FAF,MDTSTR,INFIN,IM30GS,GP02GS,
1 SGOR
  COMMON G,GM1,GP1,00G,GP102,GM102,GP10G,GM10G,GP0GP1,G03M1,
1 G03M1,00GP1,GP0GM1,SGM102,T0GM1,T0G,MGP0M2,T0GP1,GP0M12,
2 MGP0GM,M30GM1,P,GR,00R,PT,PERR,AWOKW,00A1,00KF,KF,KW
  COMMON TSL,TS4,TSW,TSP,TSa,TSWA,TSV,CTD,CTA,PV,PVOTSV,TAIW
  COMMON T,T1,DT,TSTR,DT02,TSTOP,00NT,DTOPV
  COMMON A1,A2,A3,A4,A5,A6,A7,A8,A9,A10,A11,A12,A13,A14,A15,A16,A17
  COMMON PN, PP, PPT, PD, PT, PCT0, PE0, MDE,
- MD0, MDE, MDPT, MDCT, MDPE, MDTS0, MDCT0, TE0, TP,
- TPT, TCT0, RP, RPT, RE0, RCT0, ACT0, MCT, AE,
- APE, AF, MN, MD
  COMMON PN1, PP1, PPT1, PD1, PT1, PCT01, PE01, MDE1,
- MD01, MDE1, MDPT1, MDCT1, MDPE1, MDTS01, MDCT01, TE01, TP1,
- TPT1, TCT01, RP1, RPT1, RE01, RCT01, ACT01, MCT1, AE1,
- APE1, AF1, MN1, MD1
  COMMON PN2, PP2, PPT2, PD2, PT2, PCT02, PE02, MDE2,
- MD02, MDE2, MDPT2, MDCT2, MDPE2, MDTS02, MDCT02, TE02, TP2,
- TPT2, TCT02, RP2, RPT2, RE02, RCT02, ACT02, MCT2, AE2,
- APE2, AF2, MN2, MD2
  COMMON PN3, PP3, PPT3, PD3, PT3, PCT03, PE03, MDE3,
- MD03, MDE3, MDPT3, MDCT3, MDPE3, MDTS03, MDCT03, TE03, TP3,
- TPT3, TCT03, RP3, RPT3, RE03, RCT03, ACT03, MCT3, AE3,
- APE3, AF3, MN3, MD3
  COMMON PSOP0,TS0T0,RSOR0,MSOM0
  COMMON $Y,$Y1,$Y2,$X1,$X2,$DX,$E1,$F2,$FMAX,$EP,$DE
  COMMON INSR(26),IDFRUG,ITN,IOUT,NP,IP,ITER,NVT(3),I,NT,IPAGE,
INPAGE,$N,IT(3),J,IM1,TIME,N0,NV,NCT,IFLG,IFLG1,IFLG2,IFLG3,IFLG4,
2IFLG5,IFLG6,IFLG7,IFLG8,IFLG9,I1,I2,I3,I4,I5
  COMMON J1,J2,J3,J4,J5,J6,J7,J8,J9,J10,J11,J12,J13,J14,J15,J16,J17,
1 J18,J19,J20,J21,J22,J23,J24,J25,J26
  COMMON NCTL
  DIMENSION IEXTP(19),FPS(19),Q(19),V(30,4)
  DIMENSION JJ(4)
  EQUIVALENCE (V(1),PN)
  COMPLEX*16 X(4),Y(4)
  INTEGER $N
  REAL*8 NU1,NU2,NU10,NU20,NU1N,NU2N,INFIN,KF,KW
  C
  DATA IEXTP/ 8,13,10,11,21, 9,12,25, 7,16, 5, 4, 1,14, 0, 0, 2,20,
# 17/
  MACH(D1) = DSORT(T0GM1)*((D1 / PCT0)*((-GM10G)- 1.00))
  WRITE(IDFRUG,1)
  1 FORMAT('1SMPERT',/,' ',6(' ',))
  C-----
  C COMPUTE CONSTANTS FOR EXPANSIONS
  C-----
  IF(IFLG2.NE.1)GO TO 90
  C EQUATION 1
  CAL =-1.00
  A(1) = A1 / DSORT( TE01 ) * A2
  CB1 = AE1 * A(1)

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      CC1 = PF01 * A(1)
      CD1 = 0.500 * PF01 * A(1) * AE1 / TE01
C EQUATION 2
      90 CA2 = -1.00
      A(2) = A1 / DSORT( TP1 ) * A2
      CR2 = -APE1 * A(2)
      CC2 = -PP1 * A(2)
      CD2 = 0.500 * PP1 * A(2) * APE1 / TP1
C EQUATION 3
      CA3 = -1.00
      CR3 = -(PD1 - PD1 * A16) * DOKF * A2
      CC3 = -AF1 * DOKF * A2
      CD3 = -CC3 * A16
C EQUATION 4
      CA4 = -1.00
      CR4 = -AWOKW * A2
      CC4 = -CB4 * A15
C EQUATION 5
      CA5 = -1.00
      CR5 = DTOPV
      CC5 = CR5
      CD5 = CR5
      CE5 = CR5 * ( MDPE1 + MDF1 + MDPT1 )
C EQUATION 6
      CA6 = 1.00
      CR6 = 1.00
      CC6 = -1.00
      IF (IFLG2.NE.1) GO TO 91
C EQUATION 7
      CA7 = 1.00
      CR7 = -1.00
      CC7 = -1.00
C EQUATION 8
      CA8 = -1.00
      A(3) = 1.00 - MCT1
      A(4) = 1.00 + GM102 * MCT1
      CRA = MDCTC * A(3) / A(4) ** (GOGM1 * 2.00)
C EQUATION 9
      CA9 = -1.00
      A(5) = 1.00 + GM102 * MCT1 ** 2
      CR9 = -G * A(3) * PF01 / A(4) / A(5)
C EQUATION 10
      CA10 = -1.00
      CR10 = -GM1 * A(3) * TE01 / A(4) / A(5)
C EQUATION 11
      91 CA11 = -1.00
      IF (IFLG2) 99,9999,100
100 CR11 = 0.500
      CC11 = CR11
      GO TO 101
      90 CR11 = 1.00 - A17
      CC11 = A17
C EQUATION 12
101 A(9) = -GM1 / MDT501 ** 2
      A(10) = GM1 / MDT501 ** 3
      CA12 = A(9) * MOD1

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CB12 = A(10) * MD01 **2
IF (IFLG2.EQ.1) A(6)=PD1/PE01
IF (IFLG2.NE.1) A(6)=PD1/PCT01
A(7) = A(5) ** TOG
A(8) = A(6) ** GP10G
NU10 = TOG * A(7) - GP10G * A(8)
NU20 = TMG0GS * A(7) - GP02GS * A(8)
IF (IFLG2) 94,9999,93
93 CC12 = NU10 / PD1 / PF01
   CD12 = NU20 / ( PD1 * PE01 ) ** 2
   GO TO 95
94 CC12= NU10/PD1/PCT01
   CD12=NU20/(PD1*PCT01)**2
95 IF (IFLG2.NE.1) GO TO 92
C EQUATION 13
CA13 = A(9) * MDCT1
CB13 = A(10) * MDCT1 ** 2
B(1) = PN1 / PE01
B(2) = B(1) ** TOG
B(3) = B(1) ** GP10G
NU1N=TOG*B(2)-GP10G*B(3)
NU2N=TMG0GS*B(2)-GP02GS*B(3)
CC13=NU1N/PN1/PE01
CD13=NU2N/(PN1*PE01)**2
C EQUATION 14
CA14 = -1.D0
CB14 = TSA * SGOR / DSQRT( TE01 ) * A2
CC14 = -0.5D0 * PE01 * CB14 / TE01
C EQUATION 17
92 GO TO(96,98),IFLG9
98 CA17=A2
   CB17=-R*IC101
   GO TO 97
96 CA17 = RP1
   CB17 = - G * PP1
C EQUATION 18
97 CA18= -1.D0
   CB18 = 0.5D0
   CC18 = RPI1 - RP1
C EQUATION 19
CA19 = R * RP1
CB19 = R * TP1
CC19 = -A2
IF (IFLG2) 2,9999,3
C *****
C SUPERSONIC BRANCH
C *****
2 IF (IDRUG.EQ.03) GO TO 70
   CA1 = INFIN
   CB1 = INFIN
   CC1 = INFIN
   CD1 = INFIN
   CA7 = INFIN
   CB7 = INFIN
   CC7 = INFIN
   CAR = INFIN

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CRP = INFIN		
CA9 = INFIN		
CB9 = INFIN		
CC6 = INFIN		
CA10 = INFIN		
CB10 = INFIN		
CB11 = INFIN		
CB12 = INFIN		
CA13 = INFIN		
CB13 = INFIN		
CC13 = INFIN		
CD13 = INFIN		
CA14 = INFIN		
CB14 = INFIN		
CC14 = INFIN		
C EQUATION 11		
70 SALP11 = -CC11 / CA11		
C EQUATION 18		
SALP18 = -CB18 / CA18		
SBET18 = -CC18 / CA18		
C EQUATION 19		
A(1) = -1.00 / CA19		
SALP19 = CB19 * SALP18 * A(1)		
SBET19 = CC19 * A(1)		
SGAM19 = CB19 * SBET18 * A(1)		
C EQUATION 2		
A(2) = -1.00 / CA2		
SALP2 = (CB2 + CD2 * SBET19) * A(2)		
SBET2 = CD2 * SALP19 * A(2)		
SGAM2 = (CC2 * FAF + CD2 * SGAM19) * A(2)		
C EQUATION 3		
A(3) = -1.00 / CA3		
SALP3 = CC3 * A(3)		
SBET3 = CD3 * A(3)		
SGAM3 = CB3 * EAF * A(3)		
C EQUATION 17		
A(4) = -CB17 / CA17		
SALP17 = SALP18 * A(4)		
SBET17 = SBET18 * A(4)		
C EQUATION 4		
SALP4 = CC4 * SALP11		
C EQUATION 5		
SALP5 = CA5 + CB5 * SBET2		
SBET5 = CB5 * SALP2 + CC5 * SALP3		
SGAM5 = CC5 * SBET3		
SEPS5 = CB5 * SGAM2 + CC5 * SGAM3 + CF5		
C EQUATION 4 CONTINUED		
SBET4 = CB4 * SALP17		
SGAM4 = CB4 * SBET17		
C EQUATION 5 CONTINUED		
SIOT5 = SALP5 + SALP17 * SBET5		
A(5) = -1.00 / SIOT5		
SZET5 = SGAM5 * A(5)		
SET45 = CD5 * A(5)		
SKAP5 = (SBET5 * SBET17 + SEPS5) * A(5)		
C EQUATION 4 CONTINUED		

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SETA4 = CA4 * SRET4 * SFTA5
SEPS4 = -(SBET4 * SZET5 + SALP4) / SETA4
SZET4 = -(SRET4 * SKAP5 + SGAM4) / SETA4

```

C EQUATION 6

```

A(6) = -CB6 / CA6
SALP6 = SEPS4 * A(6)
SBET6 = SZET4 * A(6)

```

C EQUATION 12

```

SGAM12 = CD12 * PCT01 ** 2
SALP12 = (CC12 * PCT01 + CA12 * SALP6) / SGAM12
SBET12 = CA12 * SBET6 / SGAM12
IF (DEBUG.EQ.03) GO TO 80

```

SEPS3 = INFIN

SGAM6 = INFIN

SEPS6 = INFIN

SZET6 = INFIN

SETA6 = INFIN

SIOT6 = INFIN

SLAM6 = INFIN

SMU6 = INFIN

SNU6 = INFIN

SALP7 = INFIN

SBET7 = INFIN

SGAM7 = INFIN

SEPS7 = INFIN

SZET7 = INFIN

SETA7 = INFIN

SIOT7 = INFIN

SKAP7 = INFIN

SALP8 = INFIN

SALP9 = INFIN

SBET9 = INFIN

SGAM9 = INFIN

SZET9 = INFIN

SALP10 = INFIN

SBET11 = INFIN

SALP13 = INFIN

SBET13 = INFIN

SGAM13 = INFIN

SALP14 = INFIN

SBET14 = INFIN

SGAM14 = INFIN

SEPS14 = INFIN

SZET14 = INFIN

SETA14 = INFIN

SIOT14 = INFIN

SGAM17 = INFIN

SEPS17 = INFIN

80 A(7) = 0.5D0 * SALP12

Y(2) = DSQRT(A(7) ** 2 - SBET12)

Y(1) = -A(7) + Y(2)

Y(2) = -A(7) - Y(2)

C-----

C SORT ROOTS

C-----

I=0

```

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DO 200 K = 1,2
IF(ABS(DIMAG(Y(K))),GT,1.D-12)GO TO 200
B(15) = DREAL(Y(K))/PD1
WRITE(IDEBUG,9)K,Y(1),Y(2),A14,B(15)
9 FORMAT(' K=',I1,' Y=',4E13.5,' A14=',E13.5,' B(15)=',E13.5)
IF(ABS(B(15)),GT,ABS(A14))GO TO 200
I = I + 1
EPS(12) = Y(K)
200 CONTINUE
IF(I,LE,1)GO TO 205
IF(ABS(DREAL(Y(2))),LT,ABS(DREAL(Y(1))))Y(1)=Y(2)
EPS(12)=Y(1)
I=1
205 IFLG8=0
K = 1
IF(I,EQ,1)GO TO 230
WRITE(ROUT,210)I
210 FORMAT('0SMPERT (SUPERSONIC)',I3,' SOLUTIONS FOUND')
IDEBUG = 06
IFLG8 = 1
K = 0
220 K=K+1
IF(K,GT,2)GO TO 60
EPS(12) = Y(K)
C=====
C COMPUTE INCREMENTS
C=====
230 EPS(6) = SALP6 * EPS(12) + SBET6
EPS(4) = SEPS4 * EPS(12) + SZET4
EPS(5) = SZET5 * EPS(12) + SETA5 * EPS(4) + SKAP5
EPS(17) = SALP17 * EPS(5) + SBET17
EPS(3) = SALP3 * EPS(17) + SBET3 * EPS(12) + SGAM3
EPS(2) = SALP2 * EPS(17) + SBET2 * EPS(5) + SGAM2
EPS(19) = SALP19 * EPS(5) + SBET19 * EPS(17) + SGAM19
EPS(18) = SALP18 * EPS(5) + SBET18
EPS(11) = SALP11 * EPS(12)
EPS(1) = 0.00
EPS(7) = 0.00
EPS(8) = 0.00
EPS(9) = 0.00
EPS(10) = 0.00
EPS(13) = 0.00
EPS(14) = 0.00
WRITE(IDEBUG,20)I,I=1,24),EPS
C=====
C COMPUTE PROPERTY VALUES
C=====
DO 240 I = 1,19
J = IEXTP(I)
IF(J,EQ,0)GO TO 240
V(J,1) = V(J,2) + EPS(I)
240 CONTINUE
TE0=TC0
MDE = MDO - MDE
PE0 = MDE * DSQRT(TF0) / (A1 * A2 * AE)
GO TO(239,238),IFLG9

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239 PPT = PPT1 * (RPT / RPT1) ** 6
    TPT = PPT * A2 * OOR / RPT
    GO TO 237
238 TPT=TCID
    PPT=RPT * R * TPT / A2
237 REF = PEO * A2 * OOR / TEO
    MD = MACH(PD)
    IF((INSTR(23),NF,2).AND.(INSTR(24),EQ,0))GO TO 241
    J16=J16+1
    CALL PRINT
    J16=J16-1
    WRITE(IOUT,242)
242 FORMAT('3')
241 IF((IFLG8,EQ,0).AND.(IDBUG,EQ,03))RETURN
    WRITE(IDBUG,32)V
C-----
C SMALL PERTURBATION RESIDUALS
C-----
    DO 250 I = 1,19
250 Q(I) = INFIN
    Q(2) = CA2 * EPS(2) + CB2 * EPS(17) + CC2 * EAPE + CD2 * EPS(19)
    Q(3) = CA3 * EPS(3) + CB3 * EAF + CC3 * EPS(17) + CD3 * EPS(12)
    Q(4) = CA4 * EPS(4) + CB4 * EPS(17) + CC4 * EPS(11)
    Q(5) = CA5 * EPS(5) + CB5 * EPS(2) + CC5 * EPS(3) + CD5 * EPS(4)
    # + CE5
    Q(6) = CA5 * EPS(6) + CB6 * EPS(4)
    Q(7) = INFIN
    Q(8) = INFIN
    Q(9) = INFIN
    Q(10) = INFIN
    Q(11) = CA11 * EPS(11) + CC11 * EPS(12)
    Q(12) = CA12 * EPS(6) + CC12 * PCT01 * EPS(12) + CD12 * PCT01 ** 2
    # * EPS(12) ** 2
    Q(13) = INFIN
    Q(14) = INFIN
    Q(17) = CA17 * EPS(17) + CB17 * EPS(18)
    Q(18) = CA18 * EPS(18) + CB18 * EPS(5) + CC18
    Q(19) = CA19 * EPS(19) + CB19 * EPS(18) + CC19 * EPS(17)
    WRITE(IDBUG,13) (I,I=1,20),Q
C-----
C EXACT RESIDUALS
C-----
    DO 40 I=1,19
40 Q(I)=1.D70
    Q(2)=MDF+A1*PP*APE/DSQRT(TPT*A2)
    Q(3)=MDF+AF*QOKF*(PP-PD*A16)*A2
    Q(4)=MDPT+AWOKW*(PP-PT*A15)*A2
    Q(5)=RPT-RPT1=(MDF+MDF+MDPT)*DTOPV
    Q(6)=MOD+MDPT-MDCT
    IF(IFLG2)254,9999,255
254 Q(11)=PT-(1.D0-A17)*PN-A17*PD
    GO TO 256
255 Q(11)=PT-0.5D0*(PN+PD)
256 R(12)=1.D0/PCI0
    R(13)=TOG+1*MDTS0**2
    Q(12)=MOD**2-9(13)*((PD*R(12))**TOG-(PD*R(12))**GP106)

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GO TO (251,252),IFLG9		
251 Q(17)=PP-PPT1*(RP/RPT1)**G		
GO TO 253		
252 Q(17)=PP-RP*R*TP/A2		
253 Q(18)=PP-0.5D0*(RPT+RPT1)		
Q(19)=TP-PP*00R/RP*A2		
GO TO 39		
C*****		
C SUBSONIC BRANCH		
C*****		
C-----		
C COMPUTE CONSTANTS FOR SOLUTION		
C-----		
C EQUATION 19		
3 SALP19 = -CB19 / CA19		
SBET19 = -CC19 / CA19		
C EQUATION 2		
B(4) = -1.00 / CA2		
SALP2 = (CB2 + CD2 * SBET19) * B(4)		
SBET2 = CD2 * SALP19 * B(4)		
SGAM2 = CC2 * EAPE * B(4)		
C EQUATION 5		
SALP5 = CB5 * SALP2		
SBET5 = CB5 * SBET2		
SGAM5 = CB5 * SGAM2 + CE5		
C EQUATION 8		
SALP8 = -CA8 / CB8		
C EQUATION 18		
SALP18 = -CA18 / CB18		
SBET18 = -CC18 / CB18		
C EQUATION 5 CONTINUED		
SKAP5 = CA5 * SALP18 + SBET5		
B(5) = 1.00 / SKAP5		
SEPS5 = -SALP5 * B(5)		
SZET5 = -CC5 * B(5)		
SETA5 = -CD5 * B(5)		
SIOT5 = - (CA5 * SBET18 + SGAM5) * B(5)		
C EQUATION 10		
SALP10 = -CB10 * SALP8 / CA10		
C EQUATION 11		
SALP11 = -CB11 / CA11		
SBET11 = -CC11 / CA11		
C EQUATION 17		
SEPS17 = CA17 + CB17 * SEPS5		
B(6) = -CB17 / SEPS17		
SALP17 = SZET5 * B(6)		
SBET17 = SETA5 * B(6)		
SGAM17 = SIOT5 * B(6)		
C EQUATION 1		
B(7) = -1.00 / CA1		
SALP1 = CB1 * B(7)		
SBET1 = CD1 * SALP10 * B(7)		
SGAM1 = CC1 * EAE * B(7)		
C EQUATION 3		
SEPS3 = CA3 + CC3 * SALP17		
B(8) = -1.00 / SEPS3		

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	SALP3 = CC3 * SRET17 * B(8)		
	SBET3 = CD3 * B(8)		
	SGAM3 = (CC3 * SGAM17 + CR3 * EAF) * B(8)		
C EQUATION 4			
	SEPS4 = CA4 + CR4 * (SBFT17 + SALP17 * SALP3)		
	B(9) = - 1.00 / SEPS4		
	SALP4 = (CR4 * SALP17 * SBET3 + CC4 * SBFT11) * B(9)		
	SBET4 = CC4 * SALP11 * B(9)		
	SGAM4 = CR4 * (SALP17 * SGAM3 + SGAM17) * B(9)		
C EQUATION 6			
	B(10) = - CR6 / CA6		
	SALP6 = SALP4 * B(10)		
	SBET6 = SBET4 * B(10)		
	SGAM6 = -CC6 / CA6		
	SEPS6 = SGAM4 * B(10)		
C EQUATION 7			
	SZET7 = CA7 * SGAM6 + CC7 * SBFT1		
	B(11) = - 1.00 / SZET7		
	SALP7 = (CA7 * SALP6 + CR7 * (SBET3 + SALP3 * SALP4)) * B(11)		
	SBFT7 = (CA7 * SBET6 + CR7 * SALP3 * SBET4) * B(11)		
	SGAM7 = CC7 * SALP1 * B(11)		
	SEPS7 = (CA7 * SEPS6 + CR7 * SGAM3 + CC7 * SGAM1 + CR7 * SALP3 * SGAM4) * B(11)		
C EQUATION 6 CONTINUED			
	SZET6 = SALP6 + SGAM6 * SALP7		
	SETA6 = SBET6 + SGAM6 * SBFT7		
	SIOT6 = SGAM6 * SGAM7		
	SKAP6 = SEPS6 + SGAM6 * SEPS7		
C EQUATION 14			
	B(12) = CC14 * SALP10		
	SALP14 = CR14 + B(12) * SGAM7		
	SBFT14 = B(12) * SALP7		
	SGAM14 = B(12) * SBFT7		
	SEPS14 = B(12) * SEPS7		
C EQUATION 9			
	B(13) = CB9 * SALP8		
	SZFT9 = CA9 + B(13) * SGAM7		
	SGAM9 = -B(13) / SZFT9		
	SALP9 = SGAM9 * SALP7		
	SBET9 = SGAM9 * SBFT7		
	SGAM9 = SGAM9 * SEPS7		
C EQUATION 14 CONTINUED			
	B(14) = - 1.00 / CA14		
	SZET14 = (SALP14 * SALP9 + SBET14) * B(14)		
	SETA14 = (SALP14 * SBET9 + SGAM14) * B(14)		
	SIOT14 = (SALP14 * SGAM9 + SEPS14) * B(14)		
C EQUATION 6 CONTINUED			
	SLAM6 = SZET6 + SIOT6 * SALP9		
	SMU6 = SETA6 + SIOT6 * SBET9		
	SNU6 = SKAP6 + SIOT6 * SGAM9		
C EQUATION 7 CONTINUED			
	SEIAT7 = SALP7 + SGAM7 * SALP9		
	SIOT7 = SBET7 + SGAM7 * SBET9		
	SKAP7 = SEPS7 + SGAM7 * SGAM9		
C EQUATION 12			
	SALP12 = DE01 - PD1 * SALP9		

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SRFT12 = -PD1 * SRFT9
SGAM12 = -PD1 * SGAM9
SA2 = CD12 * SALP12 ** 2
SB2 = CA12 * SLAM6 + CB12 * SZET14 + CC12 * SALP12 + 2.00 * CD12 *
@ SALP12 * SGAM12
SC2 = 2.00 * CD12 * SALP12 * SBET12
SD2 = CA12 * SMU6 + CR12 * SETA14 + CC12 * SRFT12 + 2.00 * CD12 *
@ SBET12 * SGAM12
SE2 = CD12 * SRFT12 ** 2
SF2 = CA12 * SNU6 + CR12 * SIOT14 + CC12 * SGAM12 + CD12 *
@ SGAM12 ** 2
C EQUATION 13
SALP13 = - PN1 * SALP9
SRFT13 = PE01 - PN1 * SRFT9
SGAM13 = - PN1 * SGAM9
SA3 = CD13 * SALP13 ** 2
SB3 = CA13 * SETA7 + CB13 * SZET14 + CC13 * SALP13 + 2.00 * CD13 *
@ SALP13 * SGAM13
SC3 = 2.00 * CD13 * SALP13 * SBET13
SD3 = CA13 * SIOT7 + CB13 * SETA14 + CC13 * SRFT13 + 2.00 * CD13 *
@ SRFT13 * SGAM13
SE3 = CD13 * SBET13 ** 2
SF3 = CA13 * SKAP7 + CR13 * SIOT14 + CC13 * SGAM13 + CD13 *
@ SGAM13 ** 2
II=IDERUG
CALL DSIMUL(SA2,SB2,SC2,SD2,SE2,SF2,SA3,SB3,SC3,SD3,SE3,SF3,II,X
1 ,Y)
C-----
C SORT ROOTS
C-----
I=0
DO 15 K=1,4
IF((DABS(DIMAG(X(K))) .GT. 1.0-12) .OR. (DABS(DIMAG(Y(K))) .GT. 1.0-12))
1 GO TO 15
R(14)=DREAL(X(K))/PD1
R(15)=DREAL(Y(K))/PN1
WRITE(IDERUG,69)K,X(K),Y(K),A14,R(14),R(15)
69 FORMAT(' K=',I1,' X=',2F13.5,' Y=',2F13.5,' A14=',F13.5,
1 ' R(14),R(15)=',2F13.5)
IF((DABS(R(14)) .GT. DABS(A14)) .OR. (DABS(R(15)) .GT. DABS(A14)))
1 GO TO 15
I=I+1
JJ(I)=K
EPS(12)=X(K)
EPS(13)=Y(K)
15 CONTINUE
IF(I.LF.1)GO TO 340
K=I
R(14)=X(JJ(I))
R(15)=Y(JJ(I))
DO 320 I=2,K
J=JJ(I)
IF(DABS(DREAL(X(J))) .GT. DABS(R(14)))GO TO 300
II=J
R(14)=X(J)
300 IF(DABS(DREAL(Y(J))) .GT. DABS(R(15)))GO TO 320

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I2=J
B(15)=Y(J)
320 CONTINUE
IF(I1.NE.I2)GO TO 340
I=1
EPS(12)=X(I1)
EPS(13)=Y(I2)
340 IFLG8=0
K=1
IF(I.EQ.1)GO TO 17
WRITE(10UT,16)I
16 FORMAT('0SMPERT:',I3,' SOLUTIONS FOUND:')
IDF3UG=06
IFLG8=1
K=0
18 K=K+1
IF(K.GT.4)GO TO 60
EPS(12)=X(K)
EPS(13)=Y(K)
C-----
C COMPUTE INCREMENTS
C-----
17 EPS(6) = SLAM6 * EPS(12) + SMU6 * EPS(13) + SNU6
EPS(7) = SETA7 * EPS(12) + SIOT7 * EPS(13) + SKAP7
EPS(9) = SALP9 * EPS(12) + SBET9 * EPS(13) + SGAM9
EPS(14) = SZET14 * EPS(12) + SETA14 * EPS(13) + SIOT14
EPS(4) = SALP4 * EPS(12) + SBET4 * EPS(13) + SGAM4
EPS(3) = SALP3 * EPS(4) + SBET3 * EPS(12) + SGAM3
EPS(1) = SALP1 * EPS(9) + SBET1 * EPS(7) + SGAM1
EPS(17) = SALP17 * EPS(3) + SBET17 * EPS(4) + SGAM17
EPS(11) = SALP11 * EPS(13) + SBET11 * EPS(12)
EPS(10) = SALP10 * EPS(7)
EPS(18) = SEPS5 * EPS(17) + SZET5 * EPS(3) + SETA5 * EPS(4) +
+ SIOT5
EPS(5) = SALP18 * EPS(18) + SBET18
EPS(8) = SALP8 * EPS(7)
EPS(2) = SALP2 * EPS(17) + SBET2 * EPS(18) + SGAM2
EPS(19) = SALP19 * EPS(18) + SBET19 * EPS(17)
WRITE(10DEBUG,20) (I,I=1,24),EPS
20 FORMAT(3(/' ',B15X,'EPS(',I2,')',4X),3(/' ',B16,A))
C-----
C COMPUTE SMALL PERTURBATION PROPERTY VALUES
C-----
DO 30 I=1,19
25 J=JEXTP(I)
IF(J.EQ.0)GO TO 30
V(J,1)=V(J,2)+EPS(I)
30 CONTINUE
GO TO(341,342),IFLG9
341 PPT = PPT1 * (RPT / RPT1) ** G
TPT = PPT * OOR / RPT * A2
GO TO 343
342 TPT=TCT0
PPT=RPT*R*TPT/A2
343 PCT0 = PEO
TCT0 = TEO

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REQ = PEO * OOR / TEO * A2
RCTO = REQ
ACTO = DSQRT(GR * TCTO)
MOCTO = RCTO * ACTO * CTA
MD = MACH(PD)
MN = MACH(PN)
IF((INSTR(23).NE.2).AND.(INSTR(24).EQ.0))GO TO 28
J16=J16+1
CALL PRINT
J16=J16-1
WRITE(IOUT,29)
29 FORMAT(1+Q1)
28 IF((IFLGB.EQ.0).AND.(IDEBUG.EQ.03))RETURN
WRITE(IDEBUG,32)V
32 FORMAT('0SMALL PERTURBATION PROPERTIES',13(/' ',8E16,A))
C-----
C SMALL PERTURBATION RESIDUALS
C-----
DO 35 I=1,19
35 Q(I)=1.070
Q(1) = CA1 * EPS(1) + CB1 * EPS(9) + CC1 * FAE + CD1 * EPS(10)
Q(2) = CA2 * EPS(2) + CB2 * EPS(17) + CC2 * FAFE + CD2 * EPS(19)
Q(3) = CA3 * EPS(3) + CB3 * EAF + CC3 * EPS(17) + CD3 * EPS(12)
Q(4) = CA4 * EPS(4) + CB4 * EPS(17) + CC4 * EPS(11)
Q(5) = CA5 * EPS(5) + CB5 * EPS(2) + CC5 * EPS(3) + CD5 * EPS(4) +
@ CFS
Q(6) = CA5 * EPS(6) + CB6 * EPS(4) + CC6 * EPS(7)
Q(7) = CA7*EPS(6) + CB7 * EPS(3) + CC7 * EPS(1)
Q(8) = CA8 * EPS(7) + CB8 * EPS(8)
Q(9) = CA9 * EPS(9) + CB9 * EPS(8)
Q(10) = CA10 * EPS(10) + CB10 * EPS(8)
Q(11) = CA11 * EPS(11) + CB11 * EPS(13) + CC11 * EPS(12)
B(15) = PF01 * EPS(12) - PD1 * EPS(9)
Q(12) = CA12 * EPS(6) + CB12 * EPS(14) + CC12 * B(15) + CD12 *
@ B(15) ** 2
B(16) = PF01 * EPS(13) - PN1 * EPS(9)
Q(13) = CA13 * EPS(7) + CB13 * EPS(14) + CC13 * B(16) + CD13 *
@ B(16) ** 2
Q(14) = CA14 * EPS(14) + CB14 * EPS(9) + CC14 * EPS(10)
Q(17) = CA17 * EPS(17) + CB17 * EPS(18)
Q(18) = CA18 * EPS(18) + CB18 * EPS(5) + CC18
Q(19) = CA19 * EPS(19) + CB19 * EPS(18) + CC19 * EPS(17)
WRITE(IDEBUG,13)(I,I=1,20),Q
13 FORMAT('0RESIDUALS FROM SMALL PERTURBATION EQUATIONS',
1 2(/' ',10(6X,I2,5X)),2(/' ',10E13.5))
C-----
C EXACT RESIDUALS
C-----
DO 140 I=1,19
140 Q(I)=1.070
Q(1)=MDF-A1*PEO*AE/DSQRT(TEO)*A2
Q(2)=MDPE+A1*PP*AFE/DSQRT(TPI)*A2
Q(3)=MDF+AF*OOKF*(PP-PD)*A2
Q(3)=MDF+AF*OOKF*(PP-PD*A16)*A2
Q(4)=MDPT+AWOKW*(PP-PT*A15)*A2
Q(5)=RPT-RPT1-(MDPE+MDF+MDPT)*DTPV

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Q(6)=MND+MDPT-MDCT
Q(7)=MDD-MDF-MDE
B(9)=1.00+GM102*MCT
B(10)=1.00+GM102*MCT**2
Q(8)=MDCT-MCT*MDCTC*B(9)**MGROGM
B(11)=B(10)/B(9)**2
Q(9)=PF0-PC*B(11)**GOGM1
Q(10)=TE0-TC*B(11)
IF(IFLG2)144,9999,145
144 Q(11)=PT-(1.00-A17)*PN-A17*PD
GO TO 146
145 Q(11)=PT-0.5D0*(PN+PD)
146 B(12)=1.00/PE0
B(13)=TOGM1*MDTS0**2
Q(12)=MND**2-B(13)*((PD*B(12))**TOG-(PD*B(12))**GP10G)
Q(13)=MDCT**2-B(13)*((PN*B(12))**TOG-(PN*B(12))**GP10G)
Q(14)=MDTS0-SGOR*PE0*TS4/DSQRT(TE0)*A2
GO TO(141,142),IFLG9
141 Q(17)=PP-PPT1*(RP/RPT1)**G
GO TO 143
142 Q(17)=PP-RP*R*TP/A2
143 Q(18)=RP-0.5D0*(RPT+RPT1)
Q(19)=TP-PP*00R/RP*A2
C*****
C OUTPUT
C*****
C-----
C CONVERT RESIDUALS TO PERCENTAGES
C-----
39 DO 49 I=1,19
J=I*EXP(I)
IF(J.EQ.0)GO TO 49
IF(I.EQ.6)GO TO 45
IF(I.EQ.7)GO TO 41
IF(I.EQ.8)GO TO 42
IF(I.EQ.12)GO TO 43
IF(I.EQ.13)GO TO 42
GO TO 46
41 J=10
GO TO 46
42 J=12
GO TO 46
43 J=9
GO TO 46
45 J=11
46 IF(V(J,2).NE.0.00)GO TO 47
Q(I)=INFIN
GO TO 49
47 Q(I)=Q(I)*1.02/V(J,2)
49 CONTINUE
WRITE(IDEBUG,21)(I,I=1,20),0
21 FORMAT('0RESIDUALS FROM EXACT EQUATIONS',
1 2(' ',10(6X,I2,5X)),2(' ',10E13.5))
IF(IFLGR.EQ.0)GO TO 60
IF(IFLG2)220,9999,18
60 WRITE(IDEBUG,31)(I,I=1,10),A:(I,I=1,30),B

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31 FORMAT('0A ARRAY'/' ',10(6X,I2,5X))/' ',10E13.5/'0','R ARRAY',
13(/' ',10(6X,I2,5X),3(/' ',10E13.5))
WRITE(IDEBUG,10) CA1, CB1, CC1, CD1, CA2, CB2, CC2, CD2,
1 CA3, CB3, CC3, CD3, CA4, CB4, CC4, CA5, CB5, CC5, CD5,
2 CE5, CA6, CB6, CC6, CA7, CB7, CC7, CA8, CB8, CA9, CB9,
3 CA10, CB10, CA11, CB11, CC11, CA12, CB12, CC12, CD12, CA13, CB13,
4 CC13, CD13, CA14, CB14, CC14, CA17, CB17, CA18, CB18, CC18, CA19,
5 CB19, CC19
10 FORMAT('15MPERT'/' ',6(=)/'0',7X,'CA1',13X,'CB1',13X,'CC1',
1 13X,'CD1'/' ',4E16.8/'0',7X,'CA2',13X,'CB2',13X,'CC2',13X,
2 'CD2'/' ',4E16.8/'0',7X,'CA3',13X,'CB3',13X,'CC3',13X,'CD3',
3 '/' ',4E16.8/'0',7X,'CA4',13X,'CB4',13X,'CC4'/' ',3E16.8/'0',
4 7X,'CA5',13X,'CB5',13X,'CC5',13X,'CD5',13X,'CE5'/' ',5E16.8
5 '/'0',7X,'CA6',13X,'CB6',13X,'CC6'/' ',3E16.8/'0',7X,'CA7',13X,
6 'CB7',13X,'CC7'/' ',3E16.8/'0',7X,'CA8',13X,'CB8'/' ',2E16.8
7 '/'0',7X,'CA9',13X,'CB9'/' ',2E16.8/'0',6X,'CA10',12X,'CB10'/' ',
8 2E16.8/'0',6X,'CA11',11X,'CB11',12X,'CC11'/' ',3E16.8/'0',
9 6X,'CA12',12X,'CB12',12X,'CC12',12X,'CD12'/' ',4E16.8/'0',6X,
A 'CA13',12X,'CB13',12X,'CC13',12X,'CD13'/' ',4E16.8/'0',6X,
B 'CA14',12X,'CB14',12X,'CC14'/' ',3E16.8/'0',6X,'CA17',12X,
C 'CB17'/' ',2E16.8/'0',6X,'CA18',12X,'CB18',12X,'CC18'/' ',
D 3E16.8/'0',6X,'CA19',12X,'CB19',12X,'CC19'/' ',3E16.8)
WRITE(IDEBUG,11) SALP1, SBET1, SGAM1, SALP2, SBET2, SGAM2, SALP3
1 , SBET3, SGAM3, SEPS3, SALP4, SBET4, SGAM4, SEPS4, SALP5, SBET5
2 , SGAM5, SEPS5, SZET5, SETA5, SIOT5, SKAP5, SALP6, SBET6, SGAM6
3 , SEPS6, SZET6, SETA6, SIOT6, SKAP6, SGAM6, SMU6, SNU6, SALP7
4 , SBET7, SGAM7, SEPS7, SZET7, SETA7, SIOT7, SKAP7, SALP8, SALP9
5 , SBET9, SGAM9, SZET9, SALP10, SALP11, SBET11, SALP12, SBET12, SGAM12
6 , SALP13, SBET13, SGAM13, SALP14, SBET14, SGAM14, SEPS14, SZET14, SETA14
7 , SIOT14, SALP17, SBET17, SGAM17, SEPS17, SALP18, SBET18, SALP19, SBET19
11 FORMAT('1',6X,'SALP1',11X,'SBET1',11X,'SGAM1'/' ',3E16.8/'0',6X,
1 'SALP2',11X,'SBET2',11X,'SGAM2'/' ',3E16.8/'0',6X,'SALP3',11X,
2 'SBET3',11X,'SGAM3',11X,'SEPS3'/' ',4E16.8/'0',6X,'SALP4',11X,
3 'SBET4',11X,'SGAM4',11X,'SEPS4'/' ',4E16.8/'0',6X,
4 'SALP5',11X,'SBET5',11X,'SGAM5',11X,'SEPS5',11X,'SZET5',11X,
5 'SETA5',11X,'SIOT5',11X,'SKAP5'/' ',8E16.8/'0',6X,'SALP6',11X,
6 'SBET6',11X,'SGAM6',11X,'SEPS6',11X,'SZET6',11X,'SETA6',11X,
7 'SIOT6',11X,'SKAP6'/' ',8E16.8/'0',6X,'SLAM6',11X,'SMU6',11X,
8 'SNU6'/' ',3E16.8/'0',6X,'SALP7',11X,'SBET7',11X,'SGAM7',11X,
9 'SEPS7',11X,'SZET7',11X,'SETA7',11X,'SIOT7',11X,'SKAP7'/' ',
A 8E16.8/'0',6X,'SALP8'/' ',F16.8/'0',6X,'SALP9',11X,'SBET9',
B 11X,'SGAM9',11X,'SZET9'/' ',4E16.8/'0',5X,'SALP10'/'
C ' ',F16.8/'0',5X,'SALP11',10X,'SBET11'/' ',2E16.8/'0',5X,'SALP12
D ' ',10X,'SBET12',10X,'SGAM12'/' ',3E16.8/'0',5X,'SALP13',10X,
E 'SBET13',10X,'SGAM13'/' ',3E16.8/'0',5X,'SALP14',10X,'SBET14',
F 10X,'SGAM14',10X,'SEPS14',10X,'SZET14',10X,'SETA14',10X,'SIOT14',
G '/' ',7E16.8/'0',5X,'SALP17',10X,'SBET17',10X,'SGAM17',10X,
H 'SEPS17'/' ',4E16.8/'0',5X,'SALP18',10X,'SBET18'/' ',2E16.8/
I '0',5X,'SALP19',10X,'SBET19'/' ',2E16.8)
WRITE(IDEBUG,12) SA2, SB2, SC2, SD2, SE2, SF2, SA3, SB3, SC3, SD3, SE3, SF3
1 , FAE, FAPE, FAF
12 FORMAT('1',7X,'SA2',13X,'SB2',13X,'SC2',13X,'SD2',13X,'SE2',
1 13X,'SF2'/' ',6E16.8/'0',7X,'SA3',13X,'SB3',13X,'SC3',13X,
2 'SD3',13X,'SE3',13X,'SF3'/' ',6E16.8/'0',7X,'FAE',12X,'FAPE',
3 13X,'FAF'
4 '/' ',8E16.8)
IF (IFLAG.EQ.1) STOP
RETJRN
9999 STOP
END

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                                QSIMUL                                DATE = 75157                                11/58/40
SUBROUTINE QSIMUL(A2,B2,C2,D2,F2,F2,A3,B3,C3,D3,F3,F3,INDEB,XX,Y)
IMPLICIT REAL*8(A-H,O-Z)
COMPLEX*16 X(2,2,4),Y(4),SIGMA,UPSLO,PHI,PSI,OMEGA,ZERO,A1,B1,C1,
-Y1,Y2,Y3,Y4,ONE,R(2,2,4,2),CINFIN,XX(4)
COMPLEX*16 SIGMA1,UPSLO1,PHI1,PSI1,OMEGA1
DIMENSION III(4),JJJ(4)
DATA ZFRO/(0.00,0.00)/,ONE/(1.00,0.00)/,CINFIN/(1.070,1.070)/
C-----
C COMPUTE QUARTIC COEFFICIENTS
C-----
ALPHI=A2*F3-A3*F2
BETAI=A2*D3-A3*D2
GAMMI=A2*F3-A3*F2
ALPHII=A3*B2-A2*B3
BETAI1=A3*C2-A2*C3
DELTII=ALPHI*BETAI1
FPSLII=A1*PHI*ALPHII+BETAI*BETAI1
ZETAI1=BETAI*ALPHII+GAMMI*BETAI1
ETAI1=GAMMI*ALPHII
BIISQ=BETAI1**2
AII=ALPHII*BETAI1
AIISQ=ALPHII**2
SIGMA=ZFRO
UPSLO=ZERO
PHI=ZERO
PSI=ZERO
OMEGA=ZFRO
DO 8 I=1,2
DO 7 J=1,2
DO 6 K=1,4
6 X(I,J,K)=CINFIN
7 CONTINUE
8 CONTINUE
SIGMA= A2 * ALPHI**2 + C2 * DELTII + E2 * BIISQ
UPSLO = 2.00 * A2 * ALPHI * BETAI + B2 * DELTII + C2 * FPSLII +
* 2.00 * E2*AI1 + D2 * BIISQ
PHI = A2 * ( BETAI ** 2 + 2.00 * ALPHI * GAMMI ) + B2 * FPSLII +
* C2 * ZETAI1 + F2 * AIISQ + F2 * BIISQ + 2.00 * D2 * ABI1
PSI = 2.00 * BETAI * GAMMI *A2 + B2 * ZETAI1 + C2 * ETAI1 +
* 2.00 * F2 * AII + D2 * AIISQ
OMEGA = A2 * GAMMI ** 2 + B2 * ETAI1 + F2 * AIISQ
SIGMA1=SIGMA
UPSLO1=UPSLO
PSI1=PSI
PHI1=PHI
OMEGA1=OMEGA
IF (INDEB.EQ.03)GO TO 9
C-----
C PRINT COEFFICIENTS
C-----
WRITE(INDEB,1)A2,B2,C2,D2,F2,F2,A3,B3,C3,D3,F3,F3,ALPHI,BETAI,
- GAMMI,ALPHII,BETAI1,DELTII,FPSLII,ZETAI1,ETAI1,BIISQ,AII,AIISQ
1 FORMAT('QSIMUL',' ',6(I=1)/'0',7X,'A2',14X,'B2',14X,'C2',14X,
1 'D2',14X,'F2',14X,'F2',' ',6E16,8/'0',7X,'A3',14X,'B3',14X,'C3',
2 14X,'D3',14X,'F3',14X,'F3',' ',6F16,8/'0',6X,'ALPHI',11X,
3 'BETAI',11X,'GAMMI',10X,'ALPHII',10X,'BETAI1',10X,'DELTII',

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4 10X,'FPSLII'/' ' ,7E16,8/'0',5X,'ZETAI',11X,'ETAI',11X,
5 'RIISQ',11X,'ARII',12X,'AIISQ'/' ' ,5E16,8)
WRITE(IDEBUG,2)SIGMA,UPSILON,PHI,PSI,OMEGA,ZERO
2 FORMAT('0',13X,'SIGMA',27X,'UPSILON',28X,'PHI',29X,'PHI'/' ' ,
18E16,8/'0',14X,'OMEGA',27X,'ZERO'/' ' ,4E16,8)
C-----
C FIND ROOTS TO QUARTIC
C-----
9 N=4
IF(DREAL(SIGMA),NE,0.00)GO TO 40
N=3
IF(DREAL(UPSILON),NE,0.00)GO TO 30
N=2
IF(DREAL(PHI),NE,0.00)GO TO 20
C LINER
N=1
10 OMEGA=OMEGA/PSI
PSI=ONE
CALL QANDC(N,OMEGA,ZERO,ZERO,ZERO,ZERO,Y1,Y2,Y3,Y4)
GO TO 45
C QUADRATIC
20 PSI=PSI/PHI
OMEGA=OMEGA/PHI
PHI=ONE
CALL QANDC(N,PSI,OMEGA,ZERO,ZERO,Y1,Y2,Y3,Y4)
GO TO 45
C CURTIC
30 PHI=PHI/UPSILON
PSI=PSI/UPSILON
OMEGA=OMEGA/UPSILON
UPSILON=ONE
CALL QANDC(N,PHI,PSI,OMEGA,ZERO,Y1,Y2,Y3,Y4)
GO TO 45
C QUARTIC
40 UPSILON=UPSILON/SIGMA
PHI=PHI/SIGMA
PSI=PSI/SIGMA
OMEGA=OMEGA/SIGMA
SIGMA=ONE
CALL QANDC(N,UPSILON,PHI,PSI,OMEGA,Y1,Y2,Y3,Y4)
C-----
C FIND ALL X VALUES FOR EACH Y ROOT
C-----
45 Y(1)=Y1
Y(2)=Y2
Y(3)=Y3
Y(4)=Y4
DO 100 I=1,4
IF(I.GT,N)GO TO 100
B1=-.5*(B2+C2*Y(I))/A2
C1=COSQRT(B1**2-(D2*Y(I)+F2*Y(I)**2+F2)/A2)
X(1,1,I)=B1+C1
X(2,1,I)=B1-C1
B1=-.5*(B3+C3*Y(I))/A3
C1=COSQRT(B1**2-(D3*Y(I)+F3*Y(I)**2+F3)/A3)
X(1,2,I)=B1+C1

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QSIMUL

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      X(2,2,T)=R1-C1
100 CONTINUE
C-----
C PRINT ROOTS TO QUARTIC AND CONICS
C-----
      WRITE(IDEBUG,2) SIGMA, UPSILON, PHI, PSI, OMEGA, ZERO
      WRITE(IDEBUG,3) Y, (T, I=1,2), X
3  FORMAT('0',15X,'Y1',30X,'Y2',30X,'Y3',30X,'Y4',/' ',8E16.8/'0',
1 2(')',26(' '), 'EQUATION',12,27(' '), '-1',/' ',2(')',11(' '),
2  'POSITIVE',11(' '), '1',11(' '), 'NEGATIVE',12(' '), '-1')
3/'0X ROOTS BASED ON Y1',/' ',8E16.8/'0X ROOTS BASED ON Y2',/' ',8
4  E16.8/'0X ROOTS BASED ON Y3',/' ',8E16.8/'0X ROOTS BASED ON Y4',/
5  ' ',8E16.8)
C-----
C CHECK ALL X AND Y VALUES IN ORIGINAL SYSTEM OF CONICS
C-----
      DO 130 L=1,2
      DO 120 I=1,4
      DO 114 J=1,2
      DO 115 K=1,2
      IF (DREAL(Y(I)).GT.1.D69) GO TO 115
      GO TO (108,110), L
108 R(J,K,I,1)=A2*X(J,K,I)**2+B2*X(J,K,I)+C2*X(J,K,I)*Y(I)+D2*Y(I)+
1  E2*Y(I)**2+F2
      GO TO 115
110 R(J,K,I,2)=A3*X(J,K,I)**2+B3*X(J,K,I)+C3*X(J,K,I)*Y(I)+D3*Y(I)+
1  F3*Y(I)**2+F3
115 CONTINUE
114 CONTINUE
120 CONTINUE
130 CONTINUE
      WRITE(IDEBUG,4) R
      4  FORMAT('0RESIDUAL ARRAY',4('/' ',8E16.8))
C-----
C SORT OUT EXTRANEOUS ROOTS
C-----
135 L=0
      DO 160 J=1,4
      DO 180 I=1,2
      DO 140 K=1,2
      IF (CDARS(R(I,1,J,K)).LT.1.D-10) GO TO 140
      GO TO 180
140 CONTINUE
      L=L+1
      IF (L.LE.4) GO TO 150
      WRITE(IOUT,145)
145 FORMAT('0QSIMUL: MORE THAN FOUR ROOTS FOUND')
      STOP
150 IIT(L)=I
      JJJ(L)=J
180 CONTINUE
160 CONTINUE
      LLL=L
      DO 200 L=1,4
      XX(L)=CTNFIN
      IF (L.GT.LLL) GO TO 190
      XX(L)=X(IIT(L),1,JJJ(L))
      GO TO 200
190 Y(L)=CTNFIN
200 CONTINUE
      WRITE(IDEBUG,220) XX,Y
220 FORMAT('0SORTED ROOTS: XX/Y',2('/' ',8E16.8))
      WRITE(IDEBUG,240) IIT, JJJ
240 FORMAT('0III AND JJJ VALUES FOR R(III,1,JJJ,K),/' ', 'III=',4I2/,
1  ' ', JJJ=',4I2)
      RETURN
      END

```

QANDC

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SUBROUTINE QANDC(N,R,C,D,F,X1,X2,X3,X4)
IMPLICIT COMPLEX*16(A,G,O-Z)
COMPLEX*16 I
DATA I/(0.00,1.00)/,CINFIN/(1.070,1.070)/
GO TO(30,20,10,5),N
C QJARTIC
5 A=((4.00*C*F-(R**2*F)-D**2)/2.00)
A1=(C*(R**2-4.00*E1)/6.00)
A2=-((C**3)/27.00)
A=A+A1+A2
R9=CDSQRT((A**2)+((R*D-4.00*F-((C**2)/3.00))**3)/27.00)
A=-A
CALL CHIART(A,R9,R)
PSTAR=R*D-4.00*F-(C**2)/3.00
R1=-PSTAR/(3.00*R)
R=(R+R1+(C/3.00))
P=CDSQRT((R**2/4.00)-C+R)
PQ=CDSQRT(0.2500*R**2-E)
AR2=.500*3*R=I
PPQ2=2.00*P*PQ
IF(CDARS(AR2-PPQ2).GT.CDARS(AR2+PPQ2))PQ=-PQ
PP=(CDARS(P))
C CALCULATING THE ZEROS
A1=(1.0,0.0)
R1=(R/2.00)+P
C1=(R/2.00)+PQ
X1=(-R1+CDSQRT(R1**2-4.00*A1*C1))/(2.00*A1)
X2=(-R1-CDSQRT(R1**2-4.00*A1*C1))/(2.00*A1)
R1=(R/2.00)-P
C1=(R/2.00)-PQ
X3=(-R1+CDSQRT(R1**2-4.00*A1*C1))/(2.00*A1)
X4=(-R1-CDSQRT(R1**2-4.00*A1*C1))/(2.00*A1)
RETURN
C CJRIC
10 CONTINUE
P=C-(R**2/3.00)
Q=D-(R*(C/3.00)+((2.00*R**3)/27.00)
Z1=-(Q/2.00)+(CDSQRT((Q**2/4.00)+(P**3/27.00)))
Z2=-(Q/2.00)-(CDSQRT((Q**2/4.00)+(P**3/27.00)))
IF(CDARS(Z1).GE.CDARS(Z2))Z=Z1
IF(CDARS(Z2).GE.CDARS(Z1))Z=Z2
IF(CDARS(Z).EQ.0.0)X1=-(R/3.00)
IF(CDARS(Z).EQ.0.0)X2=-(R/3.00)
IF(CDARS(Z).EQ.0.0)X3=-(R/3.00)
IF(CDARS(Z).EQ.0.0)RETURN
RRR=(0.00,0.00)
CALL CHIART(Z,RRR,R1)
R=-(P/(3.00*R1))
W1=-(.500)+((3.00**5)/2.00)*I
W2=-(.500)-((3.00**5)/2.00)*I
X1=-(R/3.00)+R1+W1
X2=-(R/3.00)+R1+W2
X3=-(R/3.00)+R1+W1
X4=CINFIN
RETURN
C QJARTIC
20 A1=.500
R=CDSQRT(A1**2-C)
X1=A1+R
X2=A1-R
X3=CINFIN
X4=CINFIN
RETURN
C LINPAR
30 X1=-H
X2=CINFIN
X3=CINFIN
X4=CINFIN
RETURN
END

```

```

                                CUBRT          DATE = 75157      11/58/40
SUBROUTINE CUBRT(AA,RR,RR)
IMPLICIT COMPLEX*16(A-G,Q-Z)
REAL*8 AH,H1A,H1B,HTH
REAL*8 PI,SSSS
COMPLEX*16 I
990 CONTINUE
I=(0.,1.)
II=1
Z1=AA+RR
Z2=AA-RR
IF(CDARS(Z2).GF.CDARS(Z1))A=Z2
IF(CDARS(Z1).GF.CDARS(Z2))A=Z1
R=DCONJG(A)
H1A=(A+R)/2.D0
H1B=-I*(A-R)/2.D0
HTH=DATAN2(H1B,H1A)
H=(H1A**2+H1B**2)**.5D0
PI=3.141592653589793D0
SSSS=3.D0
RR=(H*(1.D0/3.D0))*((DCOS((HTH+(II-1)*2.D0*PI)/SSSS)+I*(DSIN((HT
H+(II-1)*2.D0*PI)/SSSS)))
RETURN
END

```

```

IN IV G LEVEL 21          DREAL          DATE = 75157      11/58/40
FUNCTION DREAL(CC)
COMPLEX*16 C,CC
REAL*8 D(2),DREAL
EQUIVALENCE (C,D(1))
C=CC
DREAL=D(1)
RETURN
END

```

```

IN IV G LEVEL 21          DIMAG          DATE = 75157      11/58/40
FUNCTION DIMAG(II)
COMPLEX*16 I,II
REAL*8 D(2),DIMAG
EQUIVALENCE (I,D(1))
I=II
DIMAG=D(2)
RETURN
END

```

NOMENCLATURE

A	Area
A_{11}	Solution weighting parameter, Eq. (25)
A_{15}	Momentum correction coefficient in wall crossflow model, Eq. (7)
A_{16}	Flap correction coefficient in the flap flow model, Eq. (8)
A_{17}	Weight used in computing test section pressure, Eq. (10)
A_i, B_i, C_i, D_i, E_i	Arrays of coefficients in the small perturbation equations
E	Computational error
F	Function
k	Flow coefficient, as in k_f and k_w
M	Mach number
M_∞	Steady, asymptotic test section Mach number
\dot{m}	Mass flow rate
\dot{m}_o	Convenient quantity with units of mass flow rate defined as $\sqrt{\frac{\gamma}{R}} \frac{P_{ct_o}}{\sqrt{T_{ct_o}}} A_{ts}$
\tilde{m}	Nondimensional mass flow rate defined as $\sqrt{\frac{\gamma-1}{2}} \frac{\dot{m}}{\dot{m}_o}$, Eq. (B-3)
\hat{m}	Nondimensional mass flow rate defined as \dot{m}/\dot{m}_c , Eq. (B-1)
\dot{m}_c	Convenient quantity with units of mass flow rate defined as $\sqrt{\frac{\gamma}{R}} \frac{P_c}{\sqrt{T_c}} A_{ct}$
n	Iteration number
P	Pressure
\tilde{P}	Nondimensional pressure, P/P_o

R	Perfect gas constant
T	Temperature
t	Time
t*	Midpoint of a time interval
t _F	Final time in an area time curve
V	Volume, as in V _p or V _{ts}
v _i	Scratch variable used to develop small perturbation expansion, Eq. (27)
X,Y	Variables in numerical reversion procedure (Fig. 10)
a	Constant defined as $\sqrt{\left(\frac{2}{\gamma+1}\right)^{\frac{\gamma+1}{\gamma-1}} \frac{\gamma}{R}}$, Eqs. (4) and (5)
γ	Ratio of specific heats
δ_f	Flap gap
ϵ_i	Array of small perturbations of the variables from the exact solution (Table A-1)
ϵ_{A_e}	Perturbation in the main valve area
ϵ_{A_f}	Perturbation in the flap area
$\epsilon_{A_{pe}}$	Perturbation in the plenum exhaust valve area
ρ	Density
τ	Porosity, percent of test section wall area drilled out to allow crossflow

SUBSCRIPTS

c	Charge condition
ct	Charge tube (or supply tube)
d	Diffuser end of test section
e	Main tunnel exit, main valves

f	Flaps
i	Array index
max	Maximum value as in $A_{pe_{max}}$
n	Nozzle end of test section
p	Plenum
pe	Plenum exhaust
pt	Plenum - Test Section
t, ts	Test section, as in P_t or A_{ts}
tsw	Test section wall, as in A_{tsw} , the total wall area
w	Test section wall, as in A_w , the effective flow area
0	Stagnation condition
1	Test value in numerical reversion (Fig. 10)

SUPERSCRIPT

*	Sonic conditions
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